Understanding (PDF) Uncertainties

A personal overview of parton distribution function determinations

Alberto Guffanti NBIA & Discovery Center, Niels Bohr Institute - University of Copenhagen



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Motivation

Experimental push

 Uncertainties on Parton Distribution Functions (PDFs) are often the limiting factor in precision Standard Model studies and New Physics searches

	σ (8 TeV)	un	certainty
gg→H	19.5 pb	14.7%	
VBF	1.56 pb	2.9%	
WH	0.70 pb	3.9%	scale
ZH	0.39 pb	5.1%	
ttH	0.13 pb	14.4%	

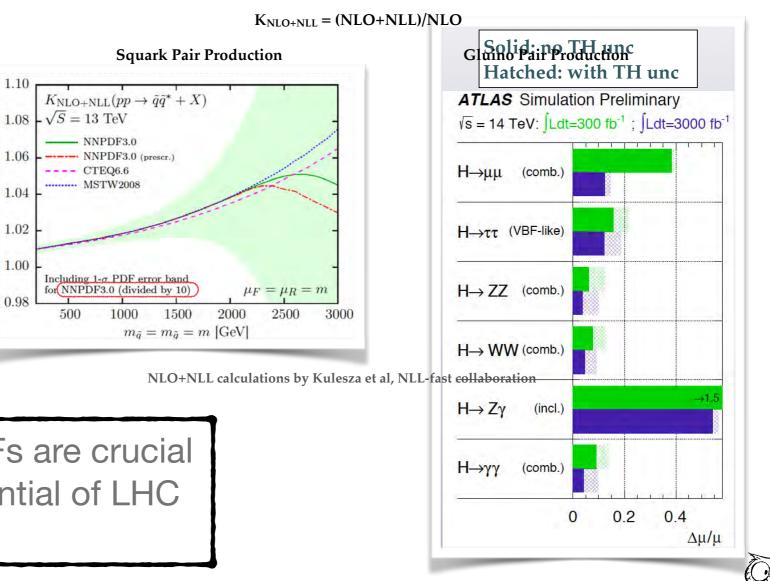
PDF uncertainties are a crucial input at the LHC, often being the limiting factor in the accuracy of theoretical

Accurate and reliable PDFs are crucial for exploiting the full potential of LHC experiments

1)

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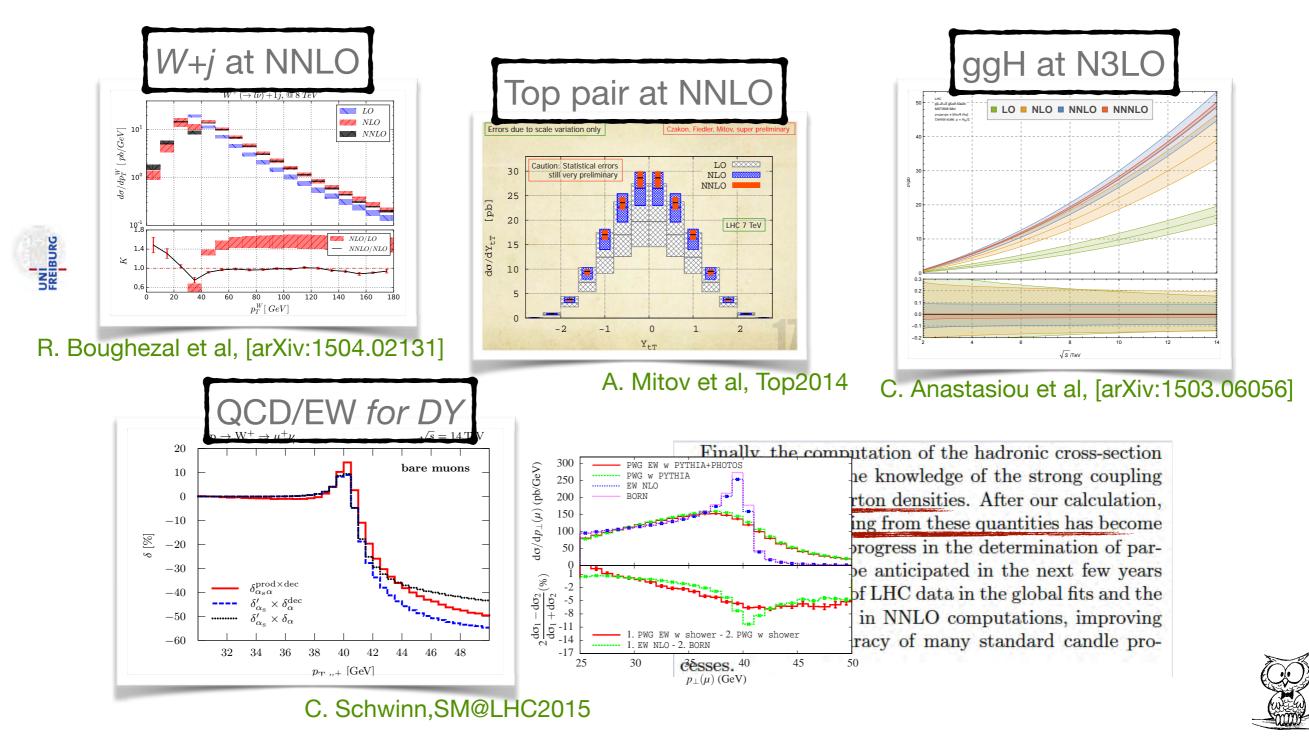
2) Very large PDF uncertainties (>100%) for BSM heavy particle production



Motivation

Theoretical push

• Impressive progress in computation of higher order corrections (QCD & EW)

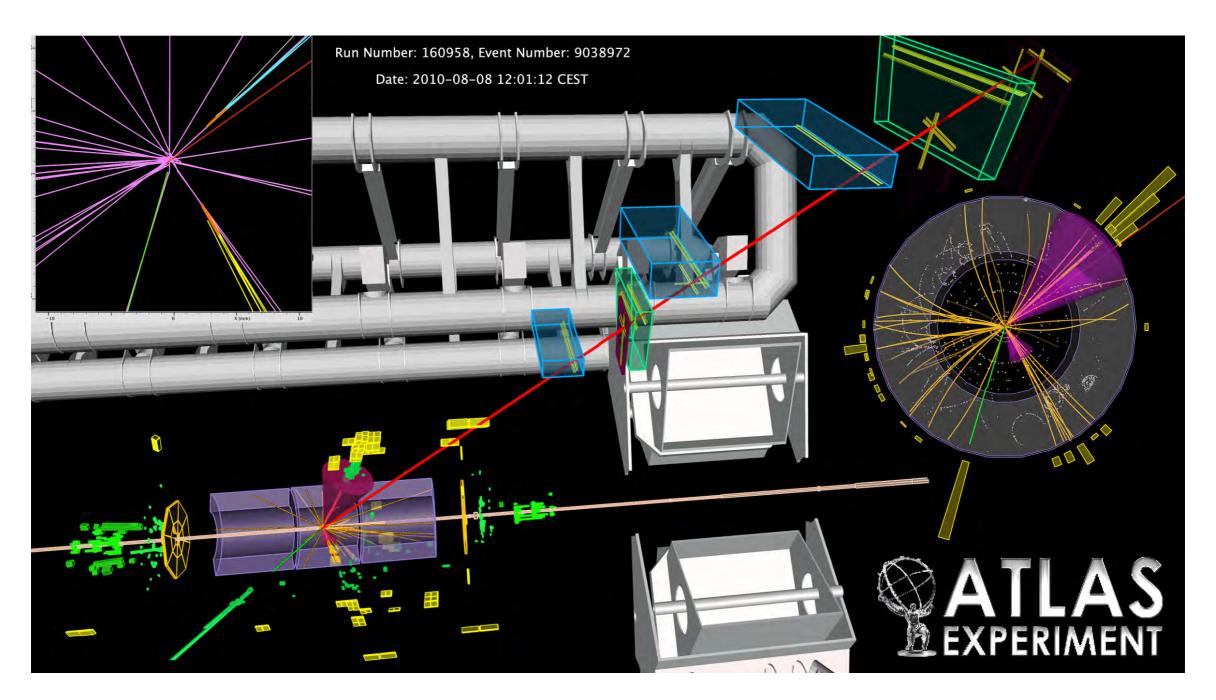




Basics of PDF determinations

LHC events

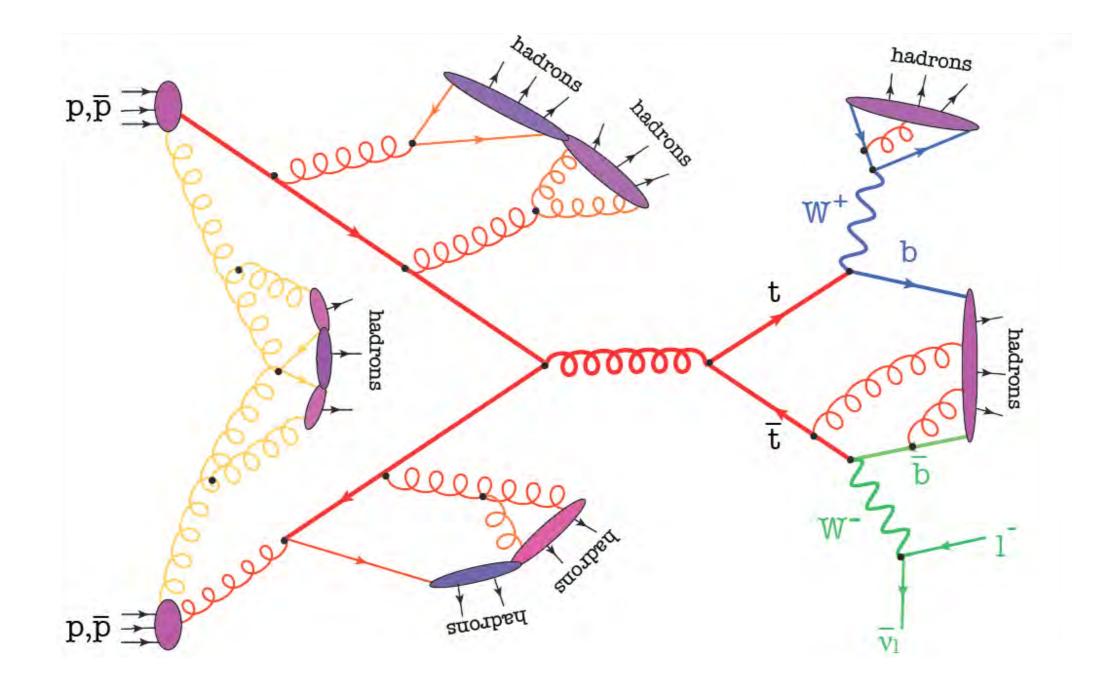
An experimentalist view





LHC events

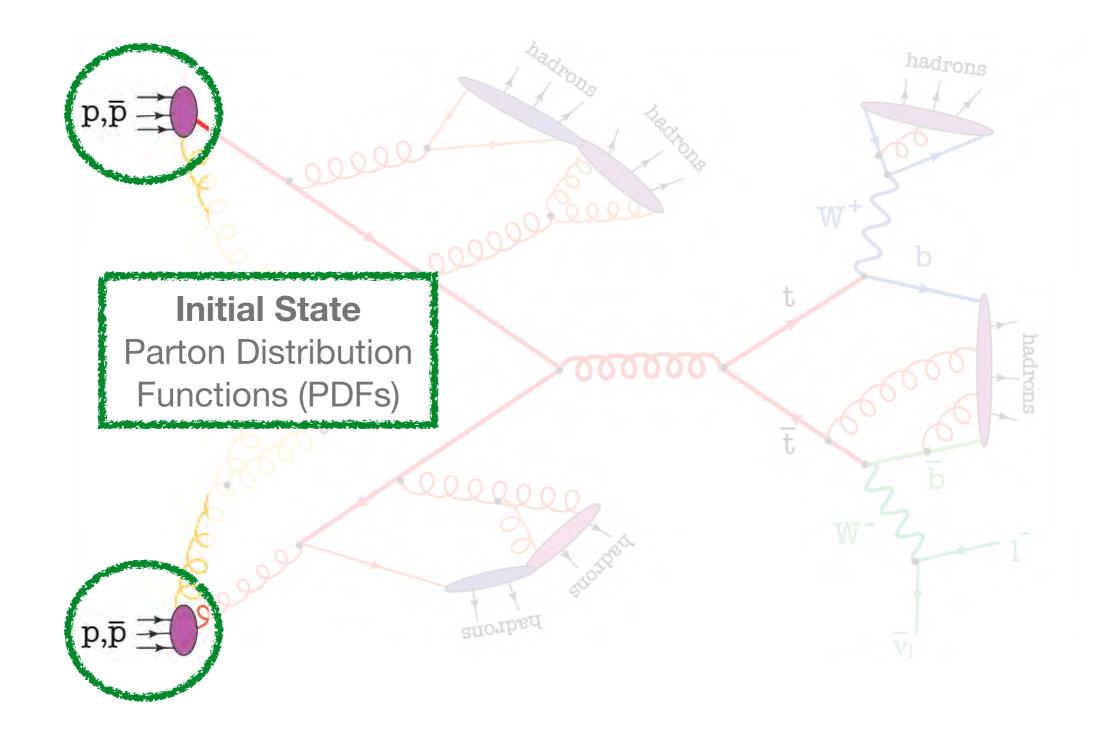
A theorist view





LHC events

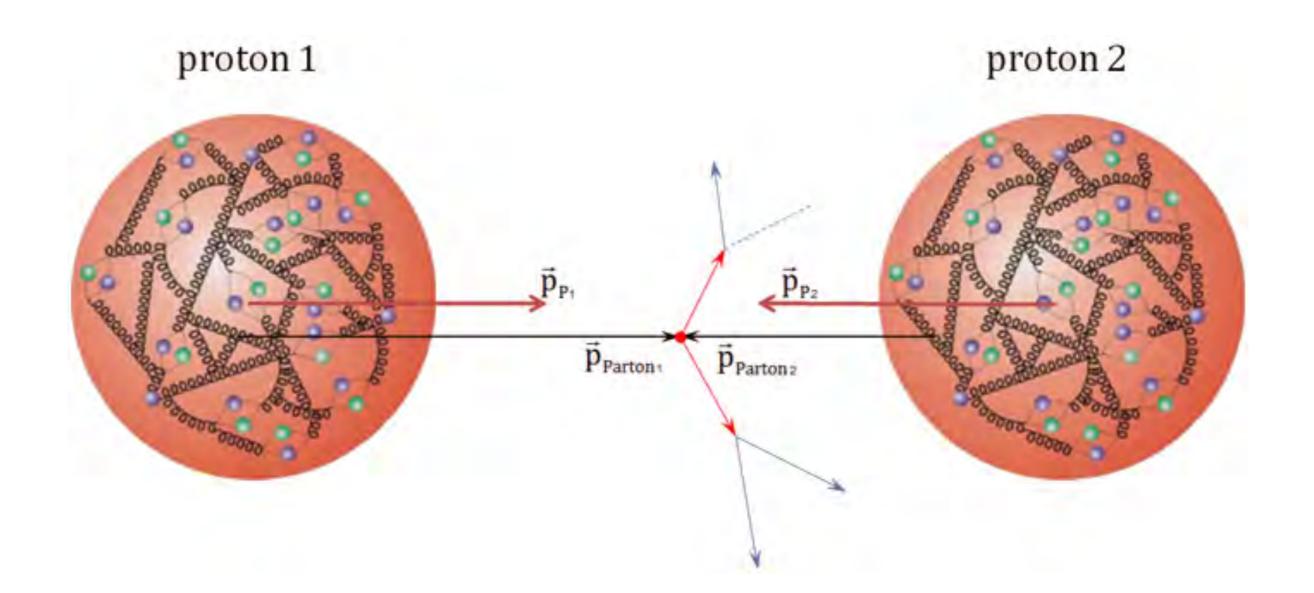
A theorist view





Parton Distribution Functions at the LHC

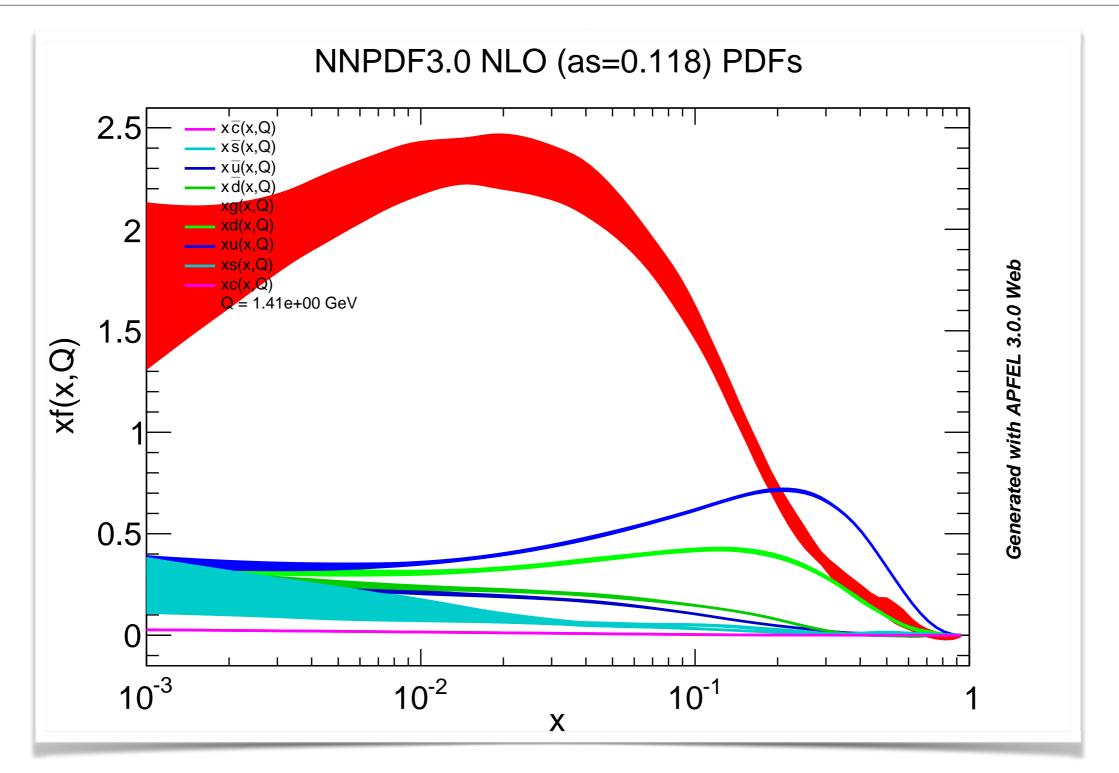
A phenomenologist view





Parton Distribution Functions at the LHC

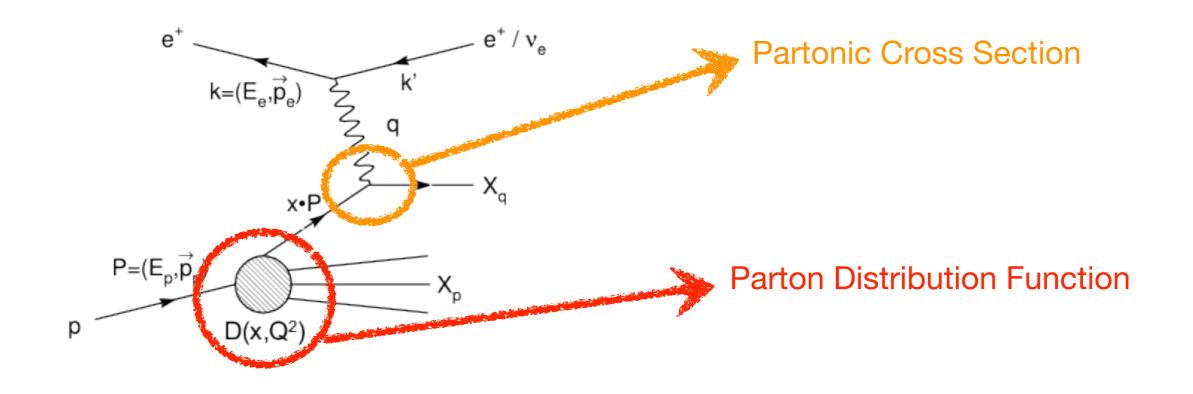
A PDF fitter view





Collinear Factorization

• Consider a process with one hadron in the initial state (Deep Inelastic Scattering)



• The cross-section can be written as (Factorization Theorem)

$$d\sigma = \sum_{a} \int_{0}^{1} \frac{d\xi}{\xi} D_{a}(x,\mu^{2}) d\hat{\sigma}_{a}\left(\frac{x}{\xi},\frac{\hat{s}}{\mu^{2}},\alpha_{s}(\mu^{2})\right) + O\left(\frac{1}{Q^{p}}\right)$$



DGLAP Evolution

- Parton Distribution Functions are non-perturbative objects and their numerical value at a given x and Q² cannot be computed in perturbative QCD (Lattice?)
- ... but their scale dependence is described by evolution equations (DGLAP)

$$\frac{\partial q_i(x,\mu^2)}{\partial \ln \mu^2} = \frac{\alpha_s(\mu^2)}{2\pi} \Big[P_{qq}(x) \otimes q_i(x,\mu^2) \Big] + \frac{\alpha_s(\mu^2)}{2\pi} \Big[P_{qg} \otimes g(x,\mu^2) \Big]$$
$$\frac{\partial g(x,\mu^2)}{\partial \ln \mu^2} = \frac{\alpha_s(\mu^2)}{2\pi} \Big[P_{gq}(x) \otimes \sum_i \left(q_i(x,\mu^2) + \overline{q}_i(x,\mu^2) \right) \Big] + \frac{\alpha_s(\mu^2)}{2\pi} \Big[P_{gg} \otimes g(x,\mu^2) \Big]$$

 where the splitting functions (P_{ij}) can be computed in QCD perturbation theory and are known up to NNLO
 [LO - Dokshitzer; Gribov, Lipatov; Altarelli, Parisi (1977)]

 [NLO - Floratos, Ross, Sachrajda; Gonzalez-Arroyo, Lopez, Yndurain; Curci, Furmanski, Petronzio (1981)]
 [NNLO - Moch, Vermaseren, Vogt (2004)]

 Moreover, Parton Distributions Functions are universal: determine them from lepton-hadron scttaring, use them for predictions in hadron-hadron collisions



DGLAP Evolution

 Parton Distribution Functions are non-perturbative objects and their numerical value at a given

HEP-EPS Prize 2015

The 2015 High Energy and Particle Physics Prize of the European Physical Society has been awarded jointly to five theoretical physicists: James D. Bjorken (SLAC National Accelerator Laboratory, Stanford, USA) "for his prediction of scaling behaviour in the structure of the proton that led to a new understanding of the strong interaction" and to Guido Altarelli (University of Roma Tre, Rome, Italy and CERN, Geneva, Switzerland), Yuri Dokshitzer (Laboratory of Theoretical and High Energy Physics, Paris, France and St. Petersburg Nuclear Physics Institute, Gatchina, Russia), Lev N. Lipatov (National Research Center "Kurchatov Institute", Petersburg Nuclear Physics Institute, Gatchina, Russia) and Giorgio Parisi (University of Rome, La Sapienza, Rome, Italy) "for developing a probabilistic field theory framework for the dynamics of quarks and gluons, enabling a quantitative understanding of high-energy collisions involving hadrons".

theory and are known up to NNLO

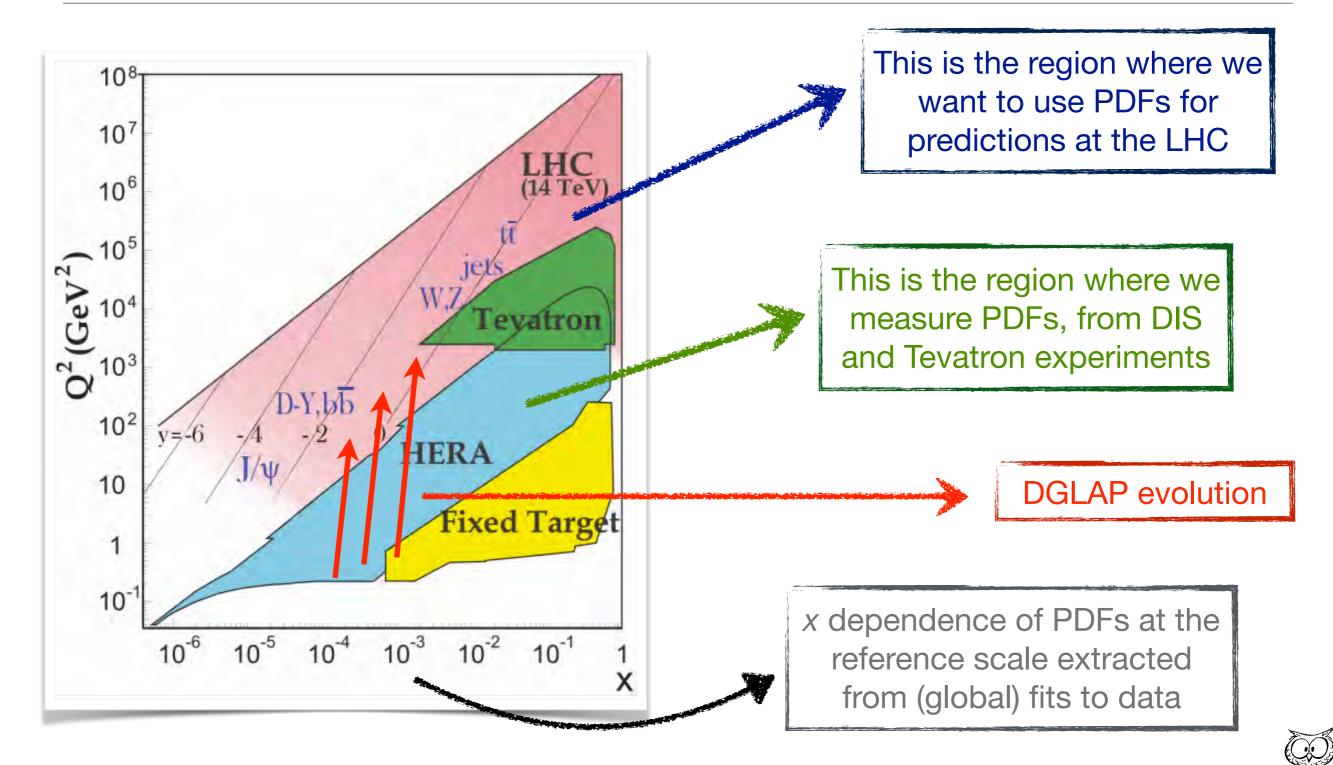
... but their scale der

[LC

NLO - Floratos, Ross, Sachrajda; Gonzalez-Arroyo, Lopez, Yndurain; Curci, Furmanski, Petronzio (1981)] [NNLO]



Where do we measure them, where do we use them



PDF determinations at the dawn of LHC Run II

State of the art, April 2015

	Dataset	Pert. Order	HQ Treatment	۵	Param.	Uncert.
ABM12 [arXiv:1310.3059]	DIS Drell-Yan	NLO NNLO	FFN (BMSN)	Fit (multiple values available)	6 indep. PDFs Polynomial (25 param.)	Hessian (Δ
CT14 [preliminary]	Global	LO NLO NNLO	GM-VFNS (S-ACOT)	External (multiple values available)	6 indep. PDFs Polynomial (27 param.)	Hessian Tolerance
HERAPDF2.0 [preliminary]	DIS (HERA I+II)	NLO NNLO	GM-VFNS (TR)	External (multiple values available)	5 indep. PDFs Polynomial (14 param.)	Hessian (Δ
MMHT14 [arXiv:1410.3989]	Global	LO NLO NNLO	GM-VFNS (TR)	Fit (multiple values available)	7 indep. PDFs Polynomial (37 param.)	Hessian Dyn. Tolerance
NNPDF3.0 [arXiv:1410.8849]	Global	LO NLO NNLO	GM-VFNS (FONLL)	External (multiple values available)	7 indep. PDFs Neural Nets (259 param.)	Monte Carlo

[LHAPDF v6.1.5 - http://lhapdf.hepforge.org/]





NNPDF Methodology Interlude

R. D. Ball, V. Bertone, S. Carrazza, C. S. Deans, L. Del Debbio, S. Forte, P. Groth-Merrild, N. P. Hartland, Z. Kassabov, J. I. Latorre, J. Rojo, M. Ubiali & AG



NNPDF Methodology ... in a Nutshell

- * Generate N_{rep} Monte Carlo replicas of the experimental data, taking into account all experimental correlations
- * Fit a set of Parton Distribution Functions, parametrized at the initial scale using Neural Networks, to each replica
- * Expectation values for observables are then given by

$$\langle \mathcal{F}[f_i(x, Q^2)]
angle = rac{1}{N_{rep}} \sum_{k=1}^{N_{rep}} \mathcal{F}\Big(f_i^{(net)(k)}(x, Q^2)\Big)$$

* and corresponding formulae for the estimators of Monte Carlo samples are used to compute uncertainties, correlations, etc.





NNPDF Methodology Main Ingredients

* Monte Carlo determination of uncertainties

- * No need to rely on linear propagation of errors
- * Possibility to test the impact of **non-gaussianly** distributed uncertainties
- * Possibility to test for non-gaussian behaviour of uncertainties of fitted PDFs
- * Parametrization of PDFs using Neural Networks
 - * Provide an unbiased parametrization
- * Determine the **best fit** PDFs using **Cross-Validation**
 - * Ensures proper fitting, avoiding overlearning





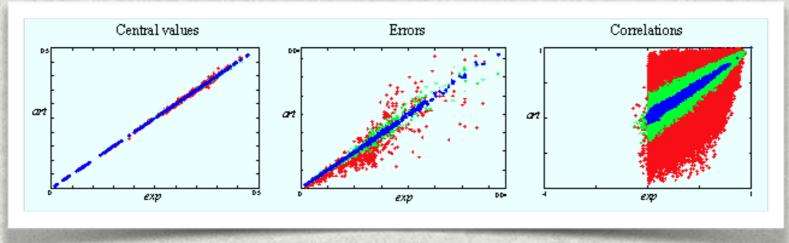
NNPDF Methodology Monte Carlo replica generation

* Monte Carlo replicas are generated according to the distribution

$$O_{i}^{(art)\,(k)} = (1 + r_{N}^{(k)}\,\sigma_{N}) \left[O_{i}^{(exp)} + \sum_{p=1}^{N_{sys}} r_{p}^{(k)}\,\sigma_{i,p} + r_{i,s}^{(k)}\,\sigma_{s}^{i} \right]$$

where r_i are (gaussianly distributed) random numbers

* Validate Monte Carlo replicas against experimental data

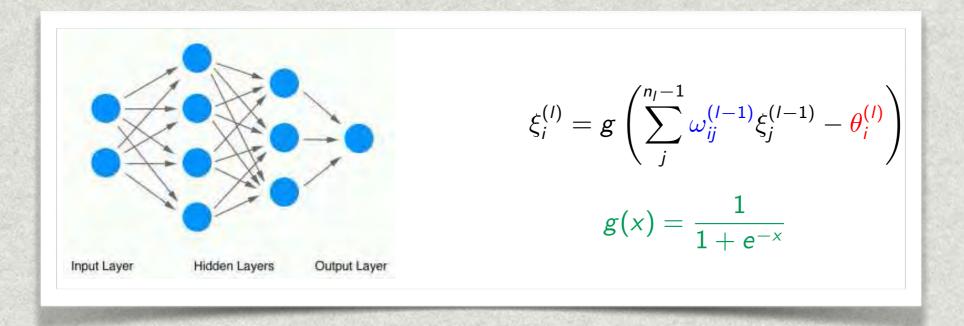


* O(1000) replicas needed to reproduce correlations in experimental data to percent accuracy





NNPDF Methodology PDF parametrization using Neural Networks



- * Artificial Neural Networks provide us with a parametrization for PDFs at the initial scale which is extremely redundant and robust against variations
- * Very efficient algorithms are available which allow us to train NN (efficient fit to large datasets in a very high dimensional parameter space)
- * ... but in the end they are just another basis of functions

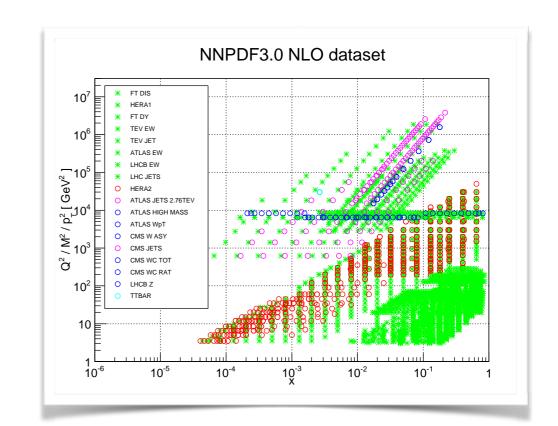




... mostly based on [arXiv:1410.8849]

Dataset

SLAC P,D DIS BCDMS P,D DIS NMC P,D DIS E665 P,D DIS	~ ~ ~	~	X
NMC P,D DIS			
·			
E665 P.D DIS	I	 Image: A set of the set of the	✓
	×	 Image: A set of the set of the	×
CDHSW NU-DIS	×	×	✓
CCFR NU-DIS	X	 Image: A second s	✓
CHORUS NU-DIS	 Image: A start of the start of	 Image: A set of the set of the	×
CCFR DIMUON	×	 Image: A set of the set of the	✓
NUTEV DIMUON	 ✓ 	 	✓
HERA I NC,CC	 	 	 ✓
HERA I CHARM	 ✓ 	 Image: A set of the set of the	✓
H1,ZEUS JETS	X	 Image: A set of the set of the	×
H1 HERA II	 ✓ 	X	×
ZEUS HERA II	✓	×	×
E605 & E866 FT DY	 	~	✓
CDF & DO W ASYM	X	~	✓
CDF & DO Z RAP	 ✓ 	 	✓
CDF RUN-II JETS	 Image: A set of the set of the	 	✓
DO RUN-II JETS	×	 	✓
ATLAS HIGH-MASS DY	 	 	✓
CMS 2D DY	 Image: A start of the start of	 	×
ATLAS W,Z RAP	 Image: A start of the start of	 	V _
ATLAS W PT	 Image: A start of the start of	 	X
CMS W ASY	 	 	 Image: A state of the state of
CMS W +c	 	×	×
LHCB W,Z RAP	 	✓	✓
ATLAS JETS	 	 	✓
CMS JETS	 	 	✓
TTBAR TOT XSEC	~	~	×
TOTAL NLO	4276	2996	3248
TOTAL NNLO	4078	2663	3045



Modern PDF sets include a substantial number of data from the LHC experiments.

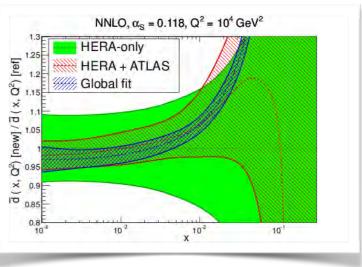


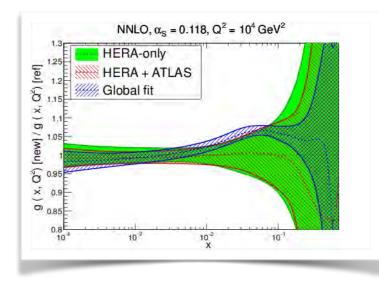
The reason for global fits

 Using a wide range of measurements coming from different experiments is crucial to constrain different PDF combinations over the whole kinematic range

Process	Subprocess	Partons	x range
$\ell^{\pm} \{p, n\} \to \ell^{\pm} X$	$\gamma^* q \rightarrow q$	q, \bar{q}, g	$x \gtrsim 0.01$
$\ell^{\pm} n/p \to \ell^{\pm} X$	$\gamma^* d/u \rightarrow d/u$	d/u	$x \gtrsim 0.01$
$pp ightarrow \mu^+ \mu^- X$	$u\bar{u}, d\bar{d} \rightarrow \gamma^*$	q	$0.015 \leq x \leq 0.3$
$pn/pp \rightarrow \mu^+\mu^- X$	$(u\bar{d})/(u\bar{u}) \rightarrow \gamma^*$	d/ū	$0.015 \leq x \leq 0.35$
$ u(ar{ u}) N ightarrow \mu^-(\mu^+) X$	$W^*q ightarrow q'$	q, \bar{q}	$0.01 \leq x \leq 0.5$
$\nu N \rightarrow \mu^- \mu^+ X$	$W^*s \rightarrow c$	5	$0.01 \leq x \leq 0.2$
$\bar{\nu} N \rightarrow \mu^+ \mu^- X$	$W^*\bar{s} \to \bar{c}$	5	$0.01 \lesssim x \lesssim 0.2$
$e^{\pm} p \rightarrow e^{\pm} X$	$\gamma^* q \rightarrow q$	g, q, \bar{q}	$0.0001 \leq x \leq 0.$
$e^+ p \rightarrow \bar{\nu} X$	$W^+ \{d, s\} \rightarrow \{u, c\}$	d,s	$x \gtrsim 0.01$
$e^{\pm}p \rightarrow e^{\pm}c\bar{c}X$	$\gamma^* c \to c, \gamma^* g \to c\bar{c}$	c, g	$0.0001 \leq x \leq 0.$
$e^{\pm}p \rightarrow \text{jet} + X$	$\gamma^*g \to q\bar{q}$	g	$0.01 \lesssim x \lesssim 0.1$
$p\bar{p} \rightarrow \text{jet} + X$	$gg, qg, qq \rightarrow 2j$	g, q	$0.01 \lesssim x \lesssim 0.5$
$p\bar{p} \to (W^{\pm} \to \ell^{\pm} \nu) X$	$ud \to W, \bar{u}\bar{d} \to W$	$u, d, \overline{u}, \overline{d}$	$x \gtrsim 0.05$
$p\bar{p} \to (Z \to \ell^+ \ell^-) X$	$uu, dd \rightarrow Z$	d	$x \gtrsim 0.05$

Parton Distributions

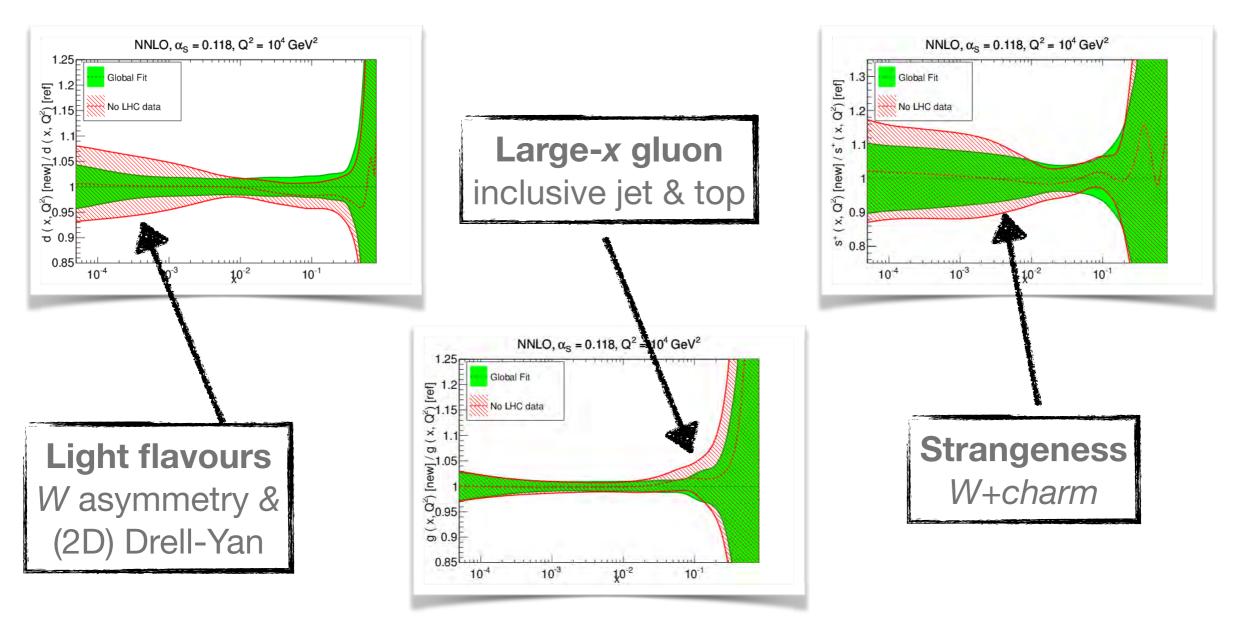






Impact of LHC data

• Impact of LHC data (still) moderate but definitely noticeable





Methodology - Closure tests

- Core idea
 - Assume the underlying PDFs are known and generate pseudodata based on a given theory set up and the chosen PDF set.
 - Decide data uncertainties: zero, as in real data, study impact of inconsistent datasets, ...
 - Fit PDFs to the generated pseudodata using the same theory setup.
 - Check if the fitted PDFs reproduce the underlying truth:
 - are true values gaussianly distributed around the fit?
 - are uncertainties a faithful reproduction of input experimental ones?
 - are the results stable upon variations of the fitting methodology?

NNPDF3.0 is the first PDF determination based on Closure Tests



Three levels of Closure tests

Level 0

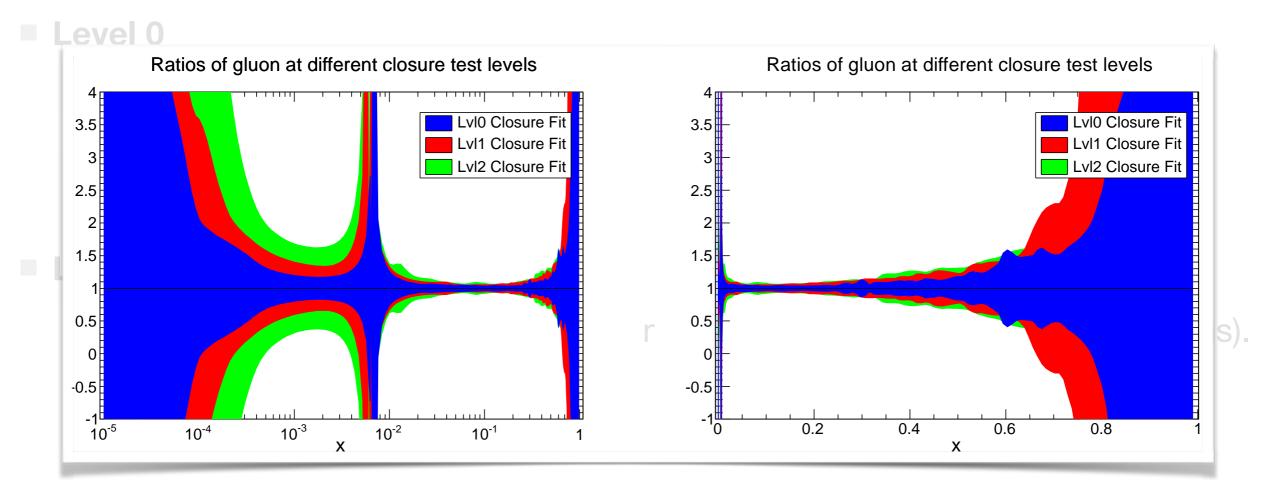
- Fake data are generated without uncertainty
- Test for efficiency and adequacy of fitting methodology
- Determine interpolation and extrapolation uncertainty

Level 1

- Fake data are generated with the same uncertainty of real data (correlations).
- No "data replicas", fit to the same data over and over.
- Determine functional uncertainty, due to infinity of equivalent minima
- Level 2
 - Replicas are fitted to fake data replicas
 - Determine data uncertainty



Three levels of Closure tests



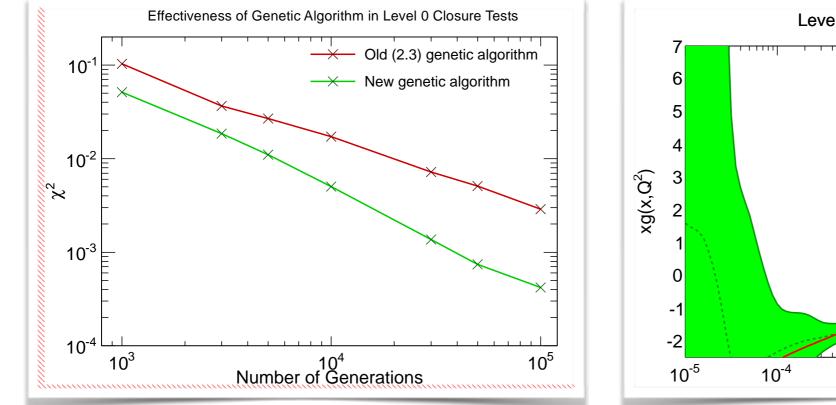
Level 2

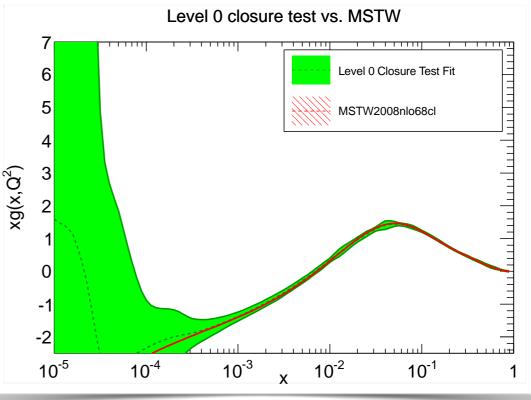
- The three sources of uncertainty are comparable in the data region
- Determine



Closure test results - Level 0

- Assume vanishing experimental uncertainty on generated data
- Perfect description of data (\chi_2 = 0) must be possible with adequate fitting methodology



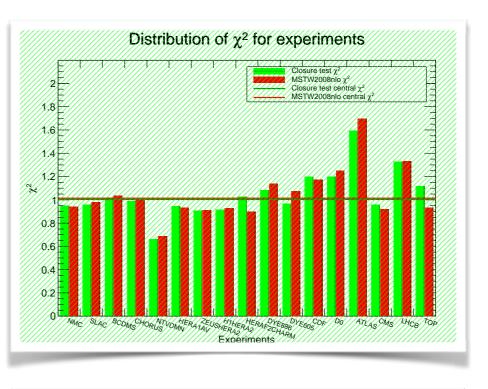




Central Values compare fitted vs. true χ^2 both for individual experiments and the entire dataset

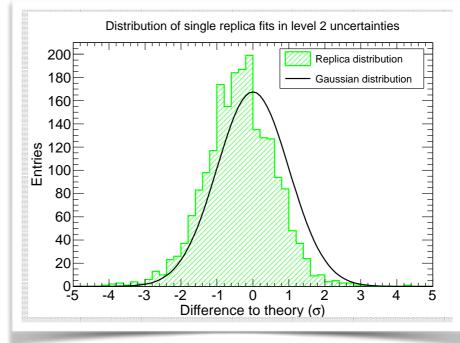


Closure test results - Level 2



Uncertainties

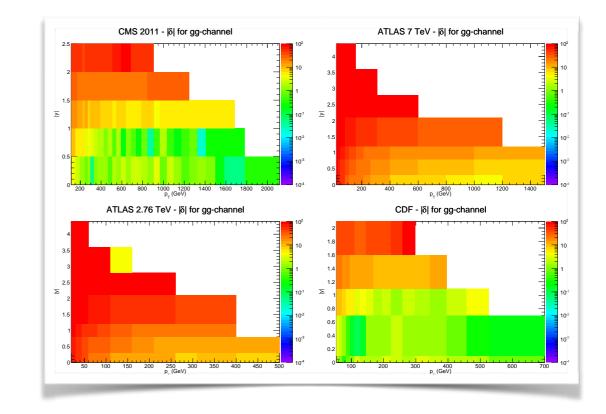
compare the distribution of deviations between fitted and true PDFs to theoretical expectation (gaussian)

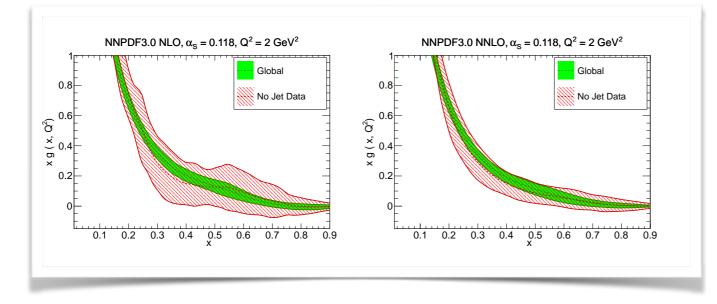




Theory improvements - Inclusive Jet data

- Systematic comparison of exact NNLO computation for the gg-channel (Gehrmann et al, 2014) and threshold approximation (De Florian et al, 2014) to determine which data points to include in the NNLO fit.
- Threshold resummation only accurate in the central rapidity and high p_T region (Carrazza & Pires, 2014)



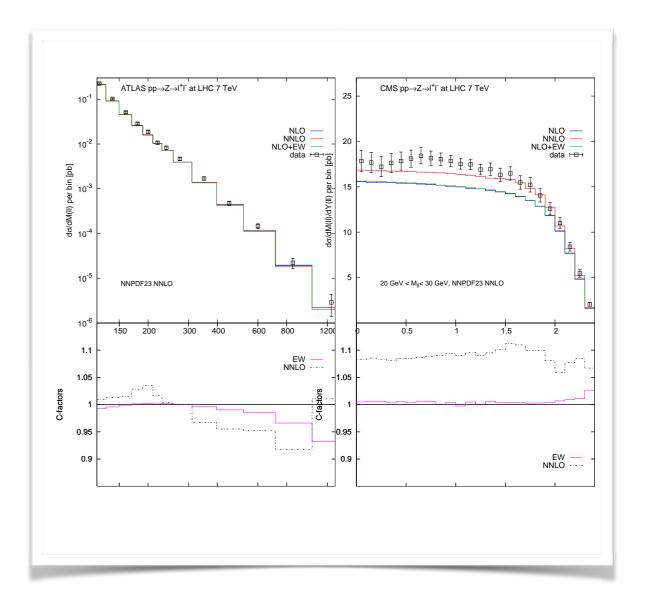


Impact of **jet data** on **large-x gluon** sizeable both at NLO at NNLO despite reduced number of data points



Theory improvements - Electroweak corrections

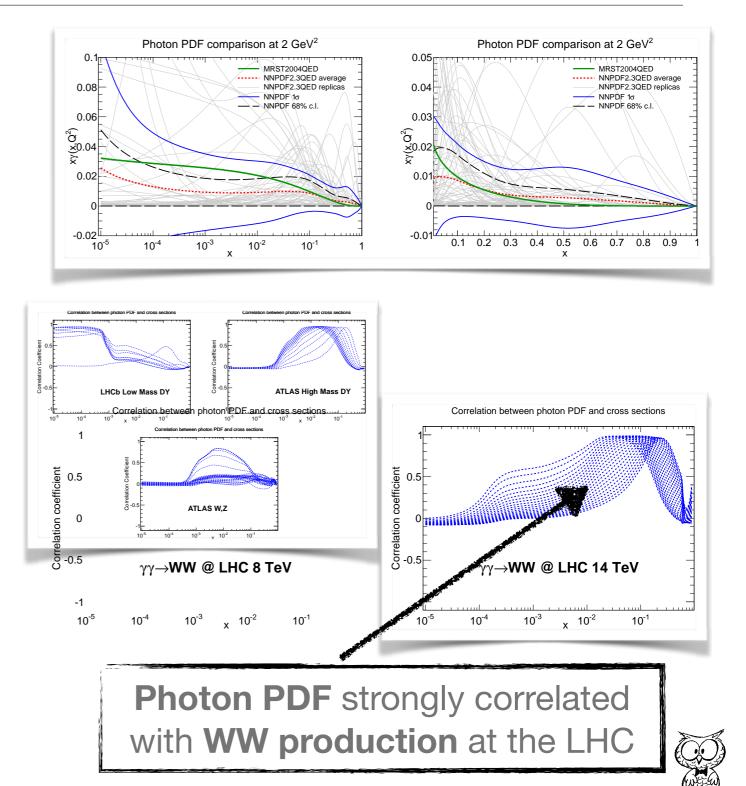
- Virtual pure EW corrections taken into account for all neutral Drell-Yan datasets (CMS double differential, ATLAS Z-peak and high-mass).
- Use data corrected for Final State Radiation
- QED corrections checked, highest invariant mass bins for ATLAS and CMS excluded from the fit because of large corrections due to photoninitiated processes
- Electroweak corrections still missing for DIS, W production, inclusive jet and top pair production





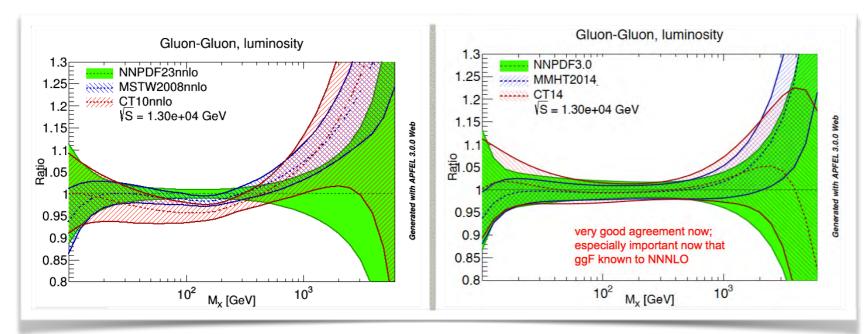
Parton Distributions with QED corrections (NNPDF2.3 QED)

- Precision LHC phenomenology, including EW effects, requires parton distributions with QED effects included in the evolution and a photon PDF
- NNPDF2.3 QED is the most recent PDF fit based on (N)NLO QCD + LO QED evolution and with a photon PDF determined from DIS and Drell-Yan (lowmass LHCb, W & Z peak and high-mass ATLAS) production R. D. Ball et al, [arXiv:1308.0598]
- LHC data are crucial for a reliable determination of the photon PDF



Improved agreement among global PDF sets

"Progress in convergence between the parton distribution functions will also be needed in order to reduce the theoretical uncertainties below the experimental measurement uncertainties"



J. Ellis, [arXiv:1504.03654]

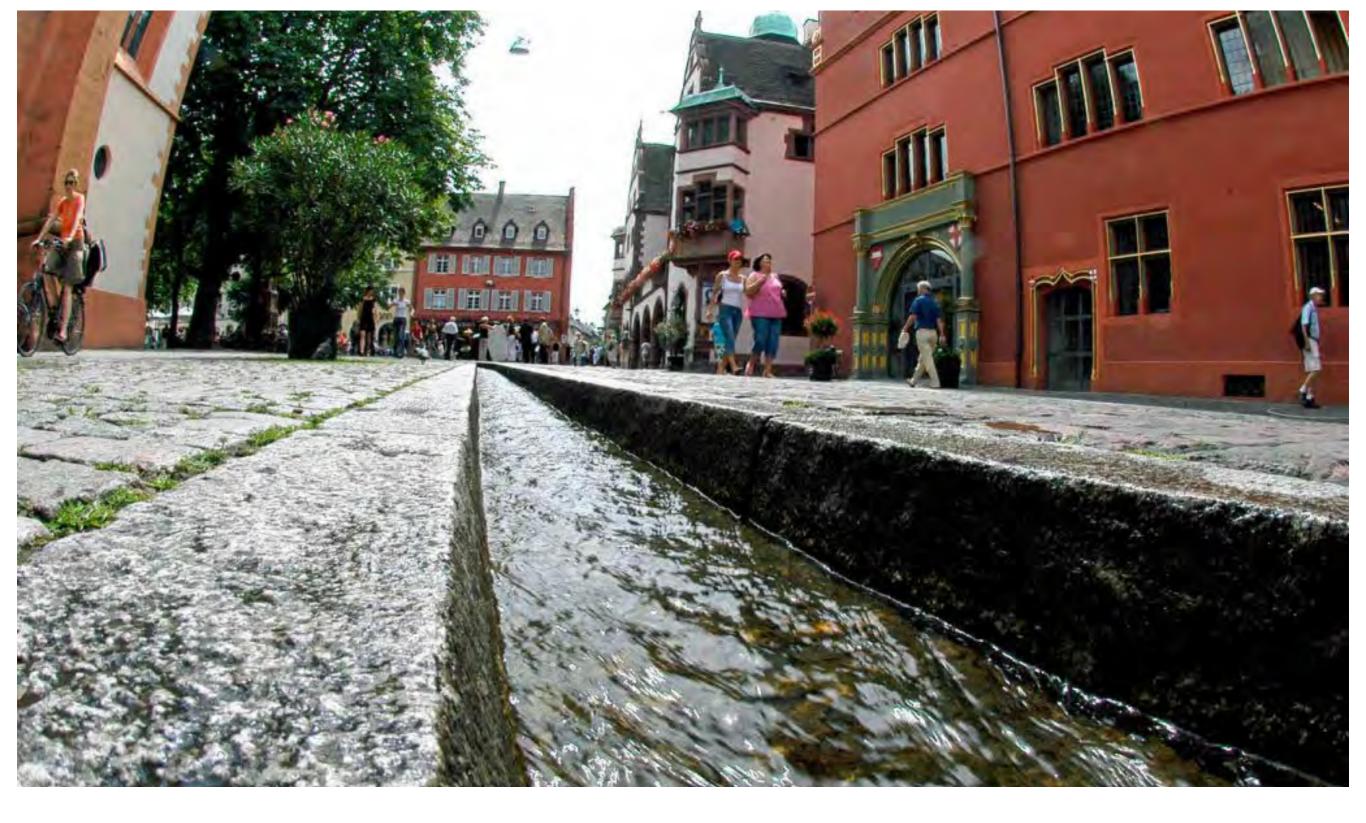
J. Houston, PDF4LHC 2015

8 TeV	18.66
Almost per	fect agreement
	newest releases
of global Pl	DF sets for ggH

NNPDF3.0 MMHT2014

18.77	CT14 18.65	MMHT2014	NNPDF3.0
12.9 7 8 TeV	18.66 pb	18.65 pb	18.77 pb
	-2.2%	-1.9%	-1.8%
	+2.0%	+1.4%	+1.8%
13 TeV	42.68 pb	42.70 pb	42.97 pb
	-2.4%	-1.8%	-1.9%
	+2.0%	+1.3%	+1.9%



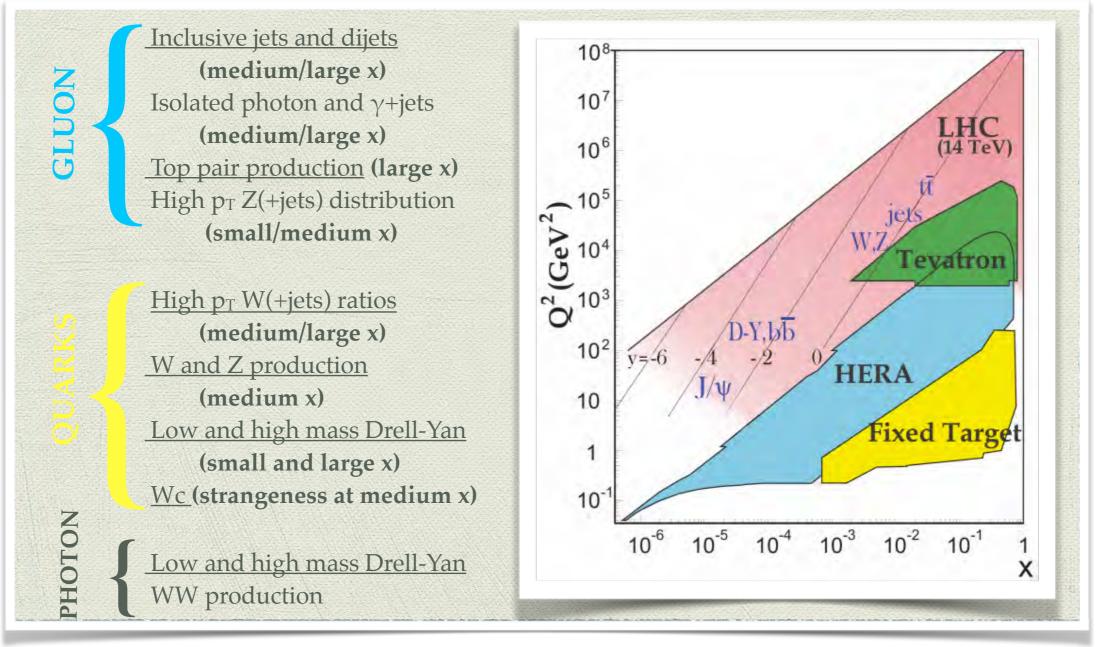


The road ahead...

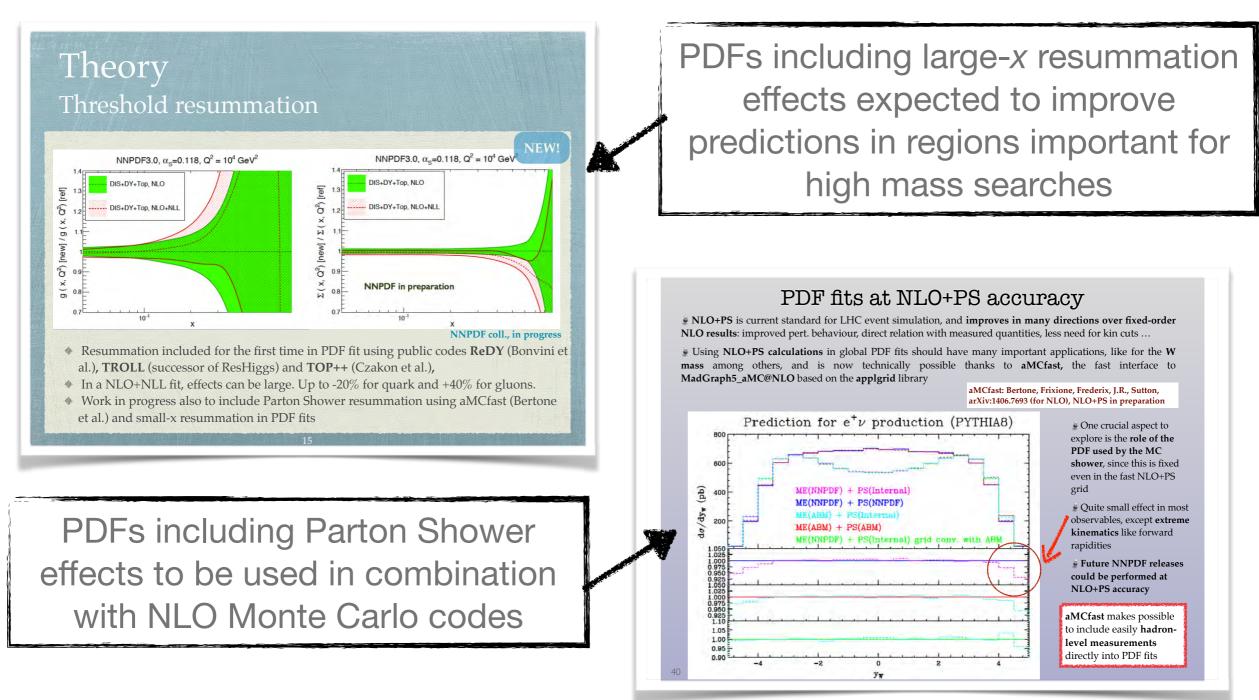
More data

• LHC data (both from Run I and Run II) to provide substantial constraints on PDFs in the (near) future





Theory refinements





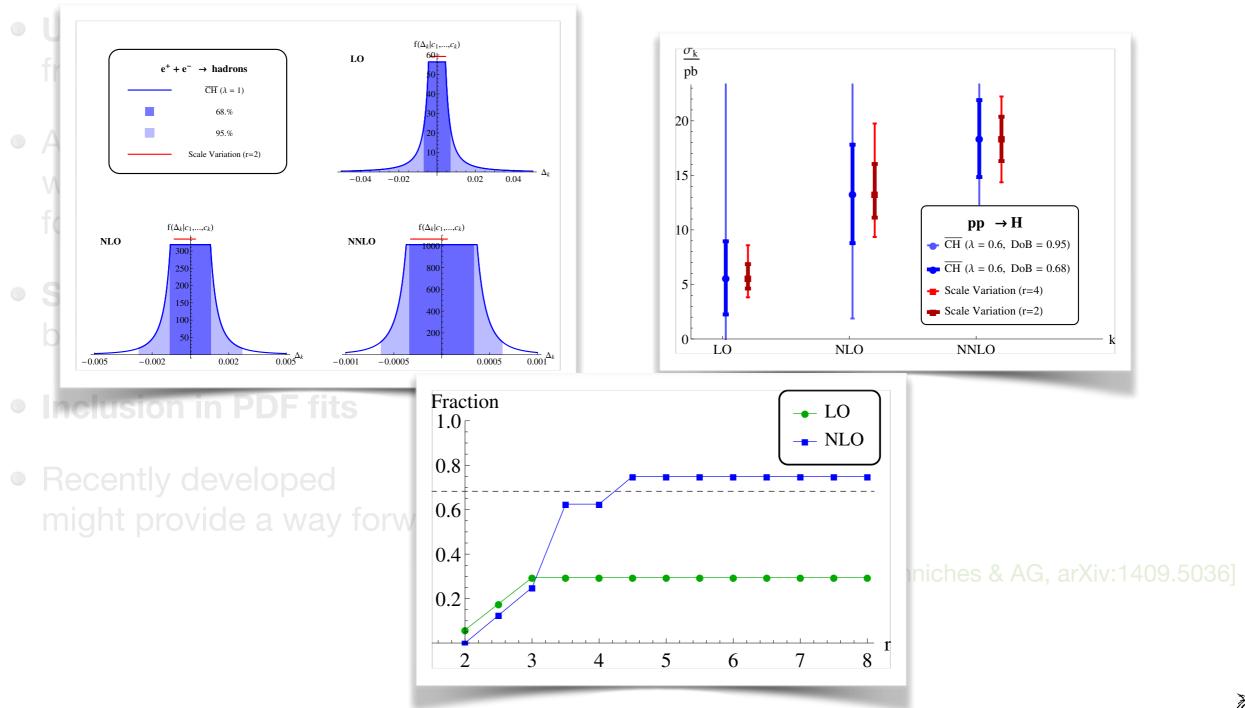
Theoretical uncertainties

- Uncertainties on PDFs only reflect experimental uncertainties propagated from the data included in the fit
- As data become more precise and constraining we should probably begin to worry about the uncertainties on the theoretical predictions we use in the fits, for example from missing higher orders
- Scale variations is the conventional way to estimate theoretical uncertainties, but the resulting uncertainty intervals have no statistical interpretation
- Inclusion in PDF fits is at best ambiguous
- Recently developed Bayesian framework to estimate theoretical uncertainties might provide a way forward ... work in progress!

[E. Bagnaschi, M. Cacciari, L. Jenniches & AG, arXiv:1409.5036]



Theoretical uncertainties





Conclusions & Outlook

Final thoughts

- An accurate knowledge of Parton Distribution Functions with a reliable estimate of their uncertainties is a crucial ingredient to exploit the potential of the LHC experiments
- Parton Distribution Function determinations evolved substantially in the last years, in every single aspect: data, theoretical input, fitting methodology ...
- ... agreement among global PDF fits (still based on different ingredients) for crucial observables like Higgs production in gluon-gluon fusion are a sign of progress ...
- The future promises
 - More high precision data from the LHC providing constraints on PDFs
 - More refined theoretical predictions (EW effects, resummations, ...)
 - ... we might finally start thinking about theoretical uncertainties

