

The global electroweak fit in a new era of precision

Roman Kogler (Universität Hamburg)

Seminar Graduiertenkolleg GRK 2044
Universität Freiburg
May 6, 2015



- ▶ Prerequisites and ingredients
- ▶ Results and status of the EW fit
- ▶ BSM constraints
- ▶ Future prospects

[illegible]

The Electroweak Sector of the SM

Electroweak interactions described by $SU(2) \times U(1)$

- ▶ 4 gauge bosons: 3 massive (Z, W^\pm), 1 massless (γ)
- ▶ 1 scalar (H)
 - extremely successful theory
 - taught in each particle physics course

The Electroweak Sector of the SM

Electroweak interactions described by $SU(2) \times U(1)$

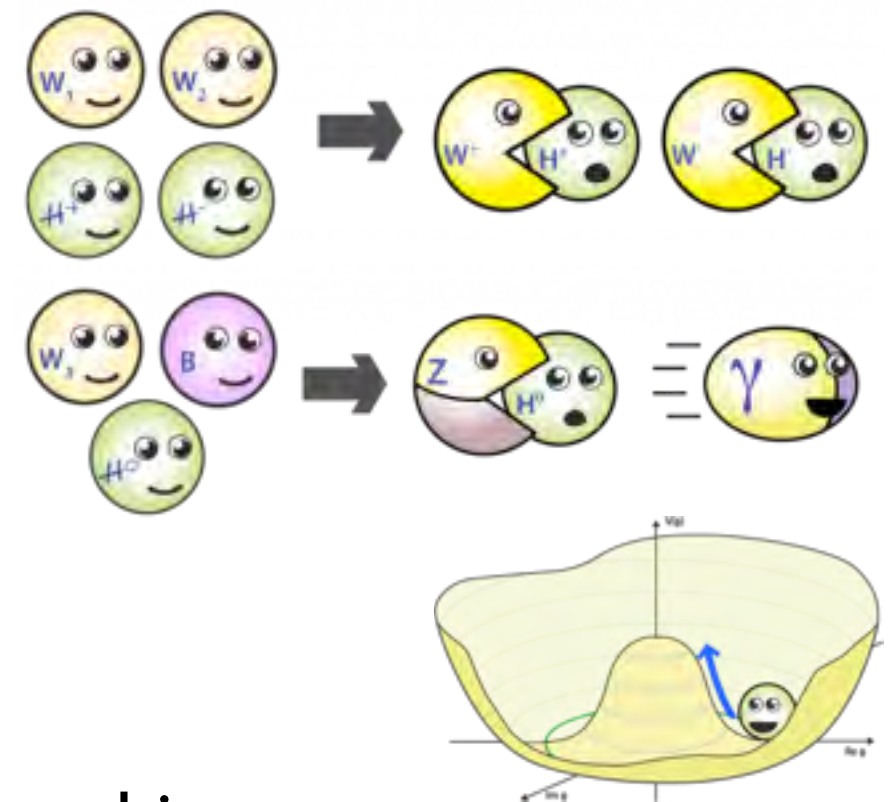
- ▶ 4 gauge bosons: 3 massive (Z, W^\pm), 1 massless (γ)
- ▶ 1 scalar (H)
 - extremely successful theory
 - taught in each particle physics course

Let's take one step back...

- ▶ it's a complicated, highly non-trivial theory
 - massive gauge bosons
 - parity (and CP) violation
 - Higgs field, results in a scalar particle

Why do we believe it?

- ▶ we physicists always had a hard time believing anything...
- ▶ we want to test the theory to ultimate precision!



[Philip Tanedo, quantumdiaries.org]

The Electroweak Sector of the SM

Electroweak sector given by 3 parameters

- ▶ g, g' : coupling constants of $SU(2)_L$ and $U(1)_Y$
- ▶ v : vacuum expectation value
- ▶ weak mixing angle : fixed by the massless photon

Use the three most precise parameters

- ▶ $\alpha : \Delta\alpha/\alpha = 3 \times 10^{-10}$
- ▶ $G_F : \Delta G_F/G_F = 5 \times 10^{-7}$
- ▶ $M_Z : \Delta M_Z/M_Z = 2 \times 10^{-5}$
- ▶ measure more than the minimal set of parameters to test the theory!

$$\begin{aligned} M_W &= \frac{v|g|}{2} \\ M_Z &= \frac{v\sqrt{g^2 + g'^2}}{2} \\ \cos \theta_W &= \frac{M_W}{M_Z} \end{aligned}$$

$$M_W^2 = \frac{M_Z^2}{2} \left(1 + \sqrt{1 - \frac{\sqrt{8}\pi\alpha}{G_F M_Z^2}} \right)$$

The Electroweak Sector of the SM

Electroweak sector given by 3 parameters

- ▶ g, g' : coupling constants of $SU(2)_L$ and $U(1)_Y$
- ▶ v : vacuum expectation value
- ▶ weak mixing angle : fixed by the massless photon

Use the three most precise parameters

- ▶ $\alpha : \Delta\alpha/\alpha = 3 \times 10^{-10}$
- ▶ $G_F : \Delta G_F/G_F = 5 \times 10^{-7}$
- ▶ $M_Z : \Delta M_Z/M_Z = 2 \times 10^{-5}$
- ▶ measure more than the minimal set of parameters to test the theory!

$$\begin{aligned} M_W &= \frac{v|g|}{2} \\ M_Z &= \frac{v\sqrt{g^2 + g'^2}}{2} \\ \cos \theta_W &= \frac{M_W}{M_Z} \end{aligned}$$

$$M_W^2 = \frac{M_Z^2}{2} \left(1 + \sqrt{1 - \frac{\sqrt{8}\pi\alpha}{G_F M_Z^2}} \right)$$

Calculate M_W and compare with experiment

- ▶ $M_W(\text{theo}) = 80.939 \pm 0.003 \text{ GeV}$
- ▶ $M_W(\text{exp}) = 80.385 \pm 0.015 \text{ GeV}$
- ▶ difference = $0.554 \text{ GeV} \sim 35\sigma$!! new physics?

Radiative Corrections

Modification of propagators and vertices

- ▶ Parametrisation of radiative corrections: electroweak form factors ρ , κ , Δr
- ▶ Effective couplings at the Z-pole:

$$g_{V,f} = \sqrt{\rho_Z^f} \left(I_3^f - 2Q^f \sin^2 \theta_{\text{eff}}^f \right)$$

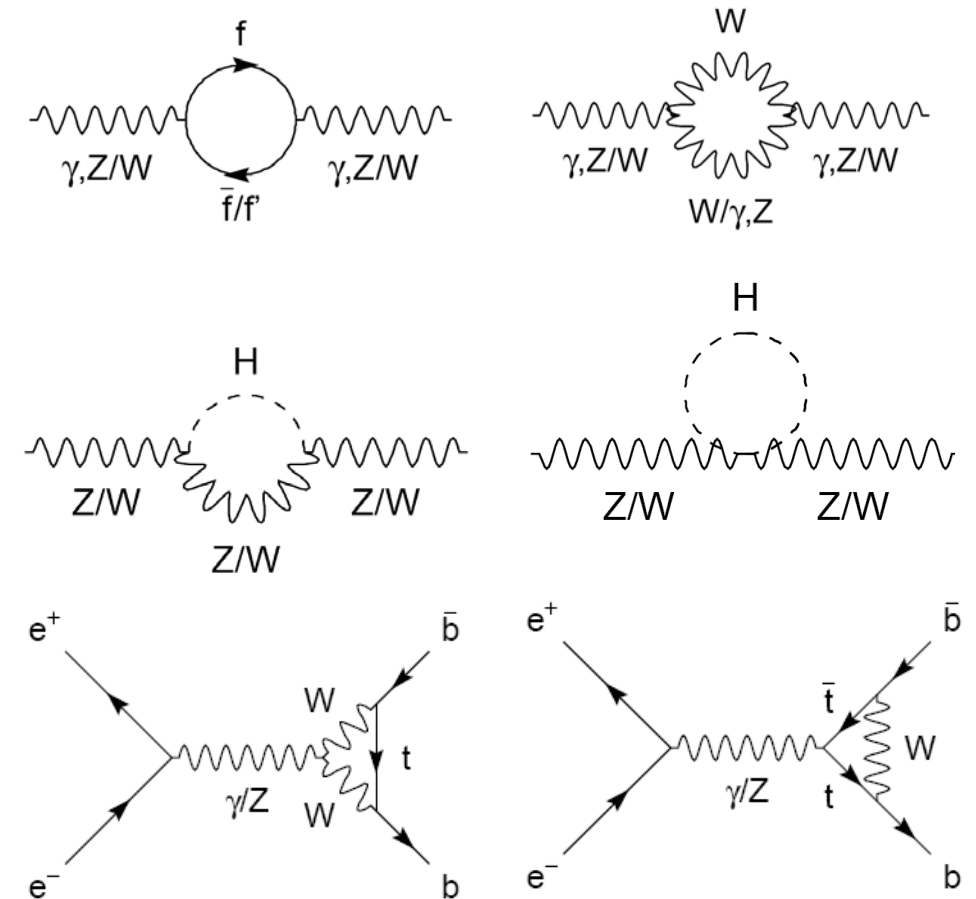
$$g_{A,f} = \sqrt{\rho_Z^f} I_3^f$$

$$\sin^2 \theta_{\text{eff}}^f = \kappa_Z^f \sin^2 \theta_W$$

- ▶ Mass of the W boson
$$M_W^2 = \frac{M_Z^2}{2} \left(1 + \sqrt{1 - \frac{\sqrt{8}\pi\alpha(1 + \Delta r)}{G_F M_Z^2}} \right)$$

- ▶ ρ , κ , Δr depend on all parameters of the theory (m_t , M_H , α_s ...)

$$\Delta r = -\frac{3\alpha c_W^2}{16\pi s_W^4} \frac{m_t^2}{M_W^2} + \frac{11\alpha}{48\pi s_W^2} \ln \frac{M_H^2}{M_W^2} + \dots$$



Free Parameters

EW sector

- ▶ G_F : $\Delta G_F/G_F = 5 \times 10^{-7}$
- ▶ M_Z : $\Delta M_Z/M_Z = 2 \times 10^{-5}$
- ▶ evolution of fine structure constant ($\Delta\alpha/\alpha = 3 \times 10^{-10}$) to scale s

$$\Delta\alpha(s) = \Delta\alpha_{\text{lep}}(s) + \Delta\alpha_{\text{had}}^{(5)}(s) + \Delta\alpha_{\text{top}}(s)$$

relative precision = $\boxed{1 \times 10^{-6}}$ $\boxed{2 \times 10^{-4}}$ $\boxed{1 \times 10^{-7}}$

Fermion masses

- ▶ m_c, m_b : precision of about 7% and 1%, sufficient (see later)
- ▶ m_t crucial parameter, experimental precision of 0.5% (more later)

Strong sector

- ▶ α_s : can be constrained using Z-pole measurements

Higgs sector

- ▶ M_H : precision of LHC measurements is 0.3%

Measure more than minimal set to constrain the theory

Measurements at e^+e^- Colliders

Z-pole measurements at LEP-I and SLC

- ▶ LEP : running near the Z-pole, four experiments, 4×10^6 Zs / experiment
- ▶ SLC : one experiment, 500.000 Zs, polarized beams

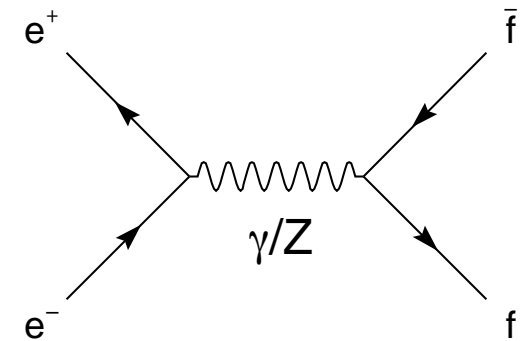
Precision measurements

- ▶ exactly known initial state
- ▶ precise beam energy, $\Delta E_{\text{beam}} = \pm 0.2$ MeV

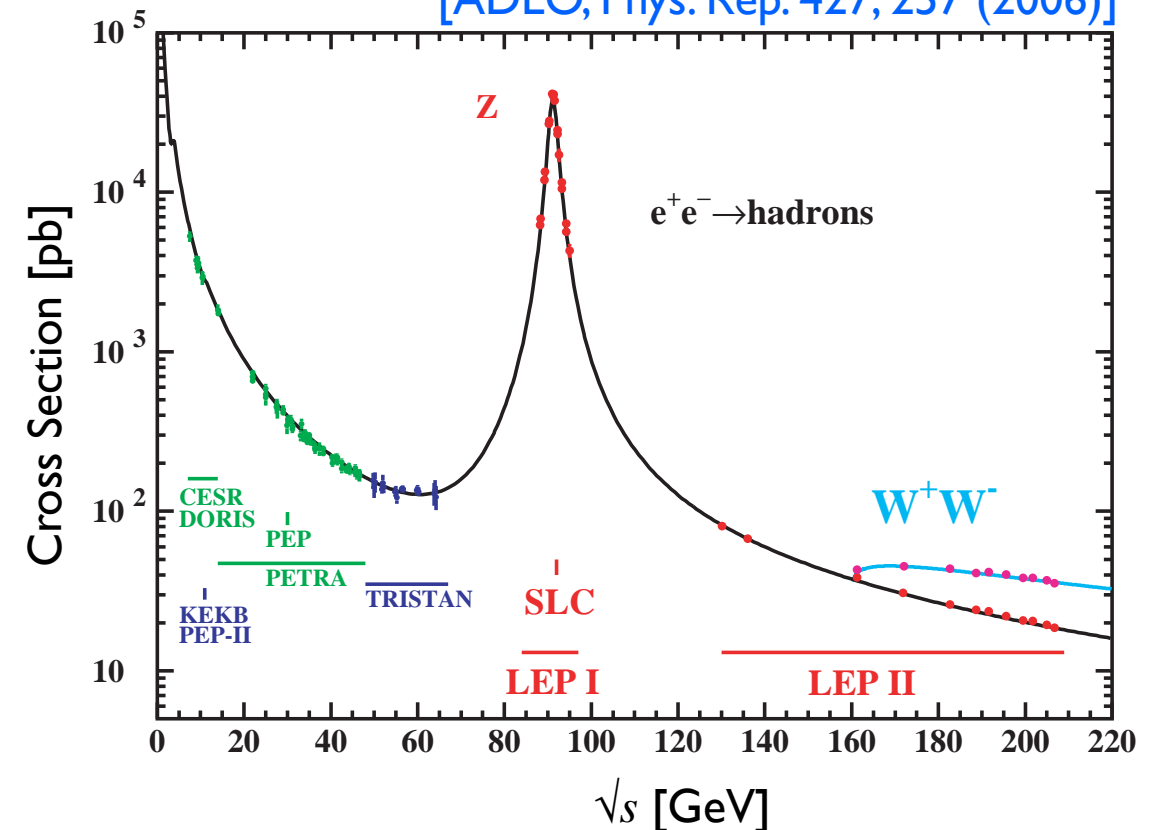
Cross section

$$\sigma_{f\bar{f}}^Z = \sigma_{f\bar{f}}^0 \frac{s\Gamma_Z^2}{(s - M_Z^2)^2 + s^2\Gamma_Z^2/M_Z^2} \frac{1}{R_{\text{QED}}}$$

$$\text{with } \sigma_{f\bar{f}}^0 = \frac{12\pi}{M_Z^2} \frac{\Gamma_{ee}\Gamma_{f\bar{f}}}{\Gamma_Z^2} \quad \text{and} \quad \Gamma_Z = \Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{\text{had}} + \Gamma_{\text{inv}}$$



[ADLO, Phys. Rep. 427, 257 (2006)]

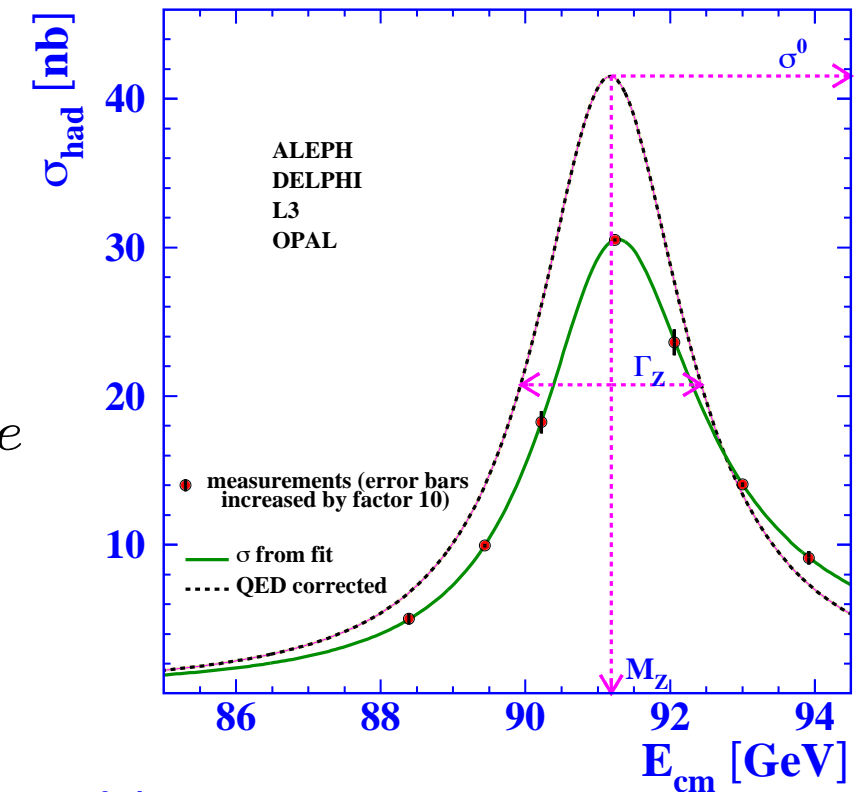


Observables

Minimal correlated set of parameters

- ▶ mass and total width of Z^0 M_Z, Γ_Z
- ▶ hadronic pole cross section σ_{had}^0
- ▶ leptonic decay ratios $R_\ell^0 = R_e^0 = \Gamma_{\text{had}}/\Gamma_{ee}$
- ▶ hadronic width ratios $R_{c,b}^0 = \Gamma_{c\bar{c},b\bar{b}}/\Gamma_{\text{had}}$

[ADLO, Phys. Rep. 427, 257 (2006)]



Asymmetries

- ▶ $A_f = \frac{g_{L,f}^2 - g_{R,f}^2}{g_{L,f}^2 + g_{R,f}^2} = 2 \frac{g_{V,f}/g_{A,f}}{1 + (g_{V,f}/g_{A,f})^2}$ directly related to $\sin^2 \theta_{\text{eff}}^{f\bar{f}}$

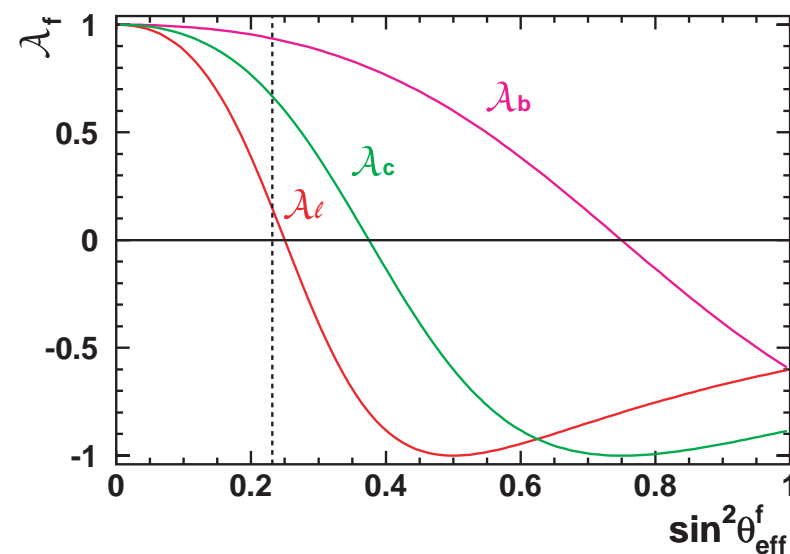
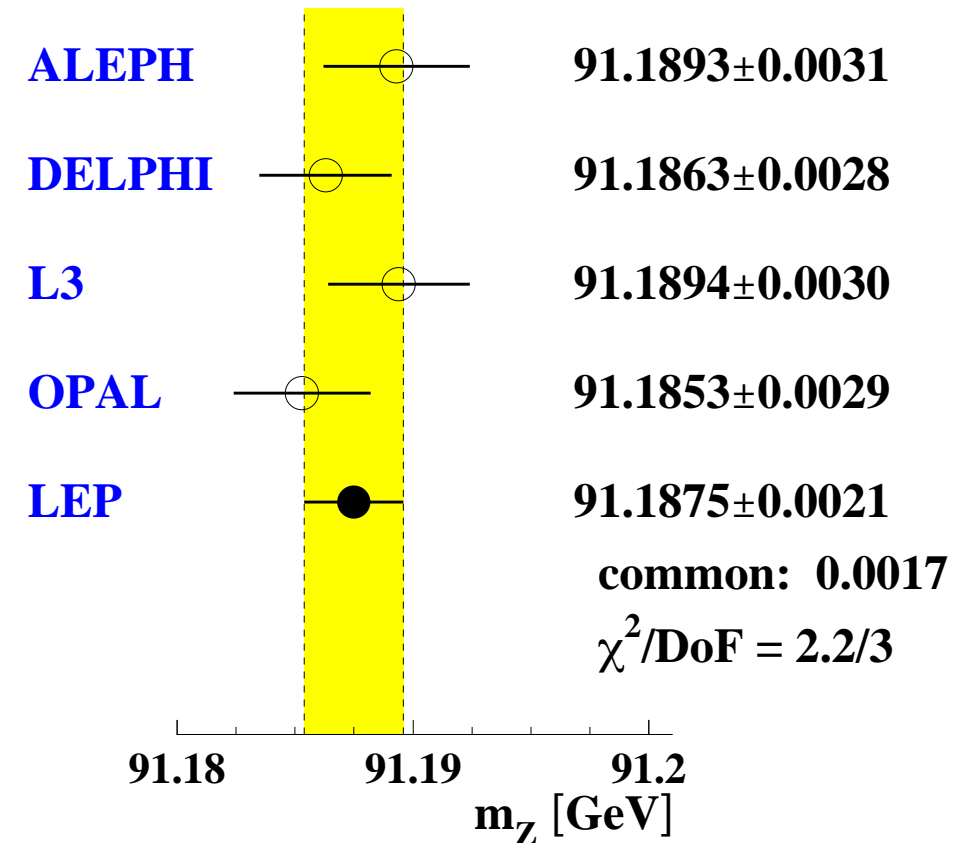
- ▶ forward/backward asymmetry $A_{FB}^f = \frac{N_F^f - N_B^f}{N_F^f + N_B^f}$, $A_{FB}^{0,f} = \frac{3}{4} A_e A_f$

- ▶ left/right asymmetry $A_{LR}^f = \frac{N_L^f - N_R^f}{N_L^f + N_R^f} \frac{1}{\langle |P|_e \rangle}$

Measurements at the Z-Pole

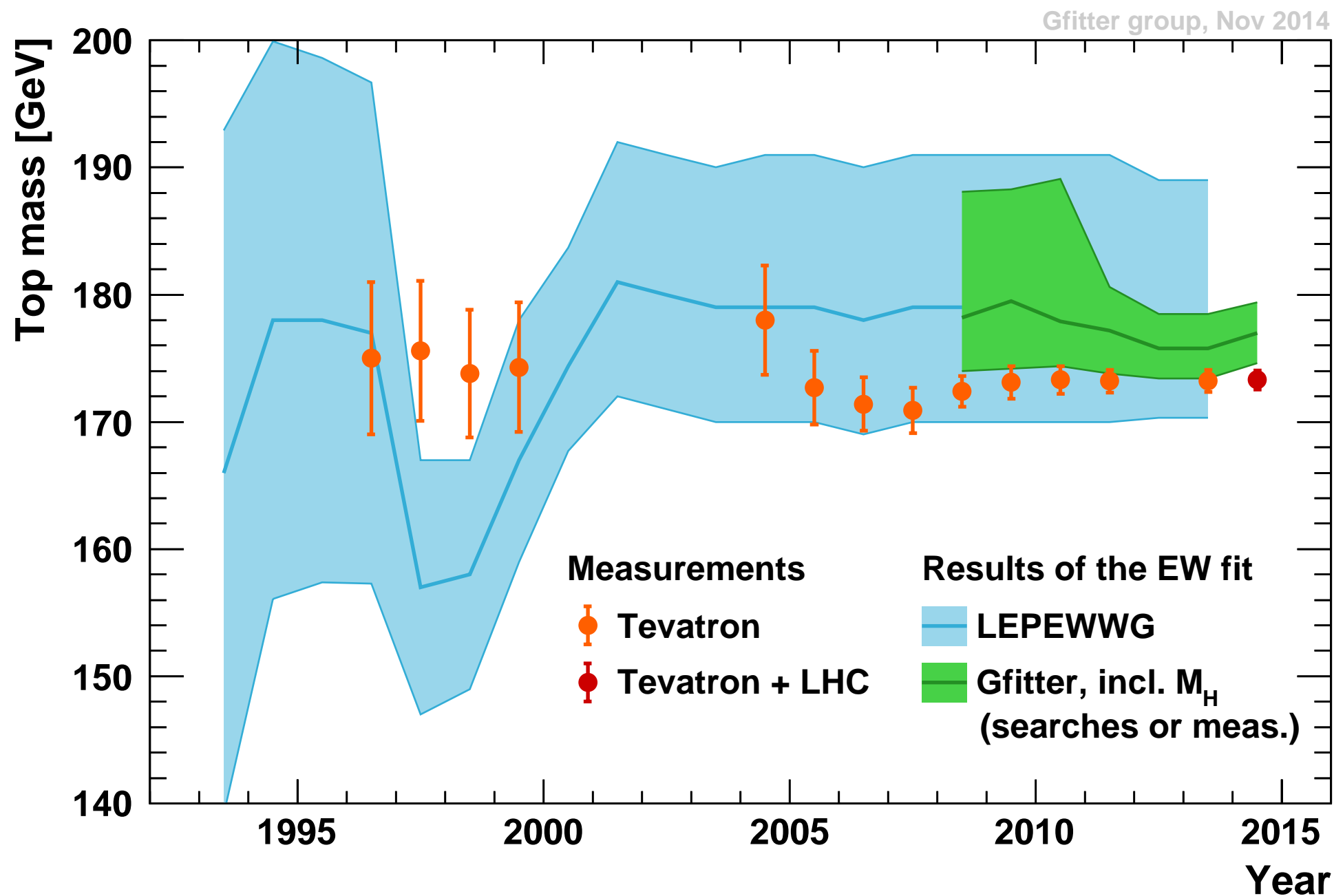
[ADLO, Phys. Rep. 427, 257 (2006)]

M_Z [GeV]	91.1875 ± 0.0021
Γ_Z [GeV]	2.4952 ± 0.0023
σ_{had}^0 [nb]	41.540 ± 0.037
R_ℓ^0	20.767 ± 0.025
$A_{\text{FB}}^{0,\ell}$	0.0171 ± 0.0010
$A_\ell^{(*)}$	0.1499 ± 0.0018
$\sin^2 \theta_{\text{eff}}^\ell(Q_{\text{FB}})$	0.2324 ± 0.0012
A_c	0.670 ± 0.027
A_b	0.923 ± 0.020
$A_{\text{FB}}^{0,c}$	0.0707 ± 0.0035
$A_{\text{FB}}^{0,b}$	0.0992 ± 0.0016
R_c^0	0.1721 ± 0.0030
R_b^0	0.21629 ± 0.00066



- ▶ precision of up to 0.002%!
- ▶ LEP/SLD measurements will stay the most precise for quite some time
- ▶ allow for precision tests of the SM and constrain new physics

Prediction of top quark mass



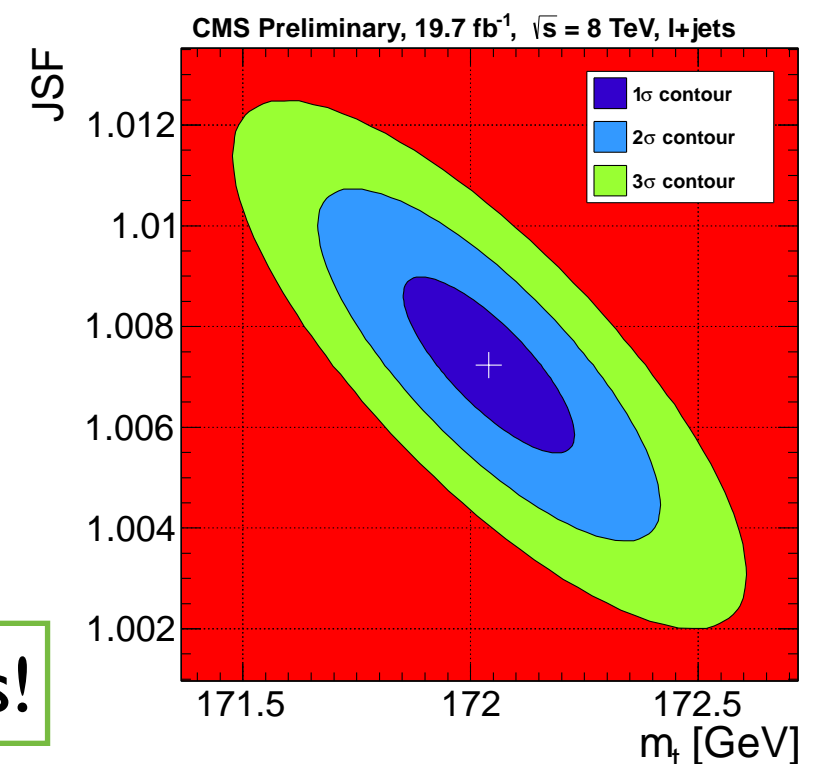
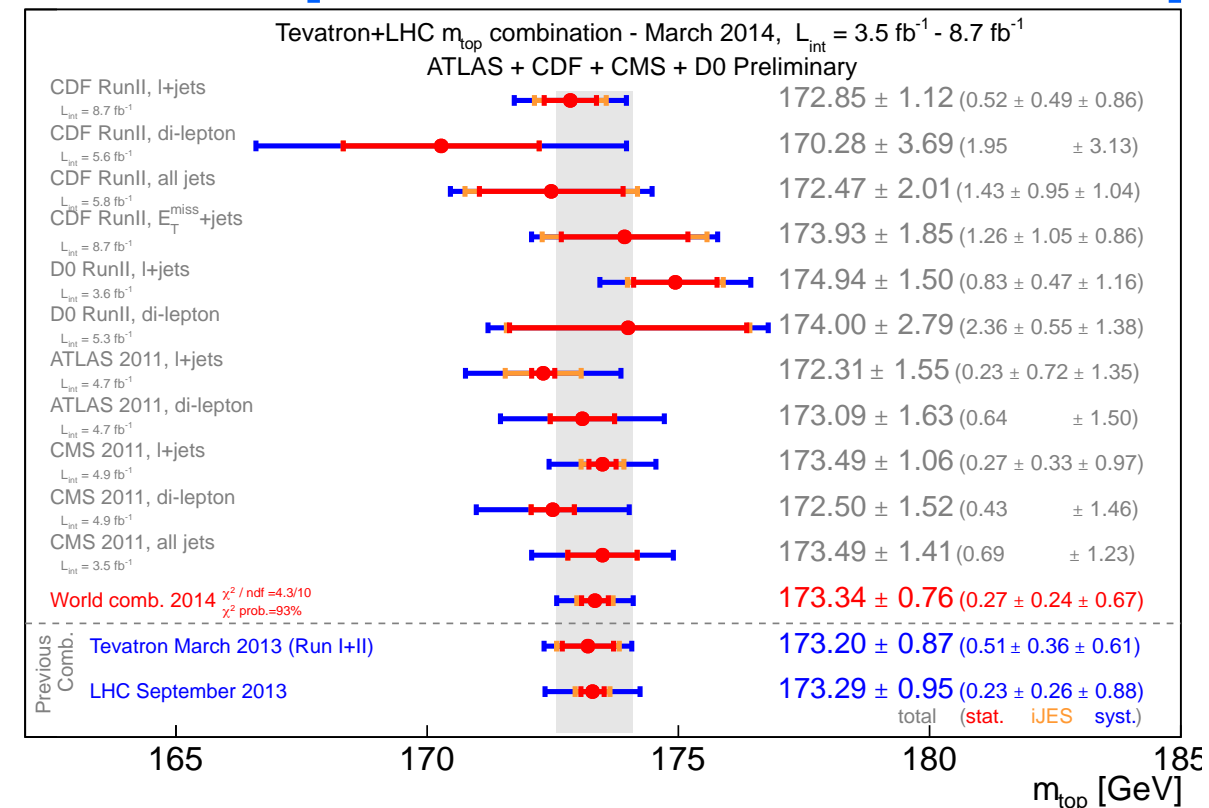
- ▶ m_t predictions from loop effects since 1990
- ▶ official LEPEWWG fit since 1993
- ▶ the fits have always been able to predict m_t correctly!

Measurements of m_t

- ▶ Tevatron pioneered measurements of a “kinematic” mass in t decays
- ▶ Tevatron
 - exceeding all expectations (expected precision: 2-3 GeV)
- ▶ LHC collaborations taking over
 - re-use of methods, high statistics
- ▶ world average: $m_t = 173.34 \pm 0.76$ GeV
 - single best measurement in WVA from CMS in $l+jets$ channel
 - recently updated [CMS-PAS-TOP-14-001]
 - $m_t = 172.04 \pm 0.19$ (stat.+JES) ± 0.75 (syst.) GeV
 - crucial: JER, pile-up, flavour dependence of JES
- ▶ Tevatron 2014: $\Delta m_t = 0.64$ GeV [D0, CDF, arXiv:1407.2682]

welcome to the community of precision measurements!

[CDF, D0, ATLAS, CMS: arXiv:1403.4427]



Interpretation of m_t measurements

What about accuracy?

► top mass definition

- **EFT, factorization:** hard function, universal jet-function, non-pert. soft function
[Moch et al, arXiv:1405.4781]

- MC mass is (may be) related to the low scale short-distance mass

in the jet function [Hoang, arXiv:1412.3649]

- but: **no quantitative statement available**

- relating m_t^{kin} to m_t^{pole} : $\Delta m_t \geq \Lambda_{\text{QCD}}$

► colour structure and hadronisation

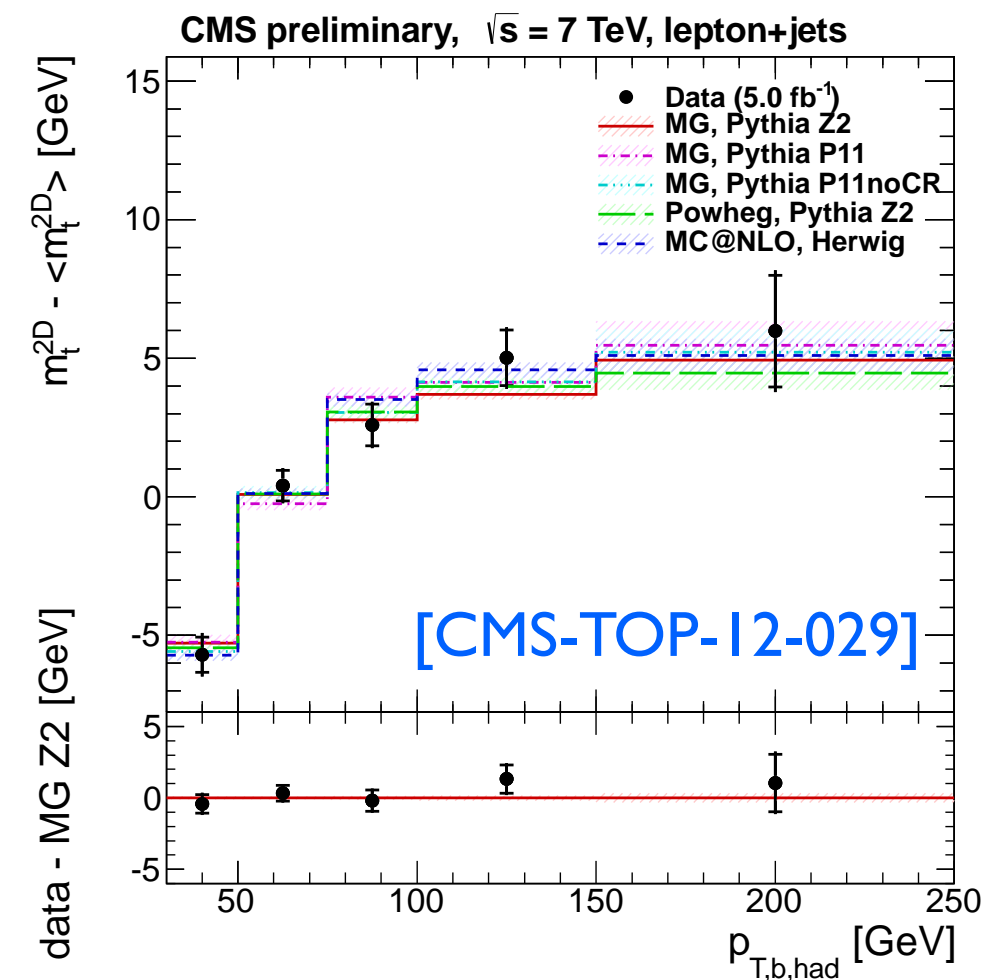
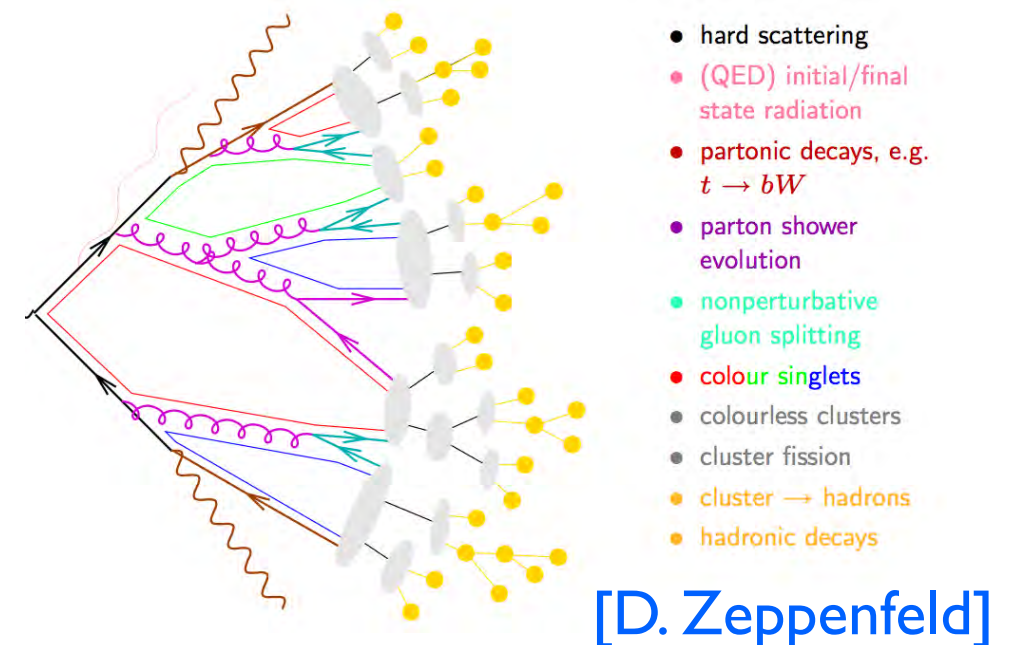
- partly included in experimental uncertainties
- study on kinematic dependencies of m_t

► calculating $m_t(m_t)$ from m_t^{pole}

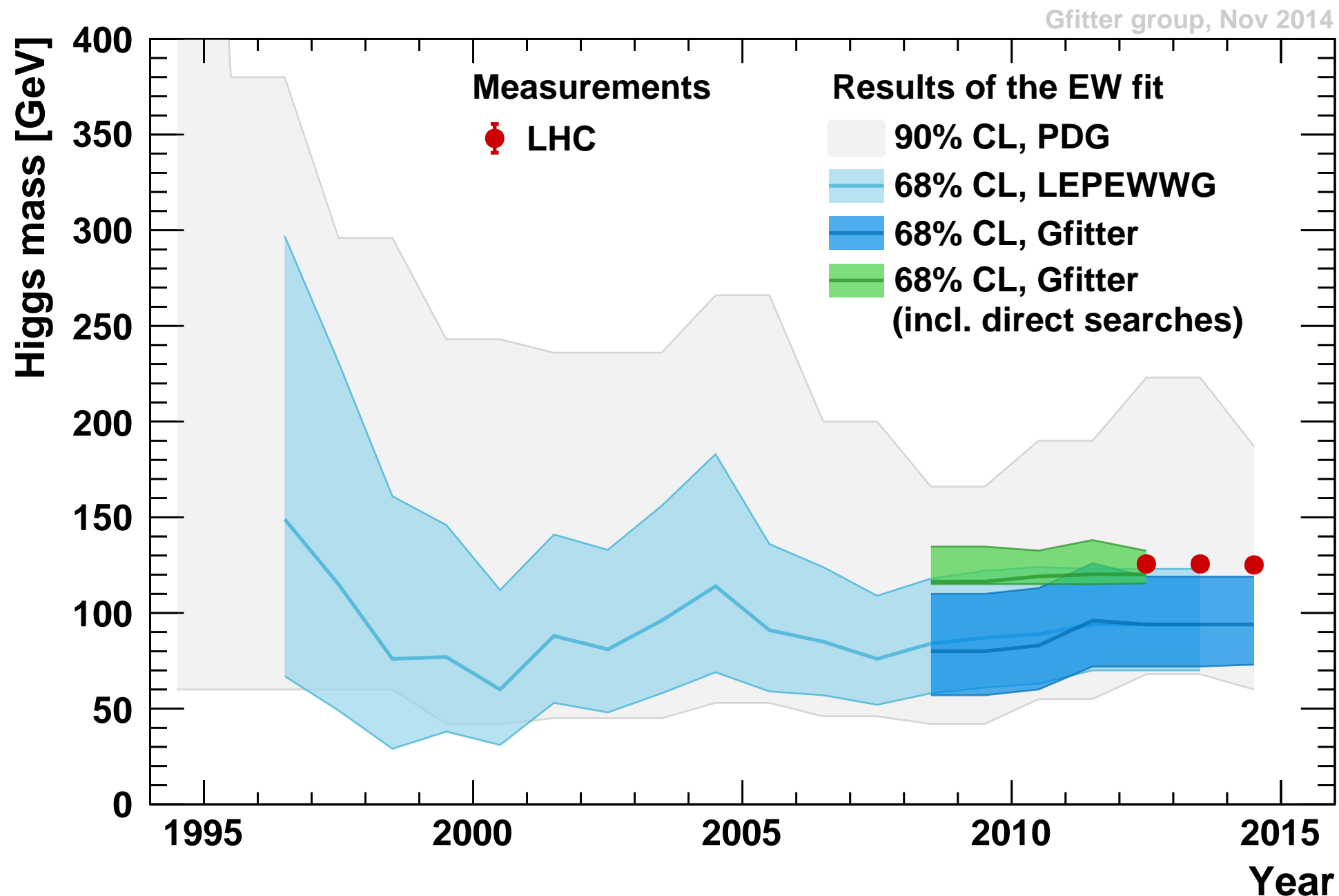
- QCD (three-loop): $\Delta m_t \approx 0.02 \text{ GeV}$

- EW (two-loop): $\Delta m_t \approx 0.1 \text{ GeV}$

[Kniehl et al., arXiv:1401.1844]



Prediction of Higgs mass

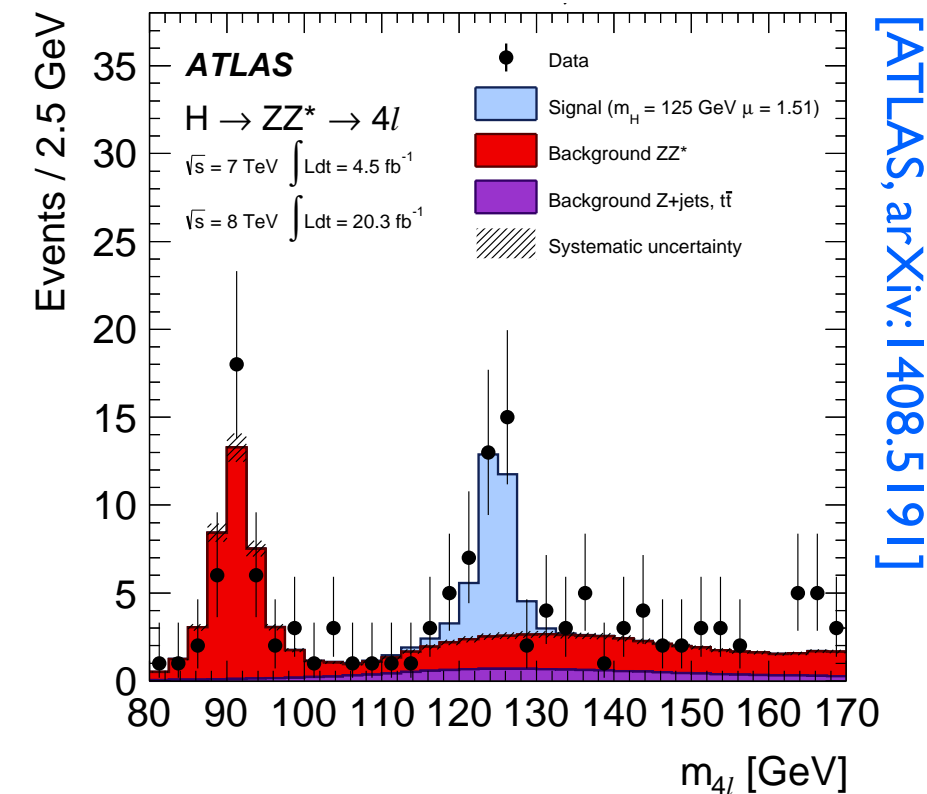


- ▶ M_H predictions from loop effects since the discovery of the top quark 1995
- ▶ weaker constraints than for m_t because of logarithmic dependence
- ▶ still, the fits have always predicted M_H correctly!

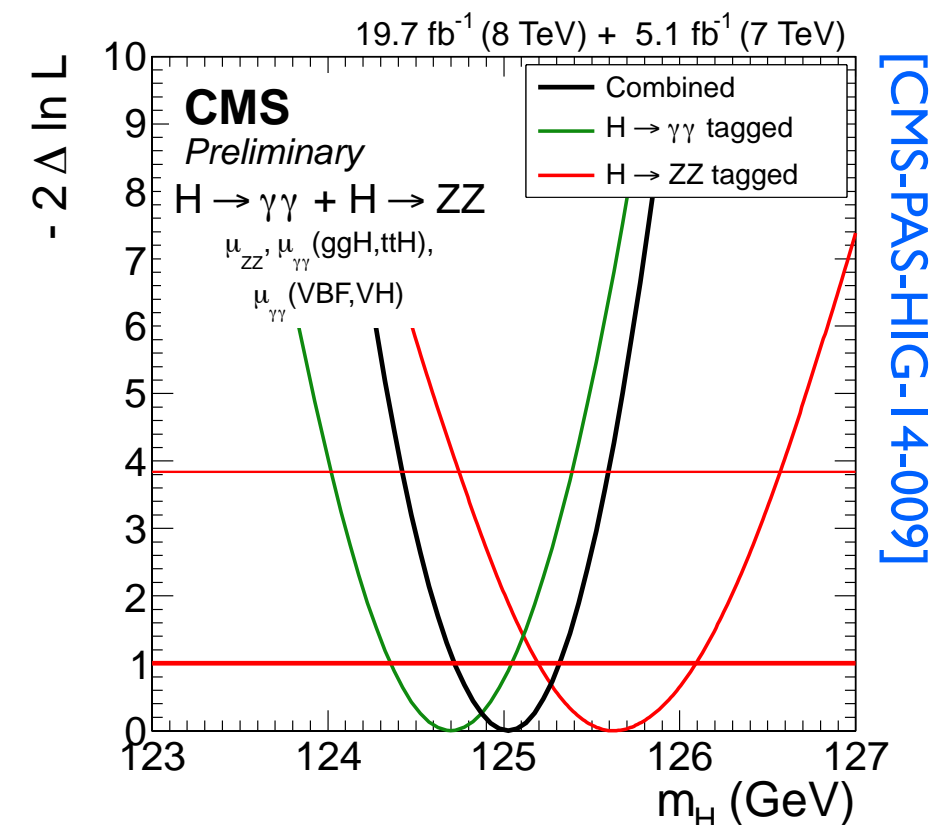
Measurements of M_H

Discovery of a Higgs boson

- ▶ cross section times branching ratios, spin, parity: compatible with SM Higgs boson
 - assume it's the SM Higgs boson
 - (or a BSM Higgs boson h in the decoupling region)
 - test the consistency of the SM including it
- ▶ best mass measurements: $H \rightarrow \gamma\gamma$, $H \rightarrow 4l$
 - ATLAS: 125.4 ± 0.4 GeV [ATLAS, 1406.3827]
 - CMS: 125.0 ± 0.3 GeV [CMS-PAS-HIG-14-009]
 - weighted average: 125.14 ± 0.24 GeV
 - change between fully uncorrelated and fully correlated systematic uncertainties is minor: $\delta M_H : 0.24 \rightarrow 0.32$ GeV
 - accuracy: 0.2% !
 - sufficient for electroweak fit (more later)



[ATLAS, arXiv:1408.5191]



[CMS-PAS-HIG-14-009]

Measurements of M_W

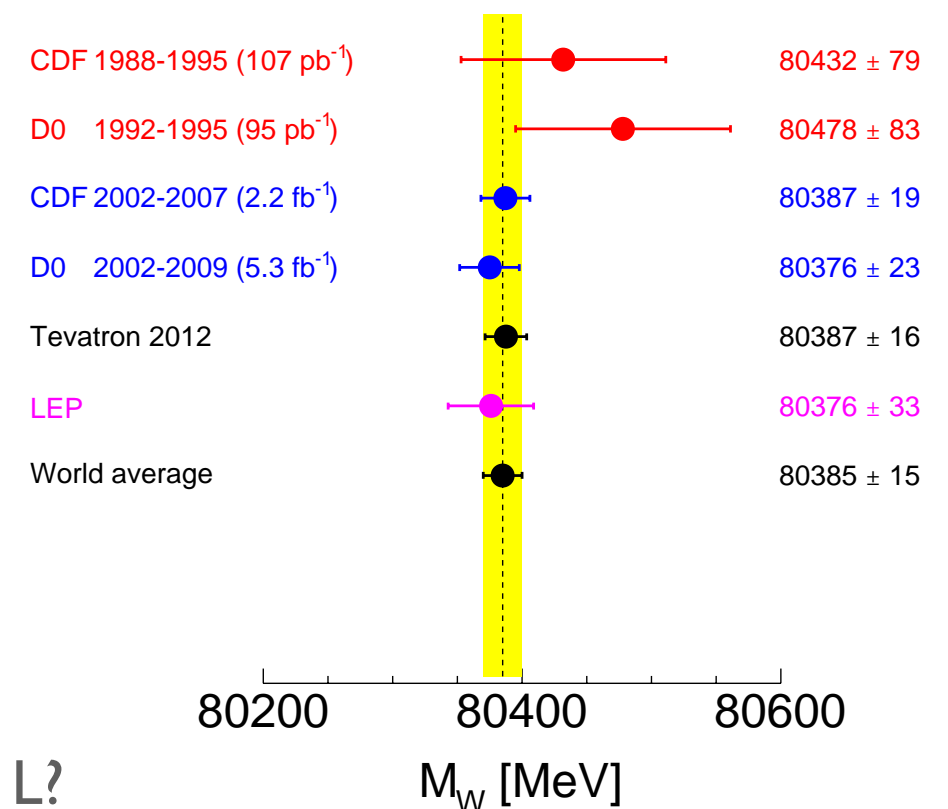
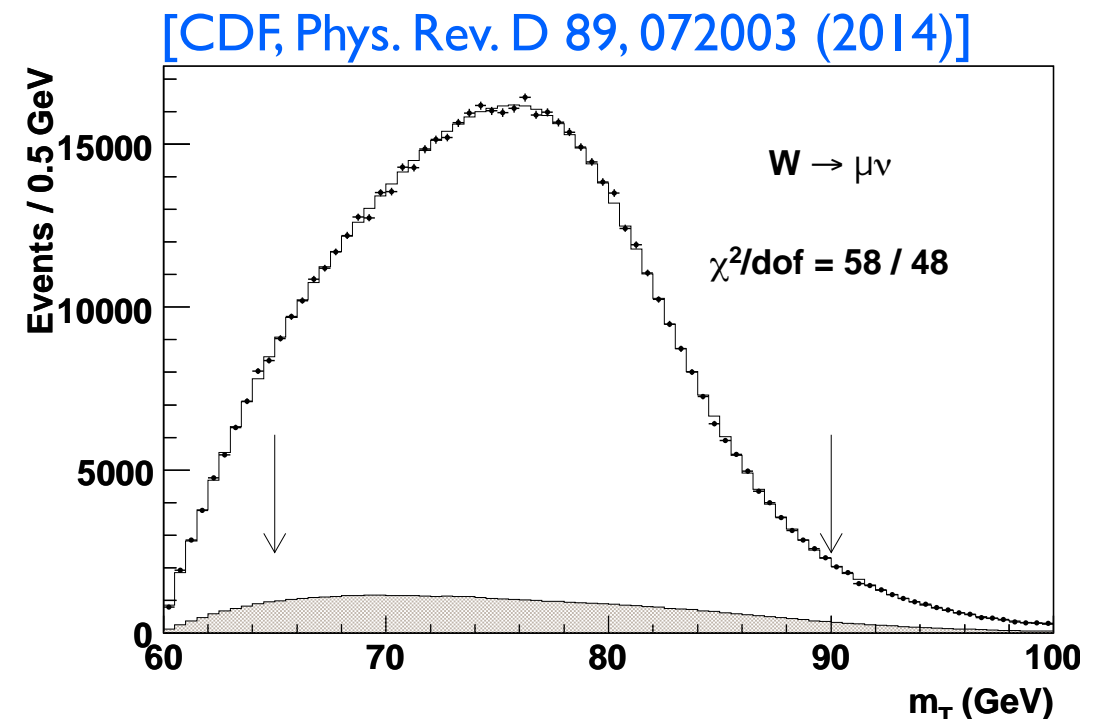
M_W : key parameter in the SM

$$\Delta r = -\frac{3\alpha c_W^2}{16\pi s_W^4} \frac{m_t^2}{M_W^2} + \frac{11\alpha}{48\pi s_W^2} \ln \frac{M_H^2}{M_W^2} + \dots$$

- ▶ final LEP-2 measurement (2013):
 - $\Delta M_W = 33 \text{ MeV}$ [ADLO, Phys. Rept. 532:119, 2013]
- ▶ Tevatron : most precise result so far
 - Jacobean peak in M_T and $p_{T,l}$ in $W \rightarrow l\nu$
 - $\Delta M = 16 \text{ MeV}$, accuracy: 0.02% !!
 - crucial: lepton energy and resolution, PDFs
- ▶ LHC : no result so far
 - (optimistic) scenarios: [arXiv:1310.6708]

ΔM_W [MeV]	LHC		
\sqrt{s} [TeV]	8	14	14
\mathcal{L} [fb $^{-1}$]	20	300	3000
Total	15	8	5

- very challenging
 - PDFs, momentum scale, hadronic recoil, pile-up at high L?



[CDF, D0, Phys. Rev. D 88, 052018 (2013)]

Experimental Input

Fit is overconstrained

- ▶ all free parameters measured
 - most input from e^+e^- colliders
 - but crucial input from hadron colliders:
 - m_t : 0.4%
 - M_W : 0.02%
 - M_H : 0.2%
 - remarkable experimental precision (<1%)
- ▶ require precision calculations!

M_H [GeV] ^(◦)	125.14 ± 0.24
M_W [GeV]	80.385 ± 0.015
Γ_W [GeV]	2.085 ± 0.042
M_Z [GeV]	91.1875 ± 0.0021
Γ_Z [GeV]	2.4952 ± 0.0023
σ_{had}^0 [nb]	41.540 ± 0.037
R_ℓ^0	20.767 ± 0.025
$A_{\text{FB}}^{0,\ell}$	0.0171 ± 0.0010
$A_\ell^{(*)}$	0.1499 ± 0.0018
$\sin^2\theta_{\text{eff}}^\ell(Q_{\text{FB}})$	0.2324 ± 0.0012
A_c	0.670 ± 0.027
A_b	0.923 ± 0.020
$A_{\text{FB}}^{0,c}$	0.0707 ± 0.0035
$A_{\text{FB}}^{0,b}$	0.0992 ± 0.0016
R_c^0	0.1721 ± 0.0030
R_b^0	0.21629 ± 0.00066
\overline{m}_c [GeV]	$1.27^{+0.07}_{-0.11}$
\overline{m}_b [GeV]	$4.20^{+0.17}_{-0.07}$
m_t [GeV]	173.34 ± 0.76
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$	2757 ± 10

LHC

Tev.

LEP

SLD

SLD

LEP

Tev.+LHC

Calculations

All observables calculated at 2-loop level

- ▶ M_W : full EW one- and two-loop calculation of fermionic and bosonic contributions

[M Awramik et al., PRD 69, 053006 (2004), PRL 89, 241801 (2002)]

+ 4-loop QCD correction [Chetyrkin et al., PRL 97, 102003 (2006)]

- ▶ $\sin^2\theta_{\text{eff}}^l$: same order as M_W , calculations for leptons and all quark flavours

[M Awramik et al, PRL 93, 201805 (2004), JHEP 11, 048 (2006), Nucl. Phys. B813, 174 (2009)]

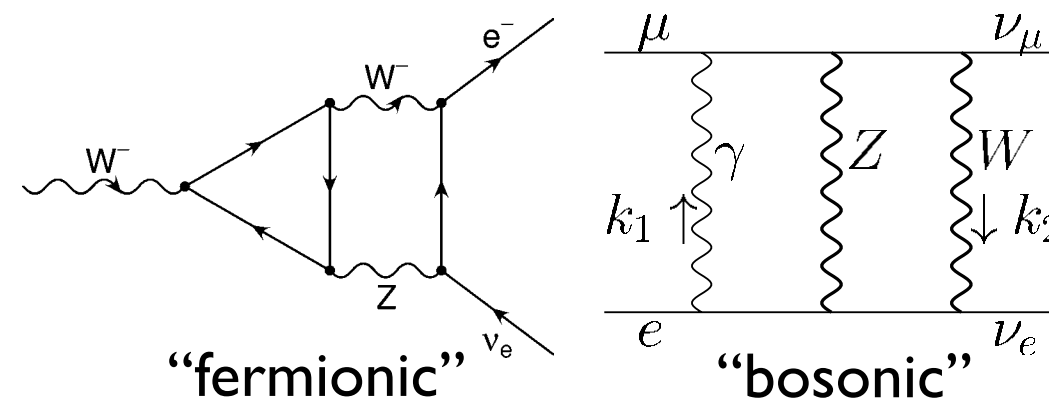
- ▶ partial widths Γ_f : fermionic corrections known to two-loop level for all flavours (includes predictions for σ_{had}^0) [A. Freitas, JHEP04, 070 (2014)]

- ▶ Radiator functions: QCD corrections at N³LO [Baikov et al., PRL 108, 222003 (2012)]

- ▶ Γ_W : only one-loop EW corrections available, negligible impact on fit

[Cho et al, JHEP 1111, 068 (2011)]

- ▶ all calculations include one- and two-loop QCD corrections and leading terms of higher order corrections



All EWPOs calculated at two-loop level or better

Theoretical Uncertainties

Estimation

- ▶ assume that perturbative expansion follows a **geometric series** ($a_n = a r^n$) :

for example: $\mathcal{O}(\alpha^2 \alpha_s) = \frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s)$

- ▶ other methods (e.g. scale variation) not always feasible

- but give **similar results**

- ▶ theoretical uncertainties smaller by a factor of 3-6 than measurements

- for the first time, **reasonable estimate for all observables**

- ▶ important missing higher order terms:

- $\mathcal{O}(\alpha^3)$, $\mathcal{O}(\alpha^2 \alpha_s)$, $\mathcal{O}(\alpha \alpha_s^2)$, $\mathcal{O}(\alpha^2 \alpha_{\text{bos}})$ (in some cases), $\mathcal{O}(\alpha_s^5)$ (rad. functions)

Observable	Exp. error	important ↓ Theo. error
M_W	15 MeV	4 MeV
$\sin^2 \theta_{\text{eff}}^l$	$1.6 \cdot 10^{-4}$	$0.5 \cdot 10^{-4}$
Γ_Z	2.3 MeV	0.5 MeV
$\sigma_{\text{had}}^0 = \sigma[e^+ e^- \rightarrow Z \rightarrow \text{had.}]$	37 pb	6 pb
$R_b^0 = \Gamma[Z \rightarrow b\bar{b}]/\Gamma[Z \rightarrow \text{had.}]$	$6.6 \cdot 10^{-4}$	$1.5 \cdot 10^{-4}$
m_t	0.76 GeV	0.5 GeV

↑
new in fit

Fit method

Free parameters

- ▶ $M_Z, \Delta\alpha_{\text{had}}, M_H, m_c, m_b, m_t, \alpha_s$
 - G_F is fixed to world average (PDG)
 - α_s is unconstrained → independent measurement

Treatment of theory uncertainties

- ▶ included as additional free parameters (10 parameters)
- ▶ different ways on how to treat their effect on the likelihood
 - **Rfit** : flat likelihood within uncertainties (box potential), corresponds to linear addition of uncertainties
 - **Gaussian likelihood** : corresponds to quadratic sum of uncertainties

Minimization

- ▶ pre-fitter : genetic algorithm (useful for many parameter fits)
- ▶ **Minuit** (standard, others are used as well)
- ▶ test of results using MC toy data

The global electroweak fit

disclaimer:

- ▶ there are several groups who routinely perform the electroweak fit
- ▶ there are small differences in the methodology, the results agree very well
- ▶ I will focus on results from the Gfitter group (www.cern.ch/gfitter)

Parameter	Input value	Free in fit	Fit Result	w/o exp. input in line	w/o exp. input in line, no theo. unc
M_H [GeV] ^(◦)	125.14 ± 0.24	yes	125.14 ± 0.24	93^{+25}_{-21}	93^{+24}_{-20}
M_W [GeV]	80.385 ± 0.015	—	80.364 ± 0.007	80.358 ± 0.008	80.358 ± 0.006
Γ_W [GeV]	2.085 ± 0.042	—	2.091 ± 0.001	2.091 ± 0.001	2.091 ± 0.001
M_Z [GeV]	91.1875 ± 0.0021	yes	91.1880 ± 0.0021	91.200 ± 0.011	91.2000 ± 0.010
Γ_Z [GeV]	2.4952 ± 0.0023	—	2.4950 ± 0.0014	2.4946 ± 0.0016	2.4945 ± 0.0016
σ_{had}^0 [nb]	41.540 ± 0.037	—	41.484 ± 0.015	41.475 ± 0.016	41.474 ± 0.015
R_ℓ^0	20.767 ± 0.025	—	20.743 ± 0.017	20.722 ± 0.026	20.721 ± 0.026
$A_{\text{FB}}^{0,\ell}$	0.0171 ± 0.0010	—	0.01626 ± 0.0001	0.01625 ± 0.0001	0.01625 ± 0.0001
$A_\ell^{(*)}$	0.1499 ± 0.0018	—	0.1472 ± 0.0005	0.1472 ± 0.0005	0.1472 ± 0.0004
$\sin^2\theta_{\text{eff}}^\ell(Q_{\text{FB}})$	0.2324 ± 0.0012	—	0.23150 ± 0.00006	0.23149 ± 0.00007	0.23150 ± 0.00005
A_c	0.670 ± 0.027	—	0.6680 ± 0.00022	0.6680 ± 0.00022	0.6680 ± 0.00016
A_b	0.923 ± 0.020	—	0.93463 ± 0.00004	0.93463 ± 0.00004	0.93463 ± 0.00003
$A_{\text{FB}}^{0,c}$	0.0707 ± 0.0035	—	0.0738 ± 0.0003	0.0738 ± 0.0003	0.0738 ± 0.0002
$A_{\text{FB}}^{0,b}$	0.0992 ± 0.0016	—	0.1032 ± 0.0004	0.1034 ± 0.0004	0.1033 ± 0.0003
R_c^0	0.1721 ± 0.0030	—	$0.17226^{+0.00009}_{-0.00008}$	0.17226 ± 0.00008	0.17226 ± 0.00006
R_b^0	0.21629 ± 0.00066	—	0.21578 ± 0.00011	0.21577 ± 0.00011	0.21577 ± 0.00004
\overline{m}_c [GeV]	$1.27^{+0.07}_{-0.11}$	yes	$1.27^{+0.07}_{-0.11}$	—	—
\overline{m}_b [GeV]	$4.20^{+0.17}_{-0.07}$	yes	$4.20^{+0.17}_{-0.07}$	—	—
m_t [GeV]	173.34 ± 0.76	yes	173.81 ± 0.85	$177.0^{+2.3(\nabla)}_{-2.4}$	$177.0 \pm 2.3(\nabla)$
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)^{(\dagger\Delta)}$	2757 ± 10	yes	2756 ± 10	2723 ± 44	2722 ± 42
$\alpha_s(M_Z^2)$	—	yes	0.1196 ± 0.0030	0.1196 ± 0.0030	0.1196 ± 0.0028

[Gfitter group, EPJ C 74, 3046 (2014)]

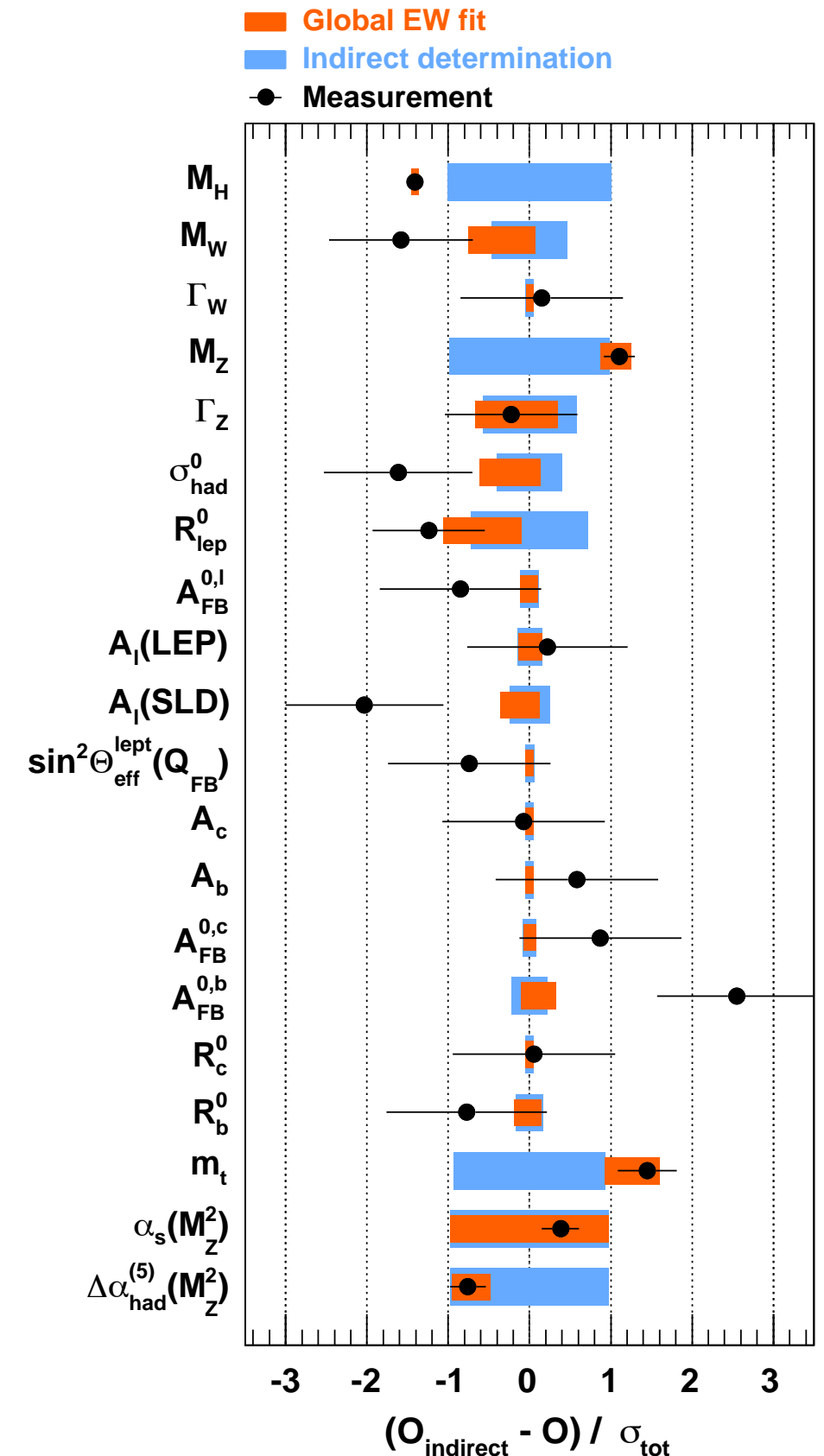
SM Fit Results

black: direct measurement (data)

orange: full fit

light-blue: fit excluding input from row

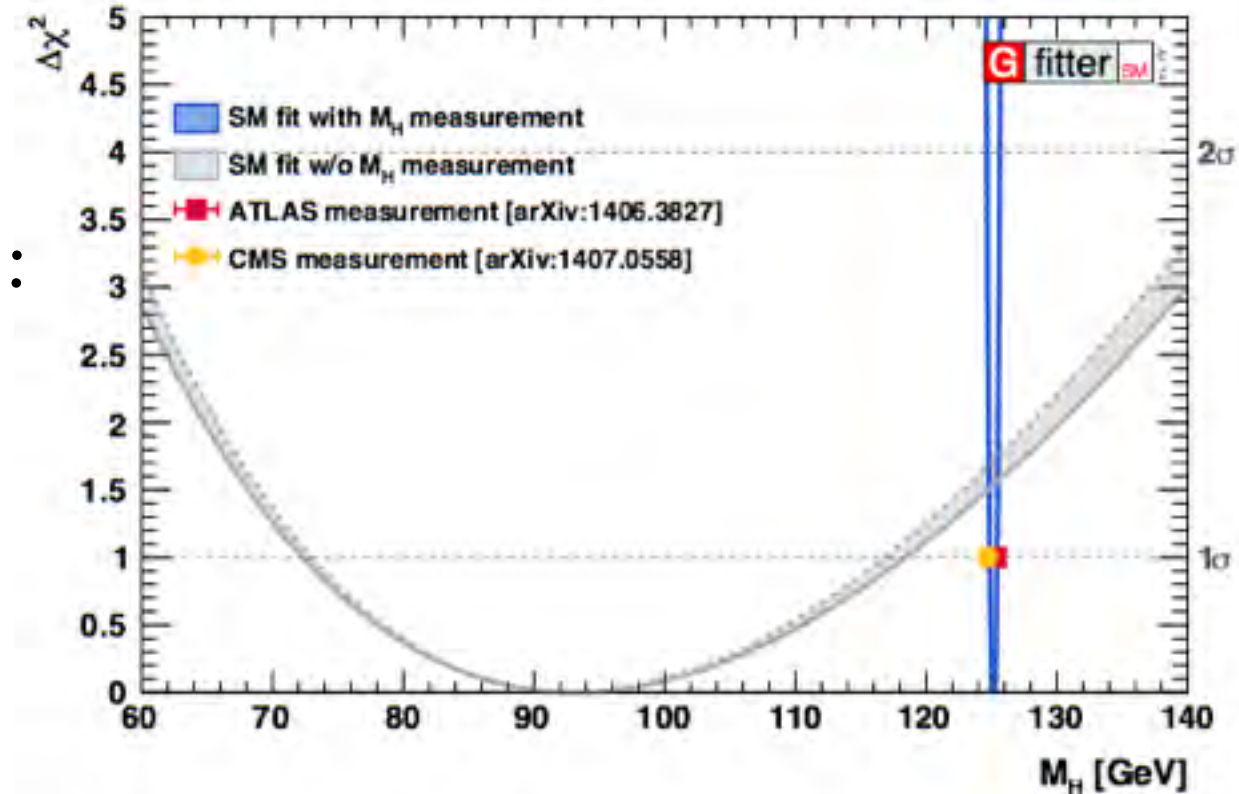
- ▶ goodness of fit, p-value:
 $\chi^2_{\min} = 17.8$ $\text{Prob}(\chi^2_{\min}, 14) = 21\%$
 Pseudo experiments: 21 ± 2 (theo)%
- $\chi^2_{\min}(\text{Z widths in 1-loop}) = 18.0$
- $\chi^2_{\min}(\text{no theory uncertainties}) = 18.2$
- ▶ no individual value exceeds 3σ
- ▶ largest deviations in b-sector:
 - $A^{0,b}_{\text{FB}}$ with 2.5σ
 \rightarrow largest contribution to χ^2
- ▶ small pulls for M_H, M_Z
 - input accuracies exceed fit requirements



Higgs results

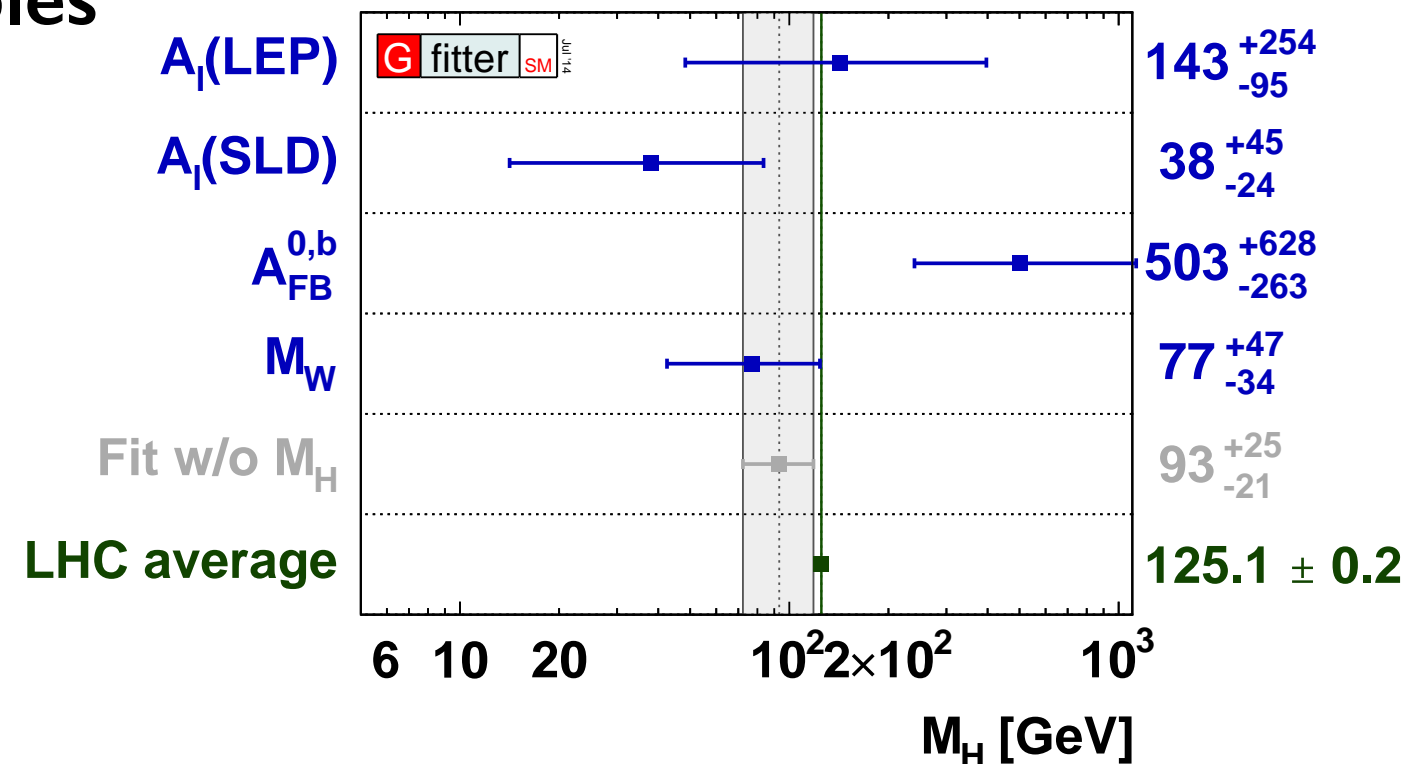
$\Delta\chi^2$ profile vs M_H

- ▶ grey band: fit without M_H measurement :
 - $M_H = 93^{+25}_{-21}$ GeV
 - consistent with measurement at 1.3σ
- ▶ blue line: full SM fit



impact of most sensitive observables

- ▶ determination of M_H , removing all sensitive observables except the given one
- ▶ known tension (3σ) between $A_l(\text{SLD})$, $A_{\text{FB}}^{0,b}$, and M_W clearly visible



Indirect determination of W mass

$\Delta\chi^2$ profile vs M_W

- ▶ also shown: SM fit with minimal input:

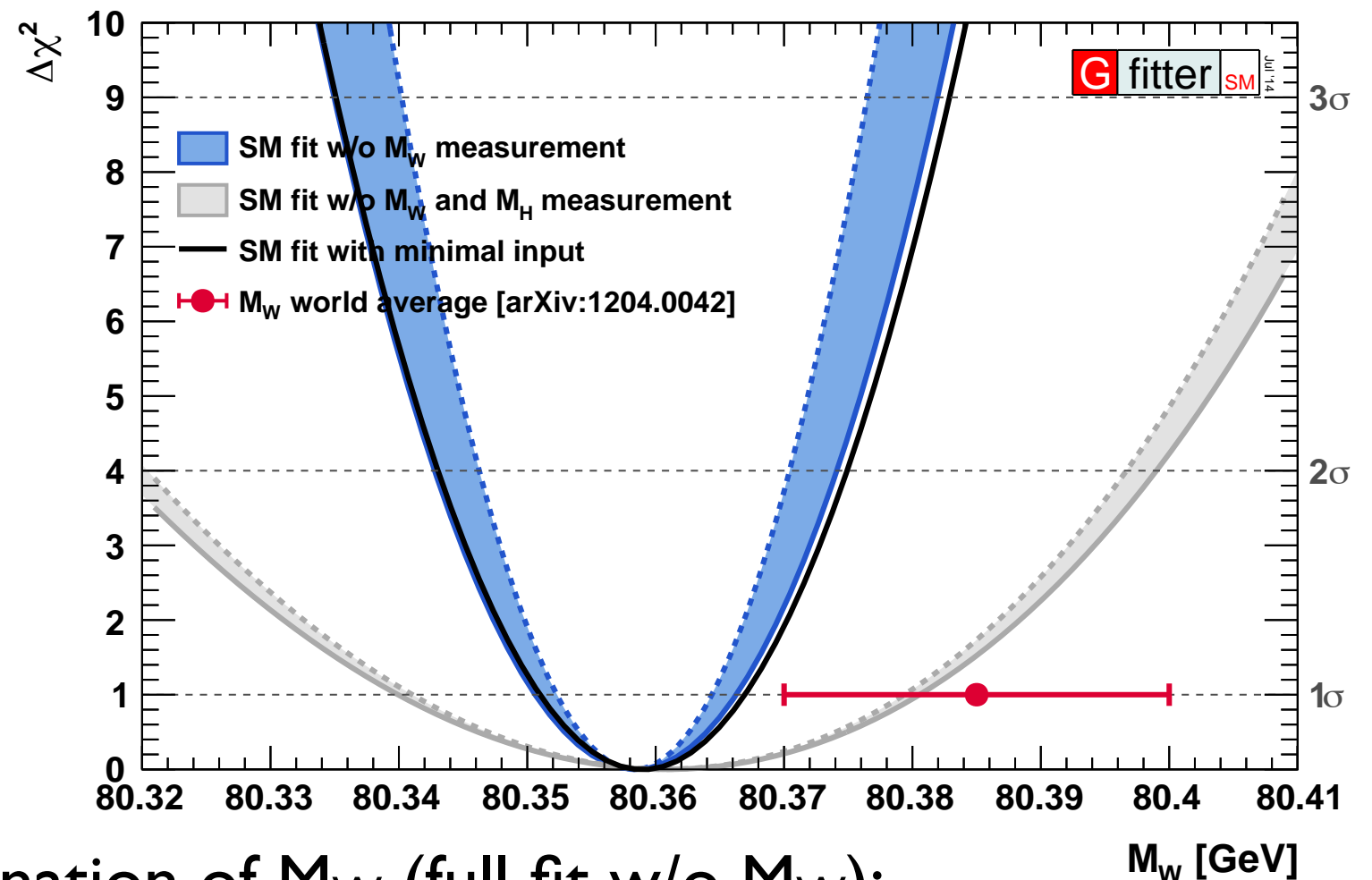
M_Z , G_F , $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$, $\alpha_s(M_Z)$, M_H , and fermion masses

- good consistency

- ▶ M_H measurement allows for precise constraint on M_W

- agreement at **1.4 σ**

- ▶ fit result for indirect determination of M_W (full fit w/o M_W):



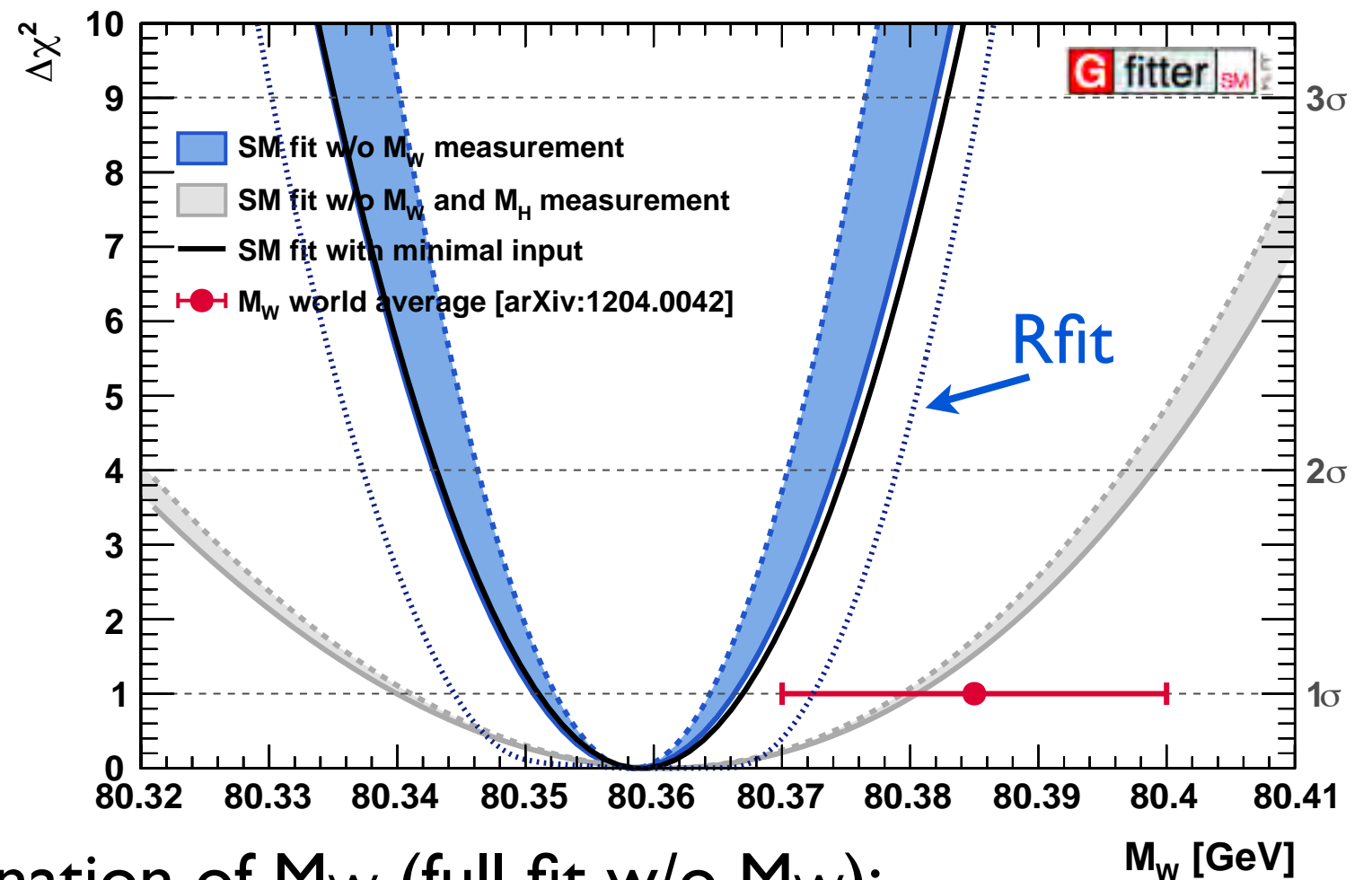
$$\begin{aligned}
 M_W &= 80.3584 \pm 0.0046_{m_t} \pm 0.0030_{\delta_{\text{theo}} m_t} \pm 0.0026_{M_Z} \pm 0.0018_{\Delta\alpha_{\text{had}}} \\
 &\quad \pm 0.0020_{\alpha_S} \pm 0.0001_{M_H} \pm 0.0040_{\delta_{\text{theo}} M_W} \text{ GeV}, \\
 &= 80.358 \pm 0.008_{\text{tot}} \text{ GeV}
 \end{aligned}$$

more precise than direct measurement (15 MeV)

Indirect determination of W mass

$\Delta\chi^2$ profile vs M_W

- ▶ also shown: SM fit with minimal input:
 $M_Z, G_F, \Delta\alpha_{\text{had}}^{(5)}(M_Z), \alpha_s(M_Z), M_H$, and fermion masses
 - good consistency
- ▶ M_H measurement allows for precise constraint on M_W
 - agreement at **1.4σ**
- ▶ fit result for indirect determination of M_W (full fit w/o M_W):



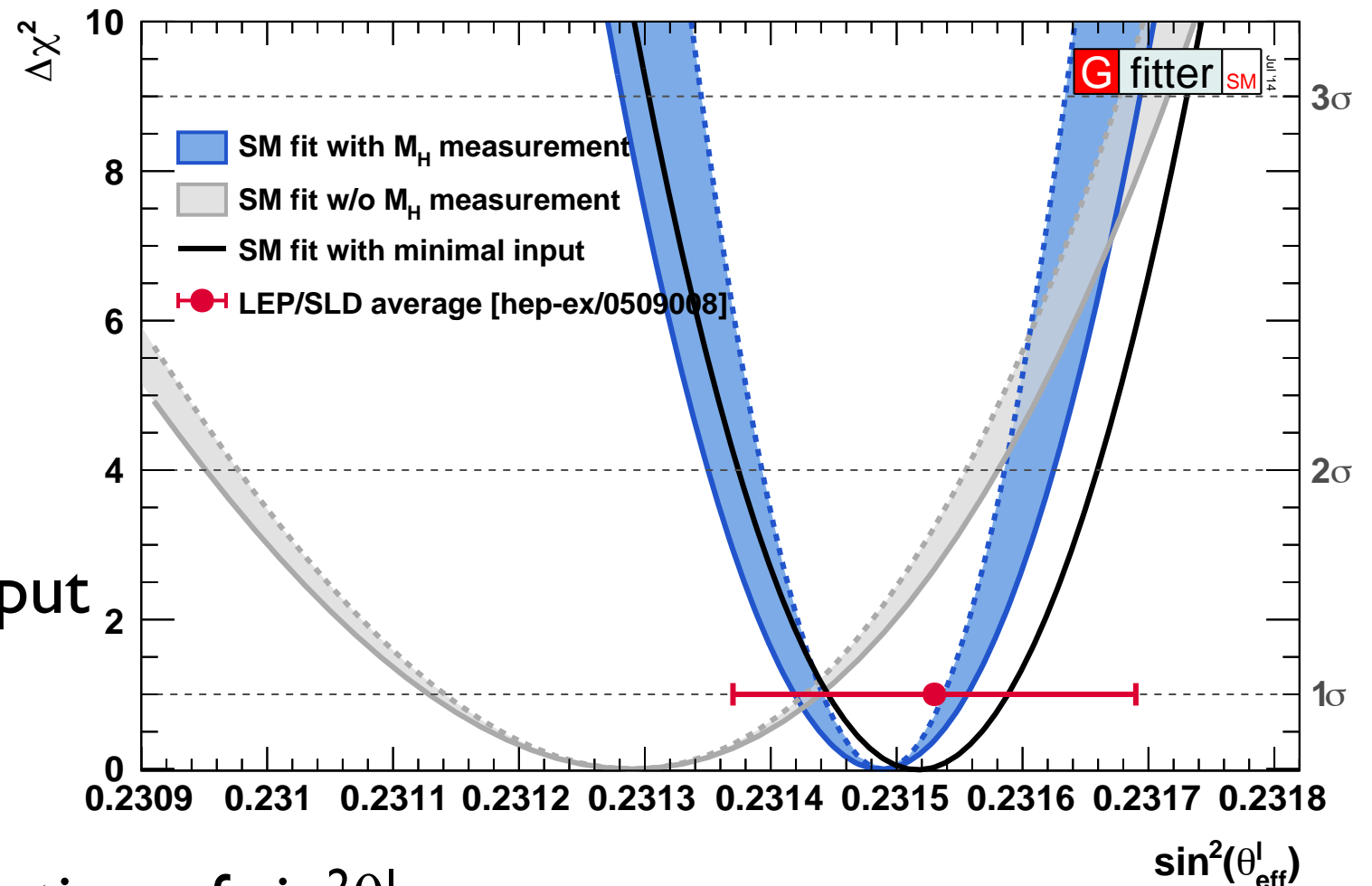
$$\begin{aligned}
 M_W &= 80.3584 \pm 0.0046_{m_t} \pm 0.0030_{\delta_{\text{theo}} m_t} \pm 0.0026_{M_Z} \pm 0.0018_{\Delta\alpha_{\text{had}}} \\
 &\quad \pm 0.0020_{\alpha_S} \pm 0.0001_{M_H} \pm 0.0040_{\delta_{\text{theo}} M_W} \text{ GeV}, \\
 &= 80.358 \pm 0.008_{\text{tot}} \text{ GeV} \quad (\text{Rfit: } \pm 13 \text{ MeV})
 \end{aligned}$$

more precise than direct measurement (15 MeV)

The effective weak mixing angle

$\Delta\chi^2$ profile vs $\sin^2\theta_{\text{eff}}^l$

- ▶ all measurements directly sensitive to $\sin^2\theta_{\text{eff}}^l$ removed from fit (asymmetries, partial widths)
 - good agreement with min input
- ▶ M_H measurement allows for precise constraint
- ▶ fit result for indirect determination of $\sin^2\theta_{\text{eff}}^l$:



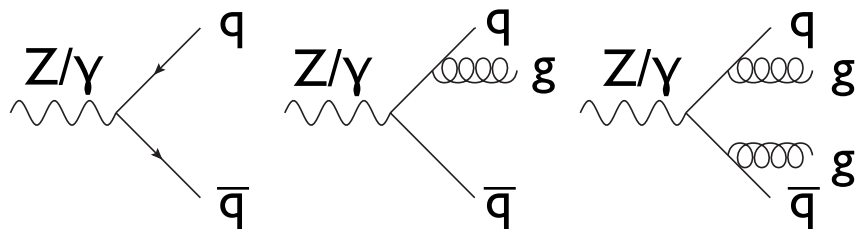
$$\begin{aligned}
 \sin^2\theta_{\text{eff}}^l &= 0.231488 \pm 0.000024_{m_t} \pm 0.000016_{\delta_{\text{theo}} m_t} \pm 0.000015_{M_Z} \pm 0.000035_{\Delta\alpha_{\text{had}}} \\
 &\quad \pm 0.000010_{\alpha_S} \pm 0.000001_{M_H} \pm 0.000047_{\delta_{\text{theo}} \sin^2\theta_{\text{eff}}^f} \\
 &= 0.23149 \pm 0.00007_{\text{tot}}
 \end{aligned}$$

more precise than determination from LEP/SLD (1.6×10^{-4})

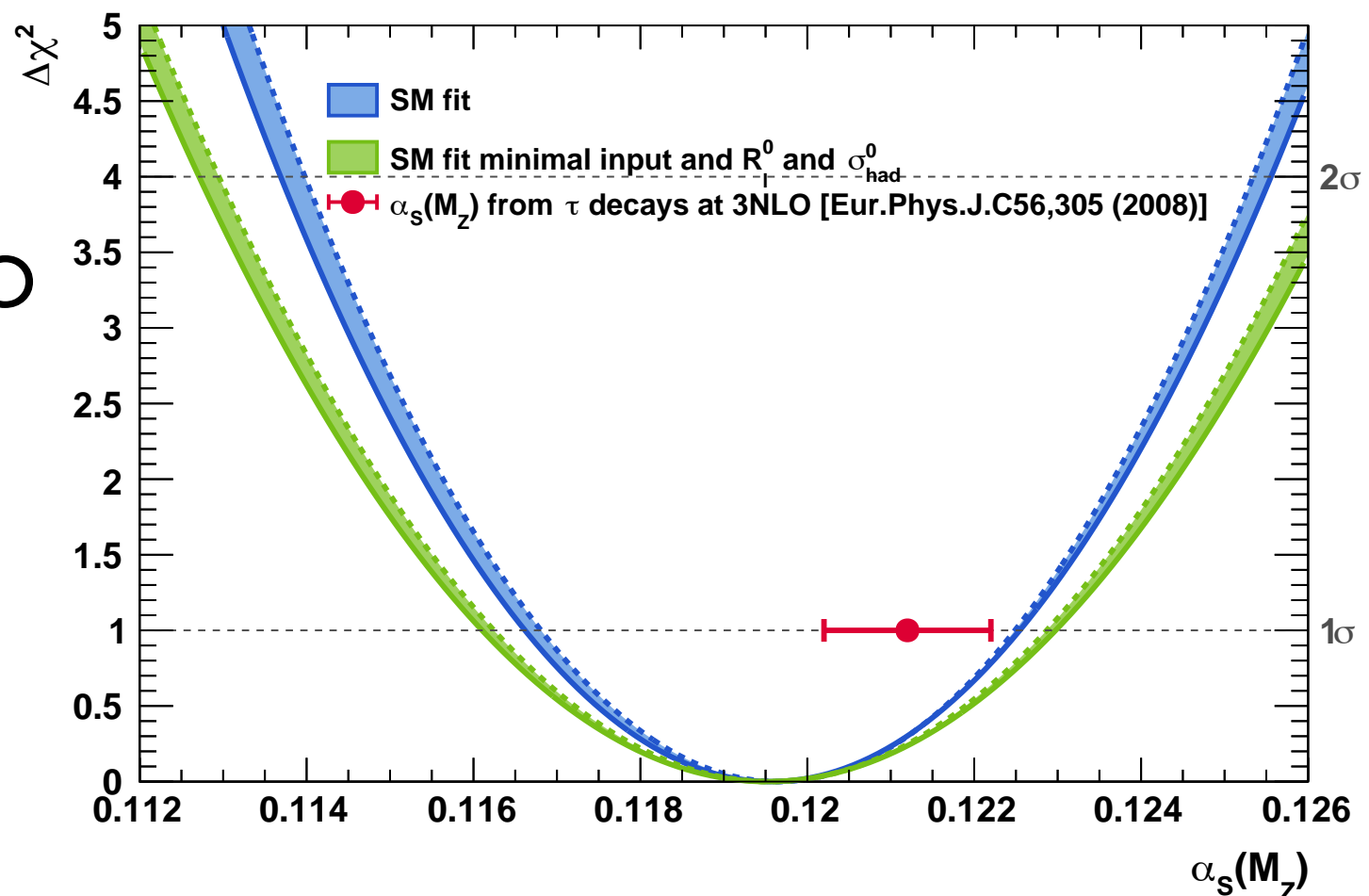
The strong coupling $\alpha_s(M_Z)$

$\Delta\chi^2$ profile vs $\alpha_s(M_Z)$

- ▶ determination of α_s at full NNLO and partial NNNLO
- ▶ also shown: minimal input with two most sensitive measurements: $R_l, \sigma_{\text{had}}^0$



- ▶ M_H has no (visible) impact



$$\begin{aligned}\alpha_s(M_Z^2) &= 0.1196 \pm 0.0028_{\text{exp}} \pm 0.0006_{\delta_{\text{theo}} \mathcal{R}_{V,A}} \pm 0.0006_{\delta_{\text{theo}} \Gamma_i} \pm 0.0002_{\delta_{\text{theo}} \sigma_{\text{had}}^0} \\ &= \underline{0.1196 \pm 0.0030_{\text{tot}}}\end{aligned}$$

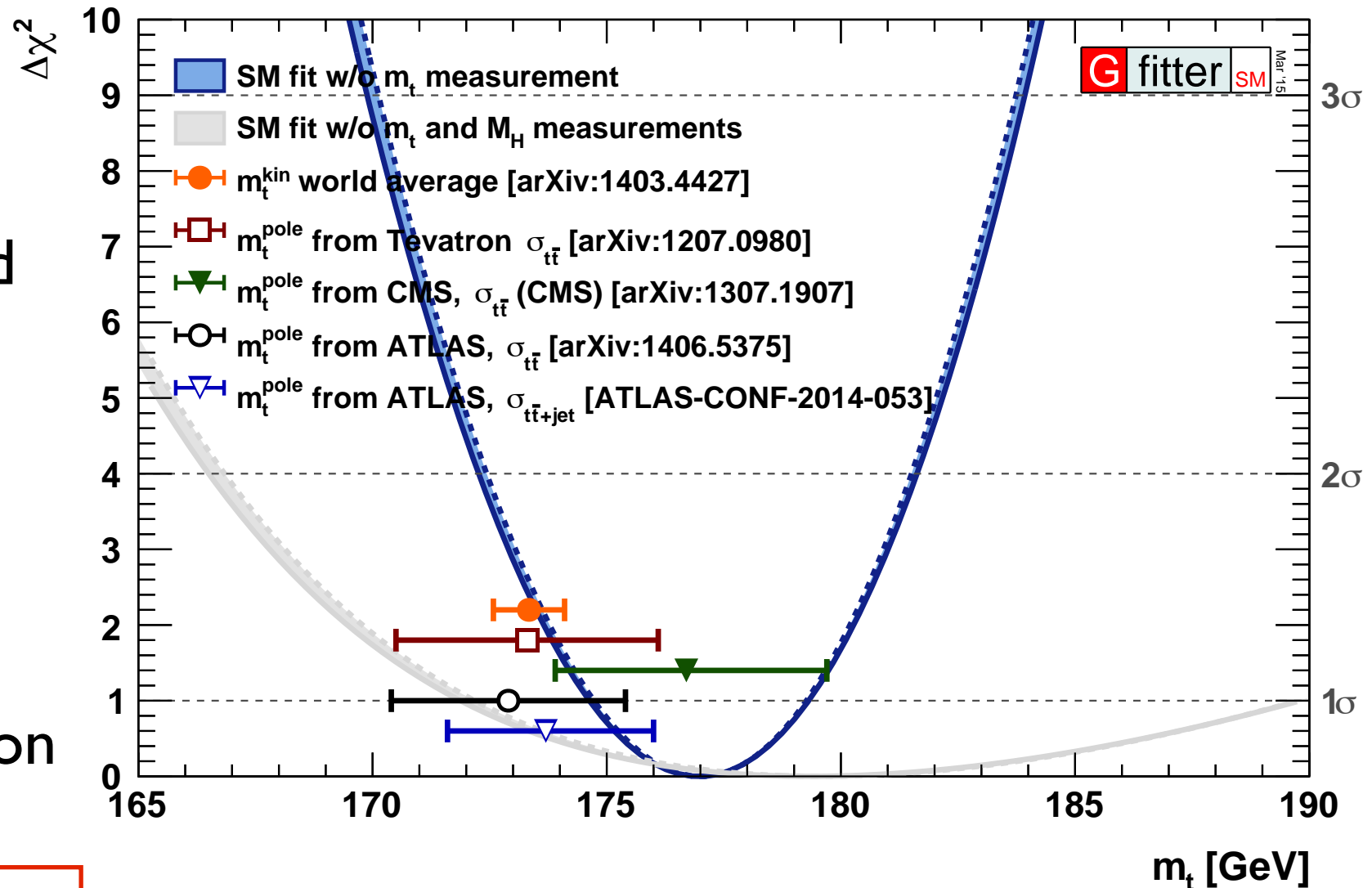
More accurate estimation of theo. uncertainties
(previously: $\delta_{\text{theo}} = 0.0001$ from scale variations)

good agreement with WA, dominated by exp. uncertainty

Indirect determination of m_t

$\Delta\chi^2$ profile vs m_t

- ▶ determination of m_t from Z-pole data (fully obtained from rad. corrections $\sim m_t^2$)
- ▶ alternative to direct measurements
- ▶ M_H allows for significantly more precise determination of m_t

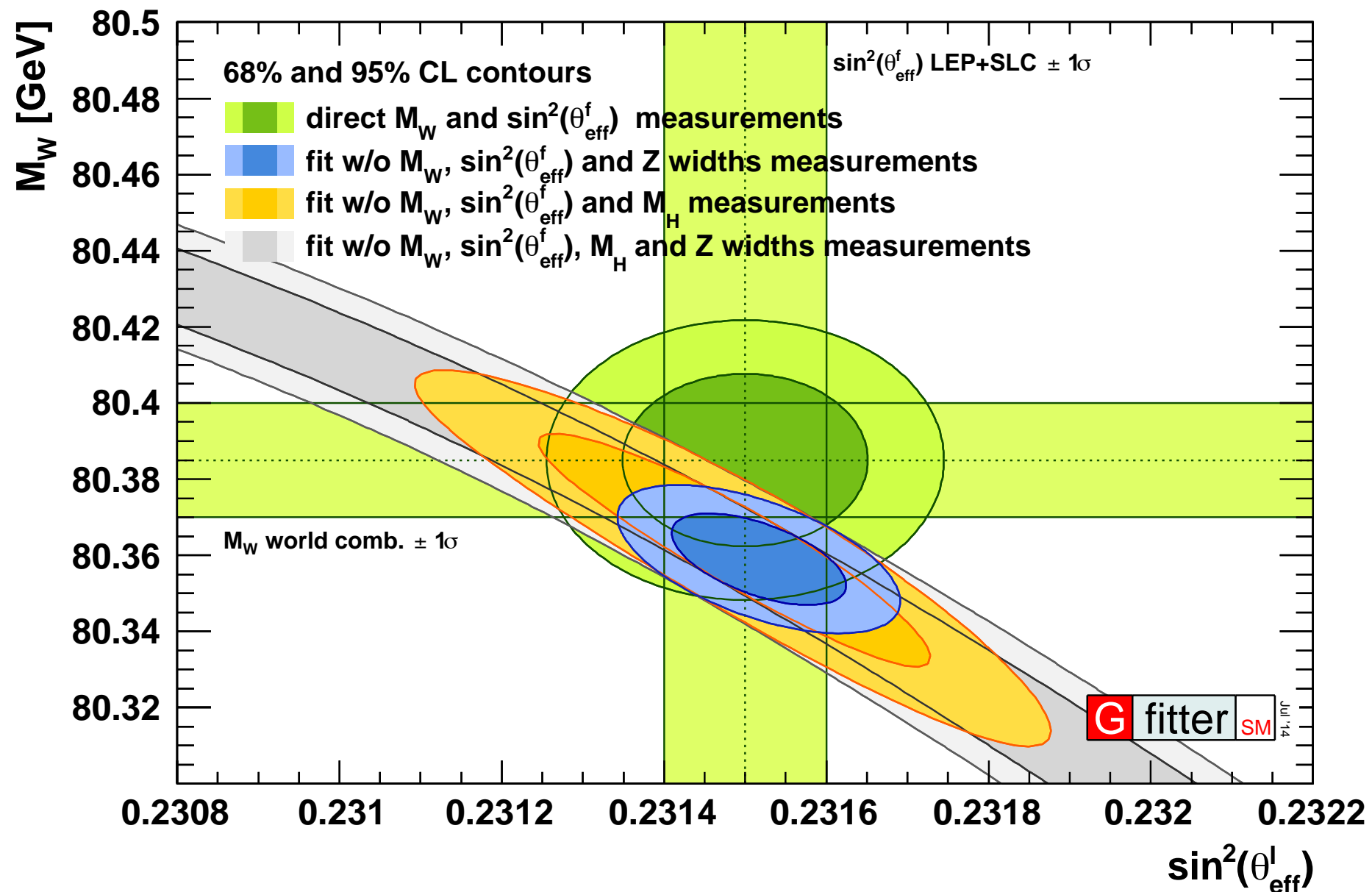


$$m_t = 177.0 \pm 2.3_{M_W, \sin^2 \theta_{\text{eff}}^f} \pm 0.6_{\alpha_s} \pm 0.5_{\Delta \alpha_{\text{had}}} \pm 0.4_{M_Z} \text{ GeV}$$

$$= 177.0 \pm 2.4_{\text{exp}} \pm 0.5_{\text{theo}} \text{ GeV}$$

- ▶ similar precision as determination from $\sigma_{t\bar{t}}$, good agreement
- ▶ dominated by experimental precision

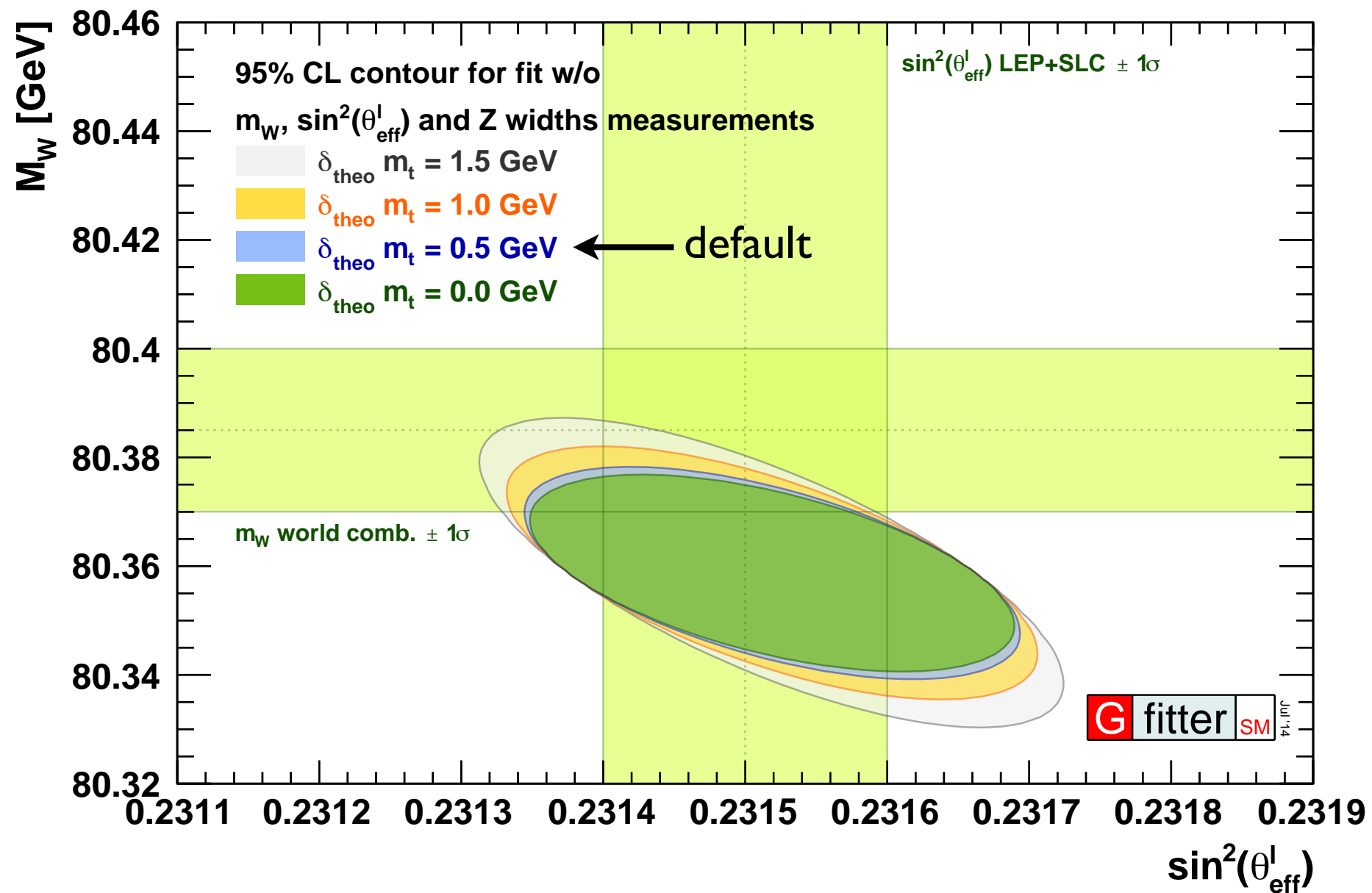
State of the SM: M_W vs $\sin^2\theta_{\text{eff}}^l$



sensitive probes of new physics

- ▶ significant reduction of parameter space due to knowledge of M_H
- ▶ predictions are more precise than the direct measurements

Theoretical uncertainty on m_t



impact of variation in $\delta_{\text{theo}} m_t$ between 0 and 1.5 GeV

- ▶ better assessment of uncertainty on m_t important for the fit
- ▶ uncertainty of 0.5 GeV small impact on result

Constraints on BSM models

- ▶ if energy scale of NP is high, BSM physics could appear dominantly through vacuum polarisation corrections

- ▶ described by STU parameters
[Peskin and Takeuchi, Phys. Rev. D46, 1 (1991)]

- ▶ SM: $M_H = 125 \text{ GeV}$, $m_t = 173 \text{ GeV}$
this defines $(S, T, U) = (0, 0, 0)$

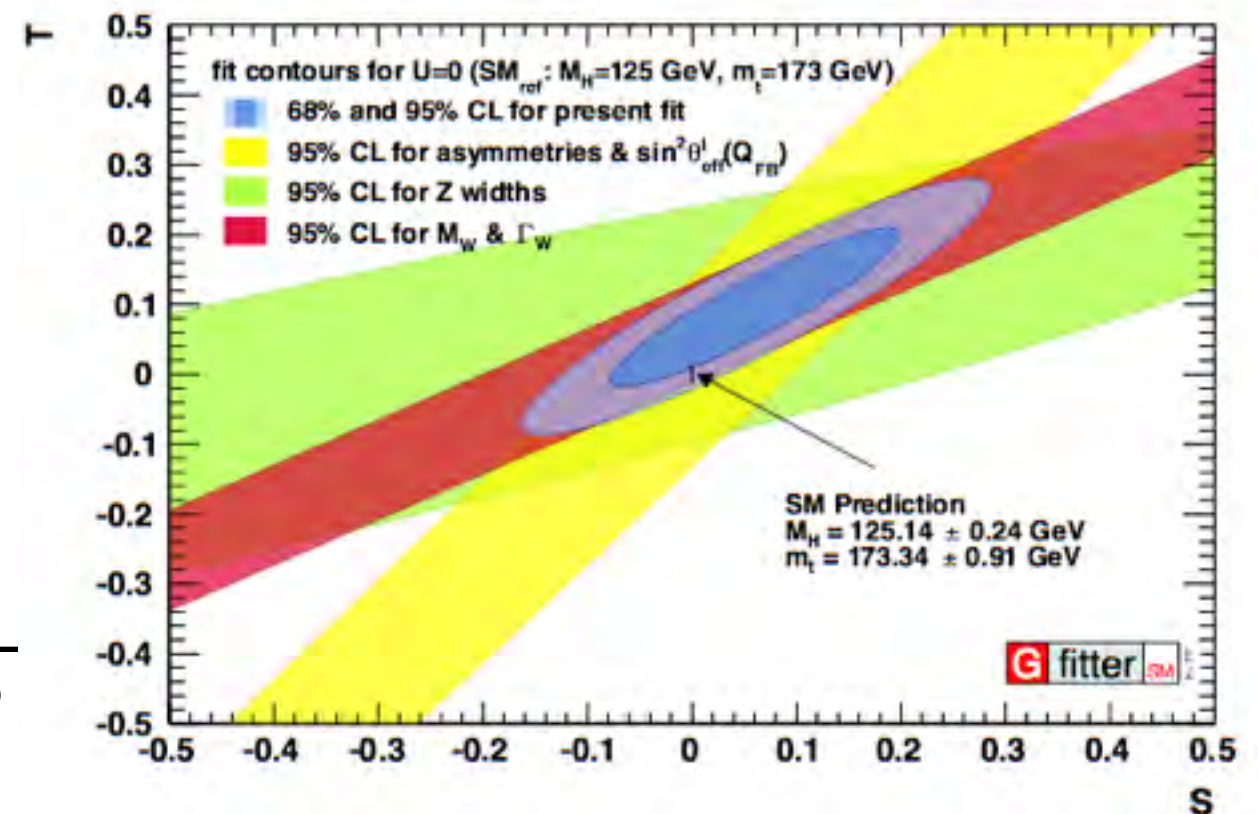
- ▶ S, T depend logarithmically on M_H

- ▶ Fit result:

	S	T	U
$S = 0.05 \pm 0.11$	S	+0.90	-0.59
$T = 0.09 \pm 0.13$	T	I	-0.83
$U = 0.01 \pm 0.11$	U		I

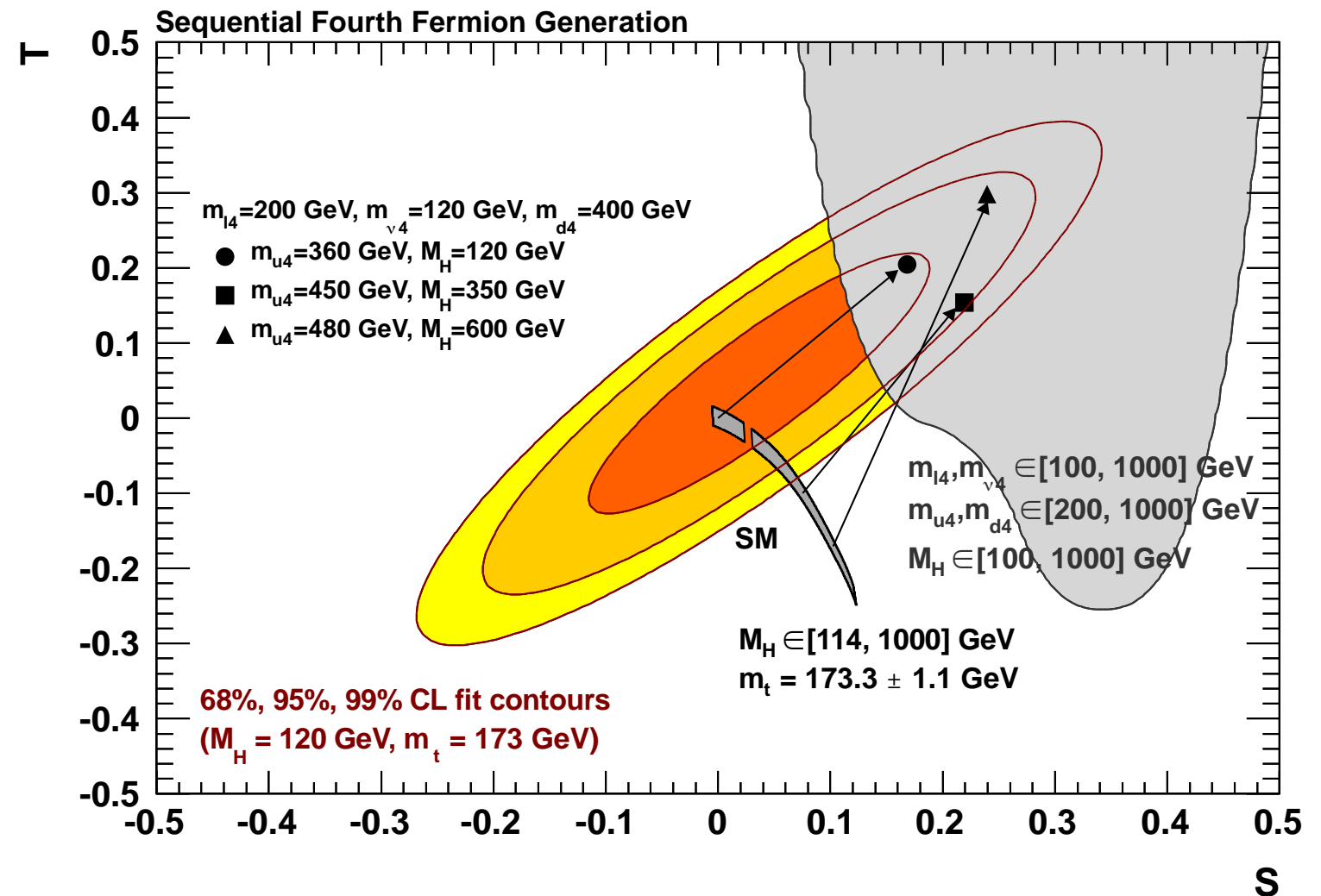
- ▶ no indication for new physics
- ▶ use this to constrain parameter space in BSM models

stronger constraints with $U = 0$:



Constraints on BSM models

- ▶ with M_H unknown, changes in S, T and U could often be compensated by changes in M_H
- ▶ rather weak limits: e.g. large parameter space for sequential fourth generation open

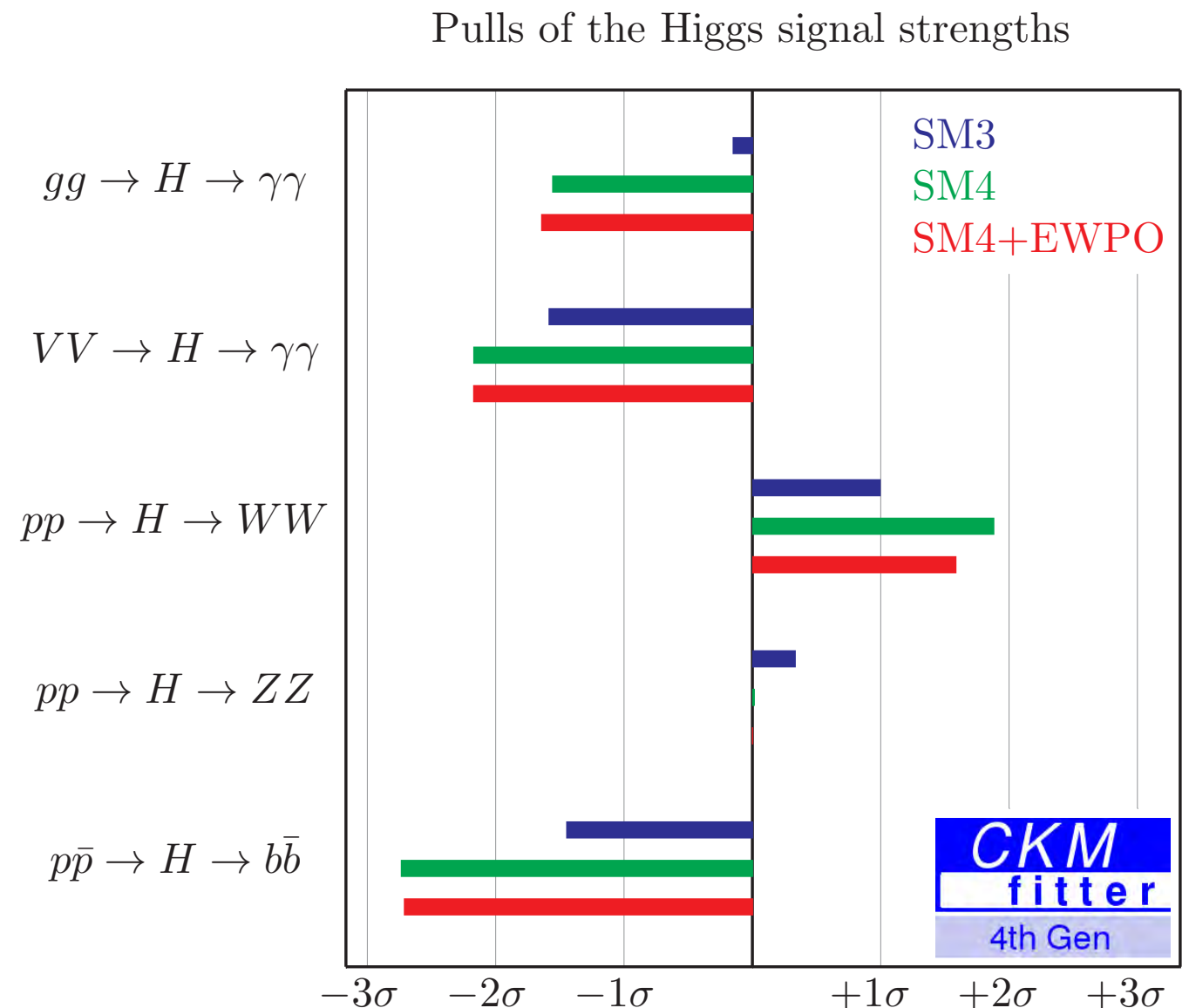


Constraints on BSM models

- ▶ with M_H unknown, changes in S, T and U could often be compensated by changes in M_H
- ▶ rather weak limits: e.g. large parameter space for sequential fourth generation open

- ▶ after discovery of a SM-like Higgs boson:
chiral 4th generation ruled out
[O. Eberhard et al., PRL 109, 241802 (2012)]

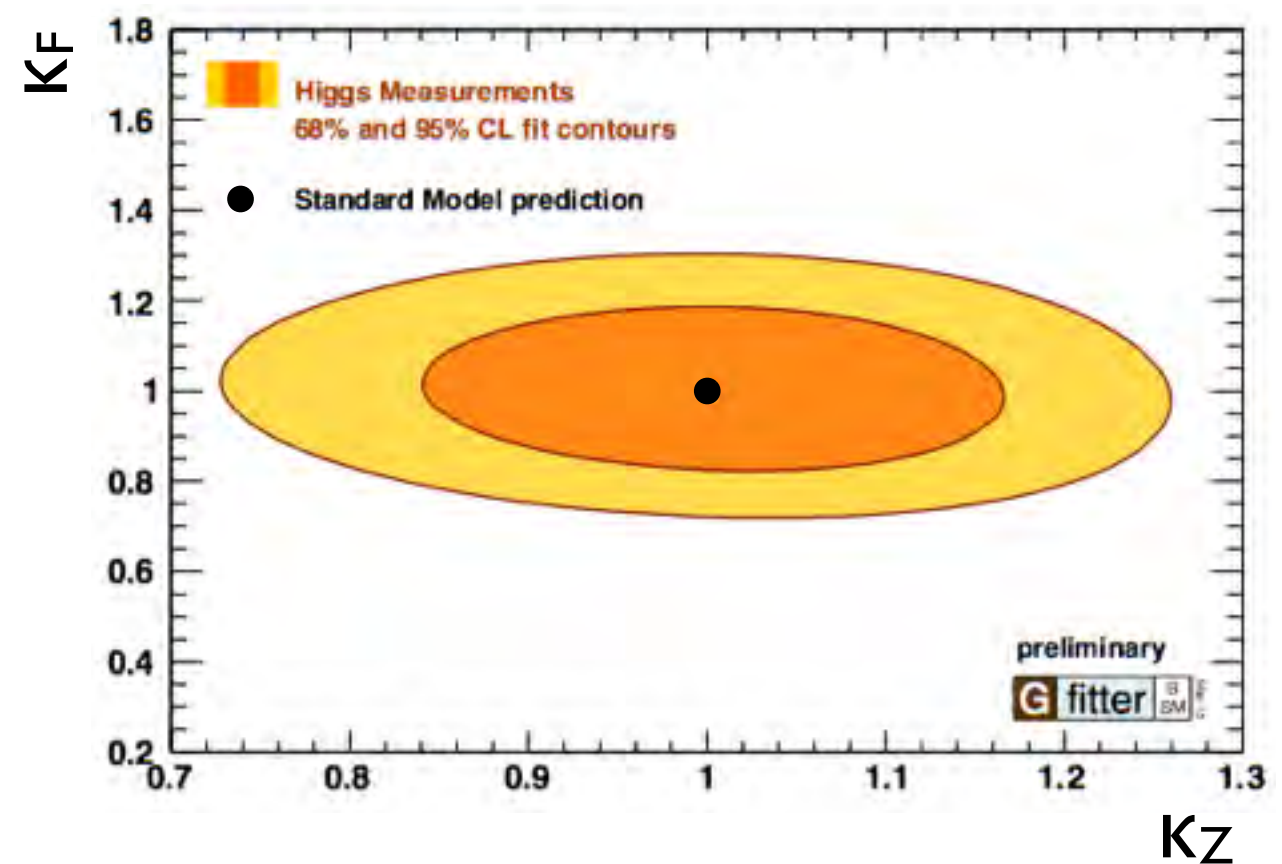
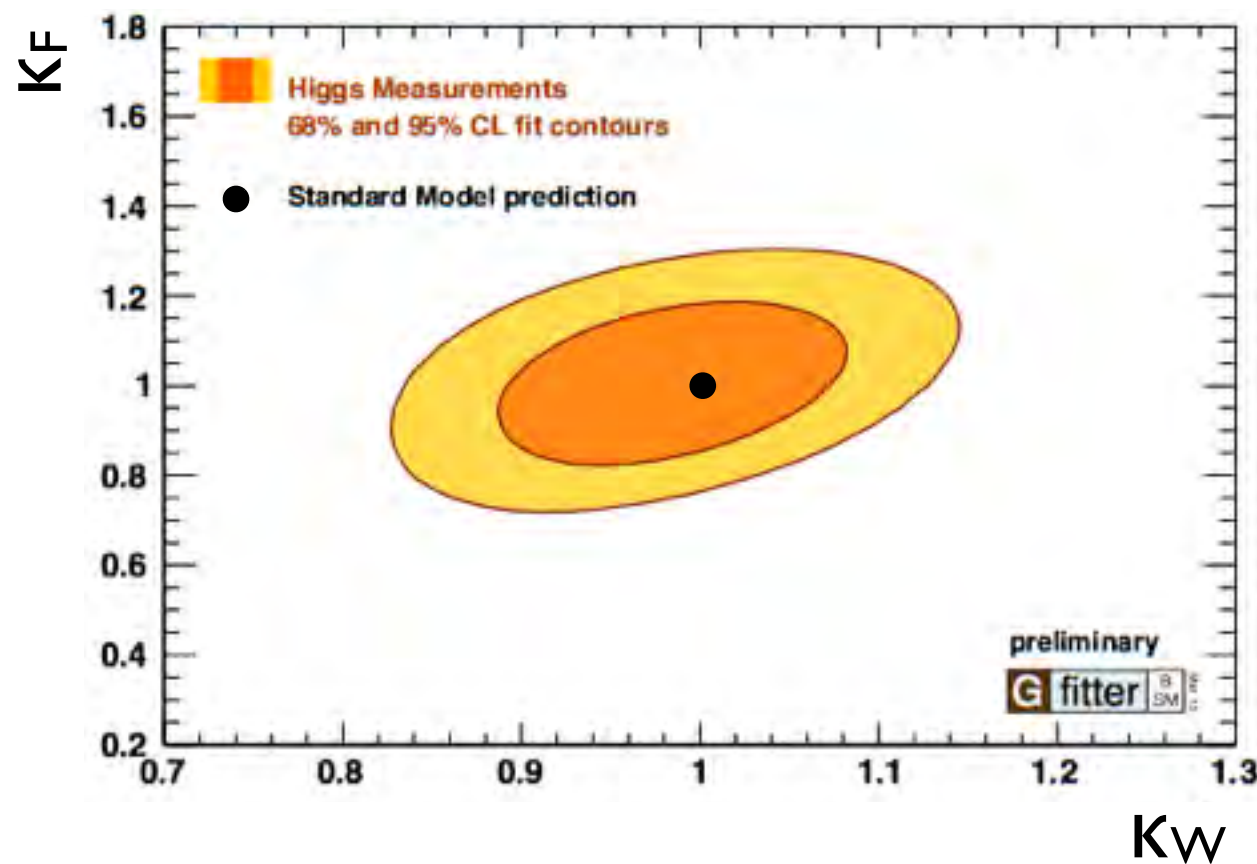
- ▶ note: mostly from Higgs signal strength, small impact of EWPO



The Scalar Sector

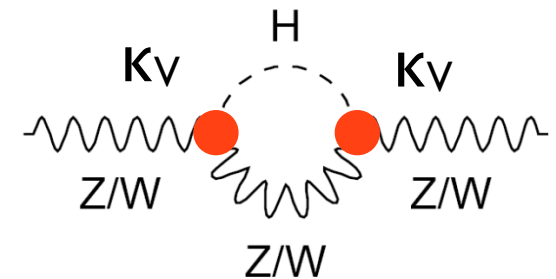
Tree Level Higgs Couplings

- ▶ study of potential deviations of Higgs couplings from SM
- ▶ leading corrections only, parametrize deviations with effective couplings
- ▶ LHC and Tevatron data included using HiggsSignals [\[P. Bechtle et al., JHEP 11, 039 \(2014\)\]](#)



- ▶ no BSM contributions on tree-level to fermion or vector-boson coupling
- ▶ stronger constraints on K_W than on K_Z
- ▶ custodial symmetry holds, $K_W = K_Z = K_V$

Constraints from EWPD



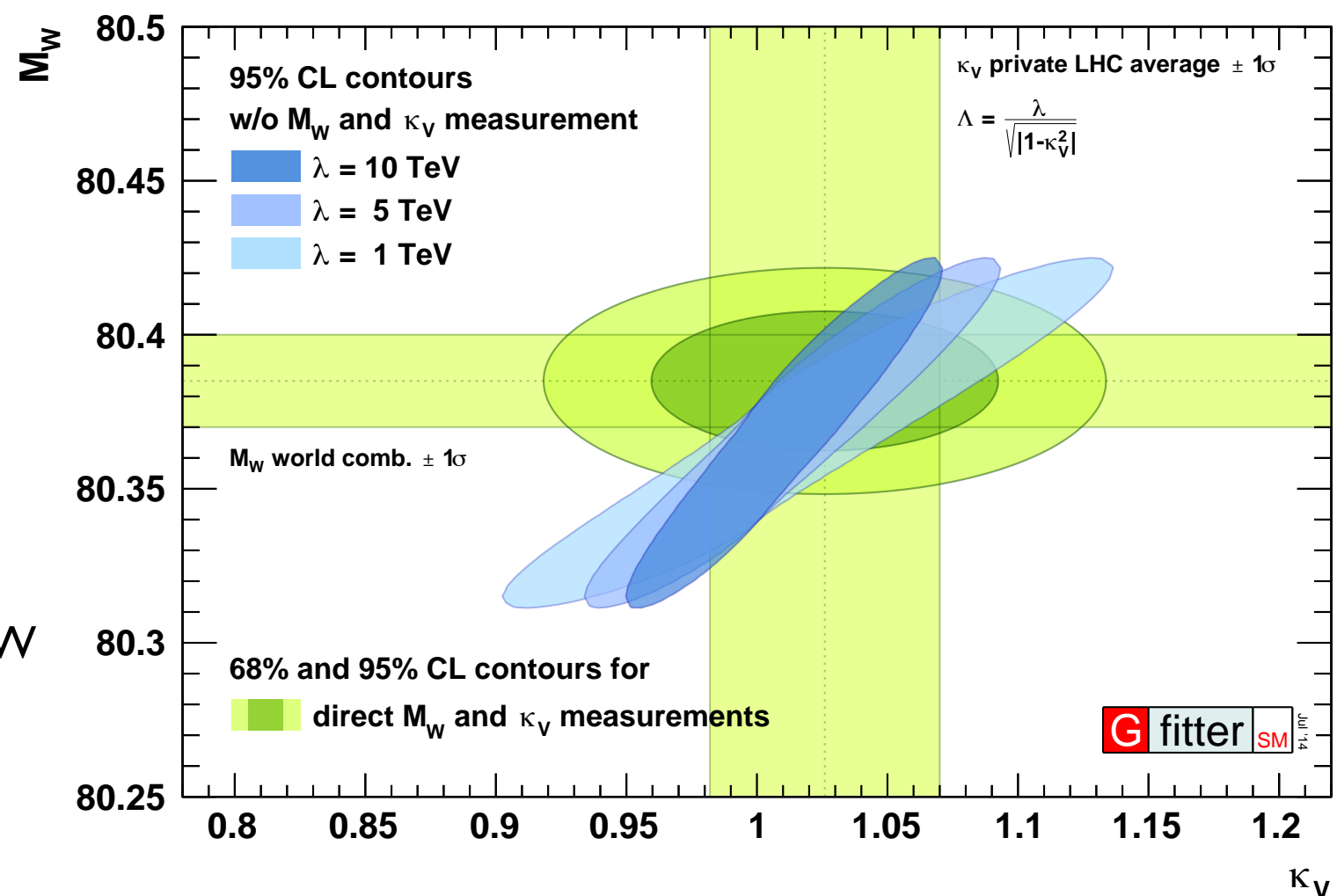
- ▶ consider specific model in “ κ parametrisation”:
 - scaling of Higgs-vector boson (κ_V) and Higgs-fermion couplings (κ_F), with no invisible/undetected widths
- ▶ main effect on EWPD due to modified Higgs coupling to gauge bosons (κ_V)
[\[Espinosa et al. arXiv:1202.3697, Falkowski et al. arXiv:1303.1812\], etc](#)

$$S = \frac{1}{12\pi} (1 - \kappa_V^2) \ln \frac{\Lambda^2}{M_H^2}$$

$$T = -\frac{3}{16\pi \cos^2 \theta_{\text{eff}}^\ell} (1 - \kappa_V^2) \ln \frac{\Lambda^2}{M_H^2}$$

$$\Lambda = \frac{\lambda}{\sqrt{|1 - \kappa_V^2|}}$$

- ▶ correlation between κ_V and M_W
 - slightly smaller values of M_W preferred



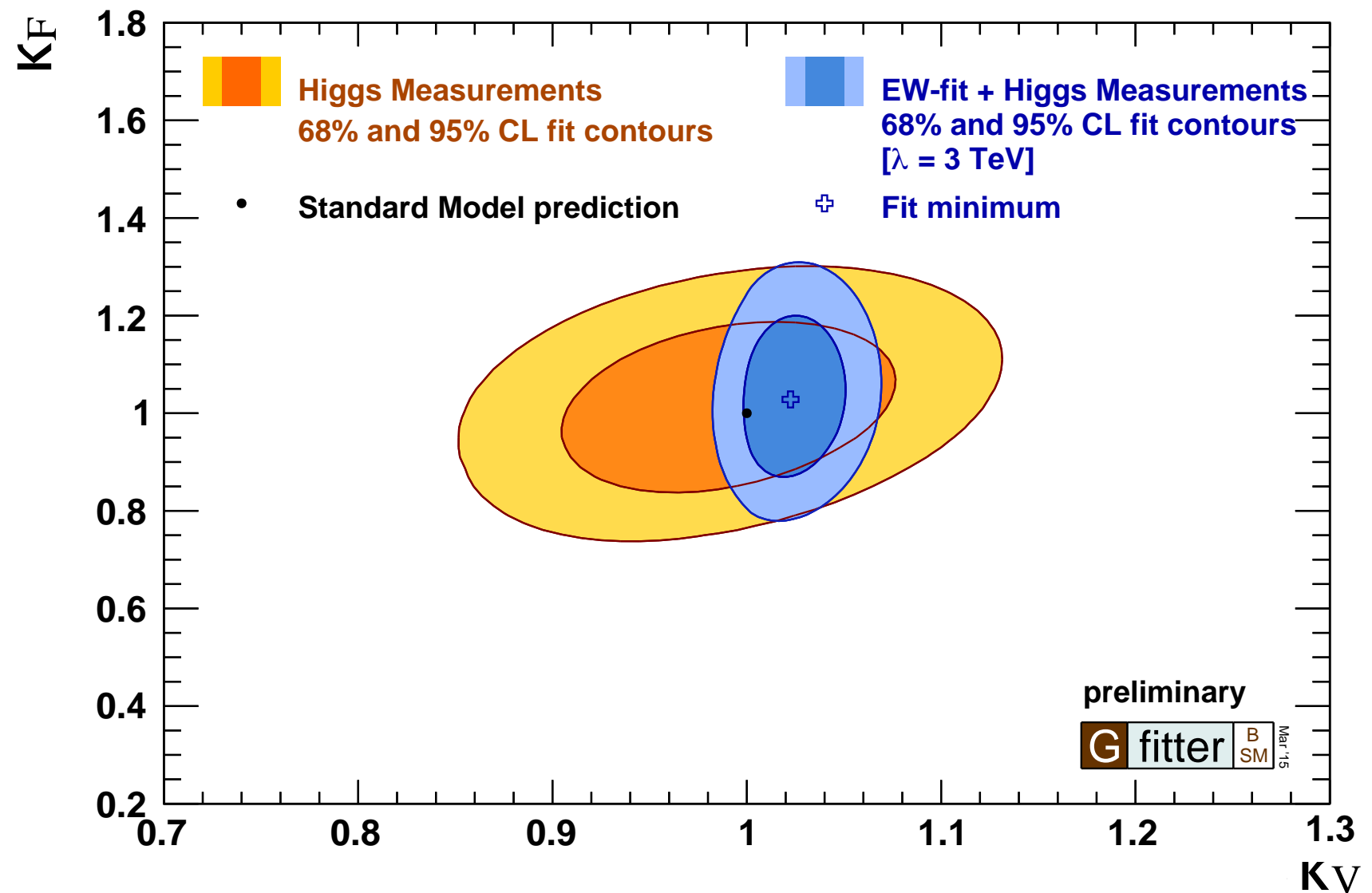
Higgs Coupling Results

Higgs coupling measurements:

- ▶ $\kappa_V = 0.99 \pm 0.08$
- ▶ $\kappa_F = 1.01 \pm 0.17$

▶ Combined result:

- ▶ $\kappa_V = 1.03 \pm 0.02$
($\lambda = 3 \text{ TeV}$)
- ▶ implies NP-scale of
 $\Lambda \geq 13 \text{ TeV}$



- ▶ some dependency for κ_V in central value [1.02-1.04] and error [0.02-0.03] on cut-off scale λ [1-10 TeV]
 - EW fit sofar more precise result for κ_V than current LHC experiments
 - EW fit has positive deviation of κ_V from 1.0
 - many BSM models: $\kappa_V < 1$

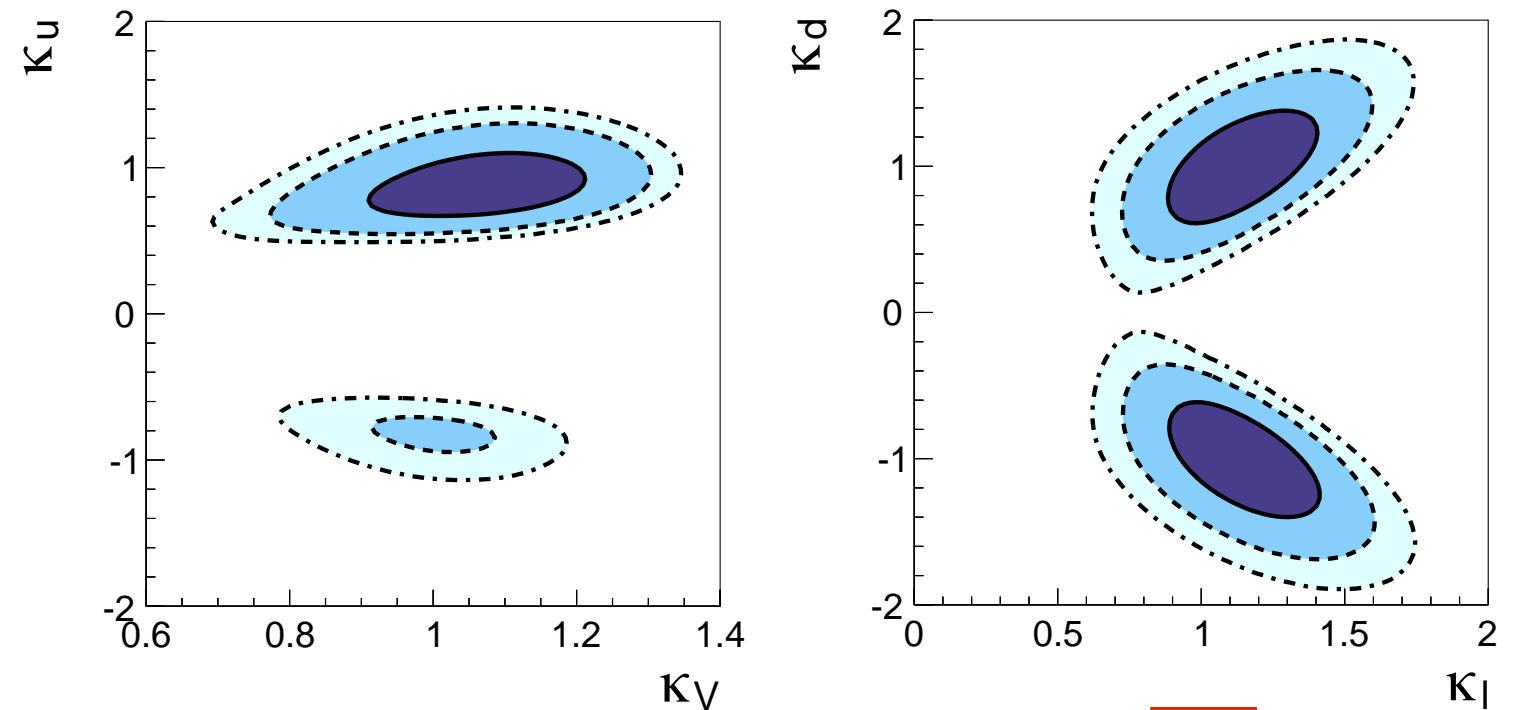
Higgs coupling results

- ▶ allowing for different couplings to up- and down-type quarks κ_u and κ_d
- ▶ stricter constraints due to EWPO, some gain also in the fermion sector

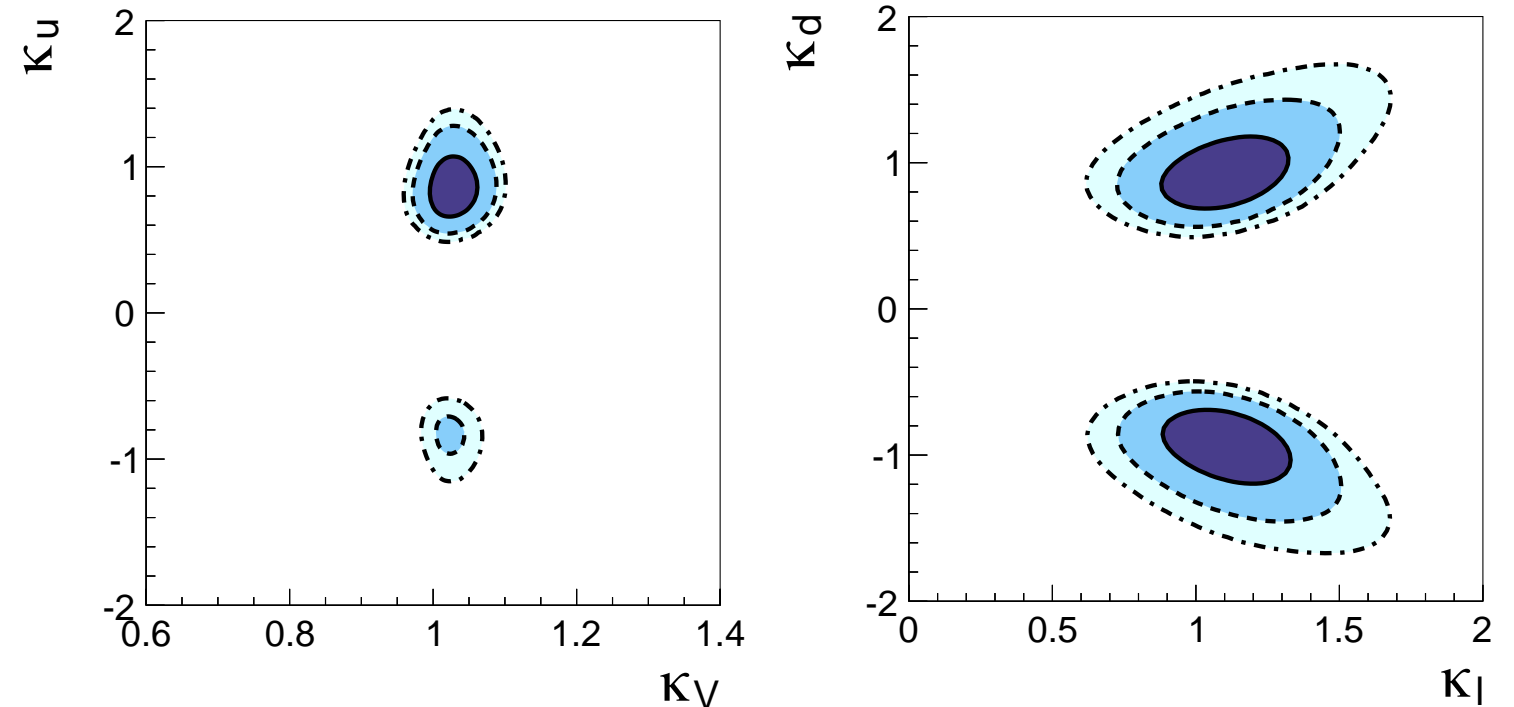
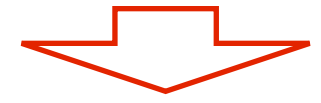
	68%	95%	Correlations			
κ_V	1.03 ± 0.02	[0.99, 1.07]	1.00			
κ_ℓ	1.10 ± 0.14	[0.82, 1.38]	0.14	1.00		
κ_u	0.88 ± 0.12	[0.66, 1.15]	0.09	0.23	1.00	
κ_d	0.92 ± 0.15	[0.65, 1.26]	0.28	0.35	0.81	1.00

- ▶ also possible to constrain coefficients of **dimension-6 operators**
 - contributions to EWPO have been worked out
 - theoretically sounder than constraints from S,T,U

only Higgs signal strength



+ EWPO

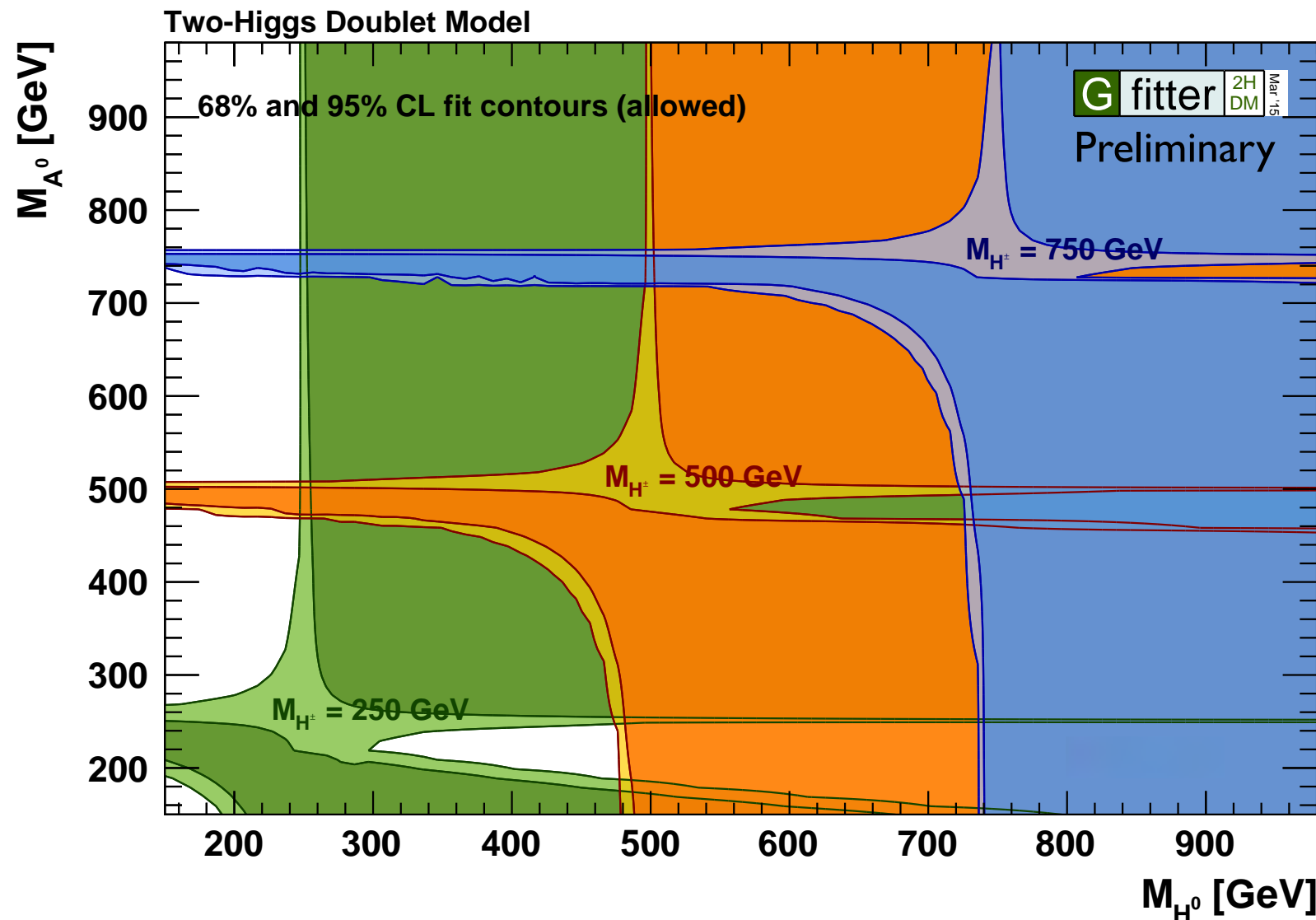


[Marco Ciuchini et al, arXiv:1410.6940]

Two Higgs Doublet Models

- ▶ extend the scalar sector by another doublet
- ▶ studies of Z_2 Type-I and Type-2 2HDMs
 - difference in the coupling to down-type quarks
 - Type-2 related to MSSM, but less constrained

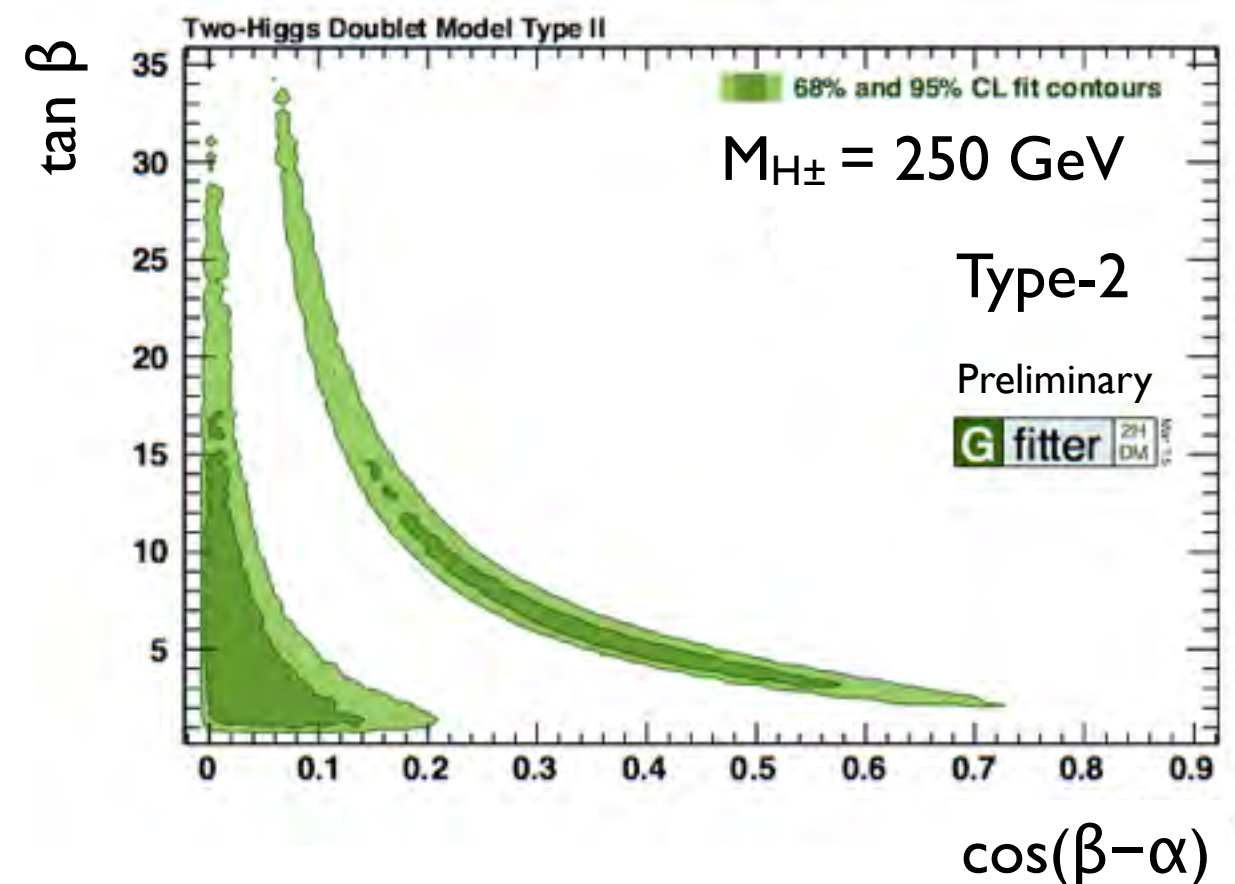
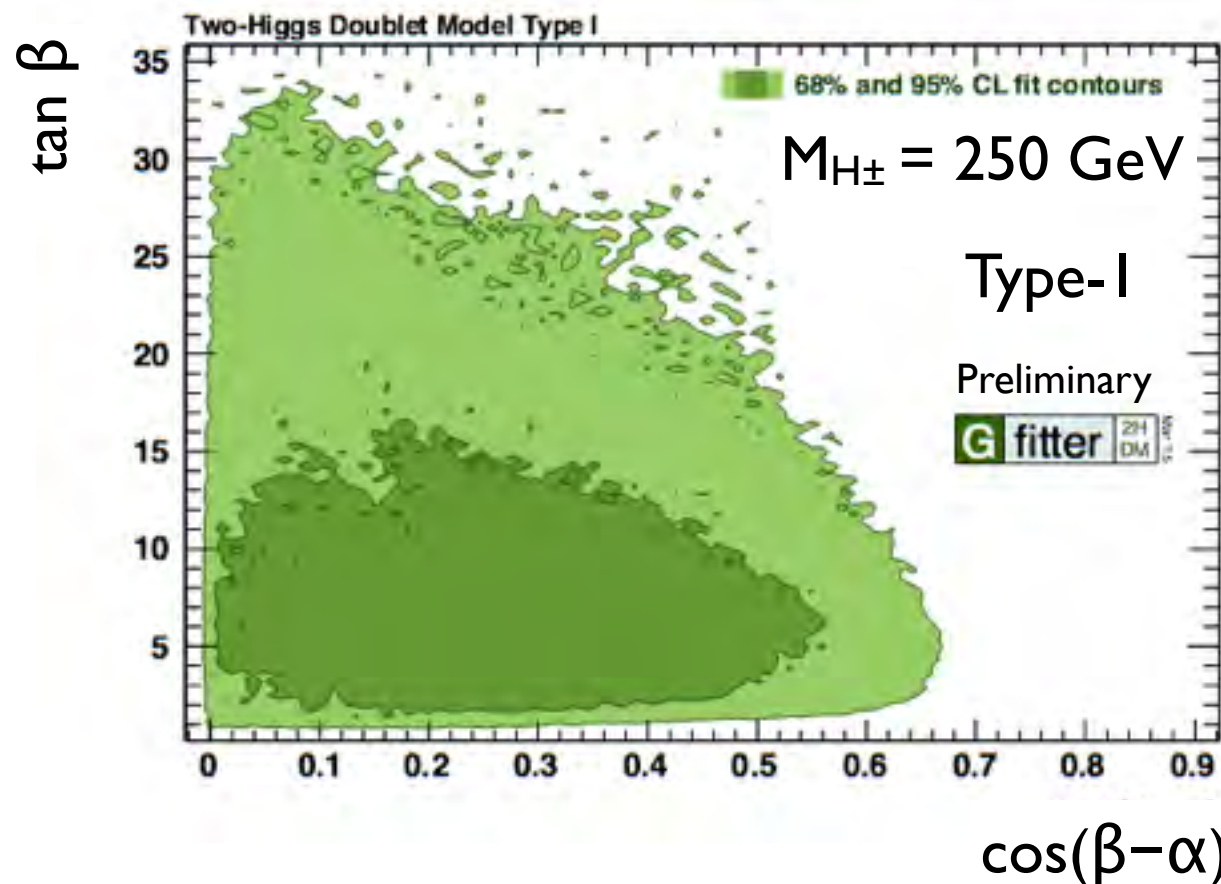
	Type I and Type II
Higgs	C_V
h	$\sin(\beta - \alpha)$
H	$\cos(\beta - \alpha)$
A	0



- ▶ constraints derived from EWPD using S,T,U formalism
- ▶ lightest scalar $M_h = 125.1 \text{ GeV}$
- ▶ weak constraints on masses, since $\tan\beta$ and $\cos(\beta - \alpha)$ are unconstrained

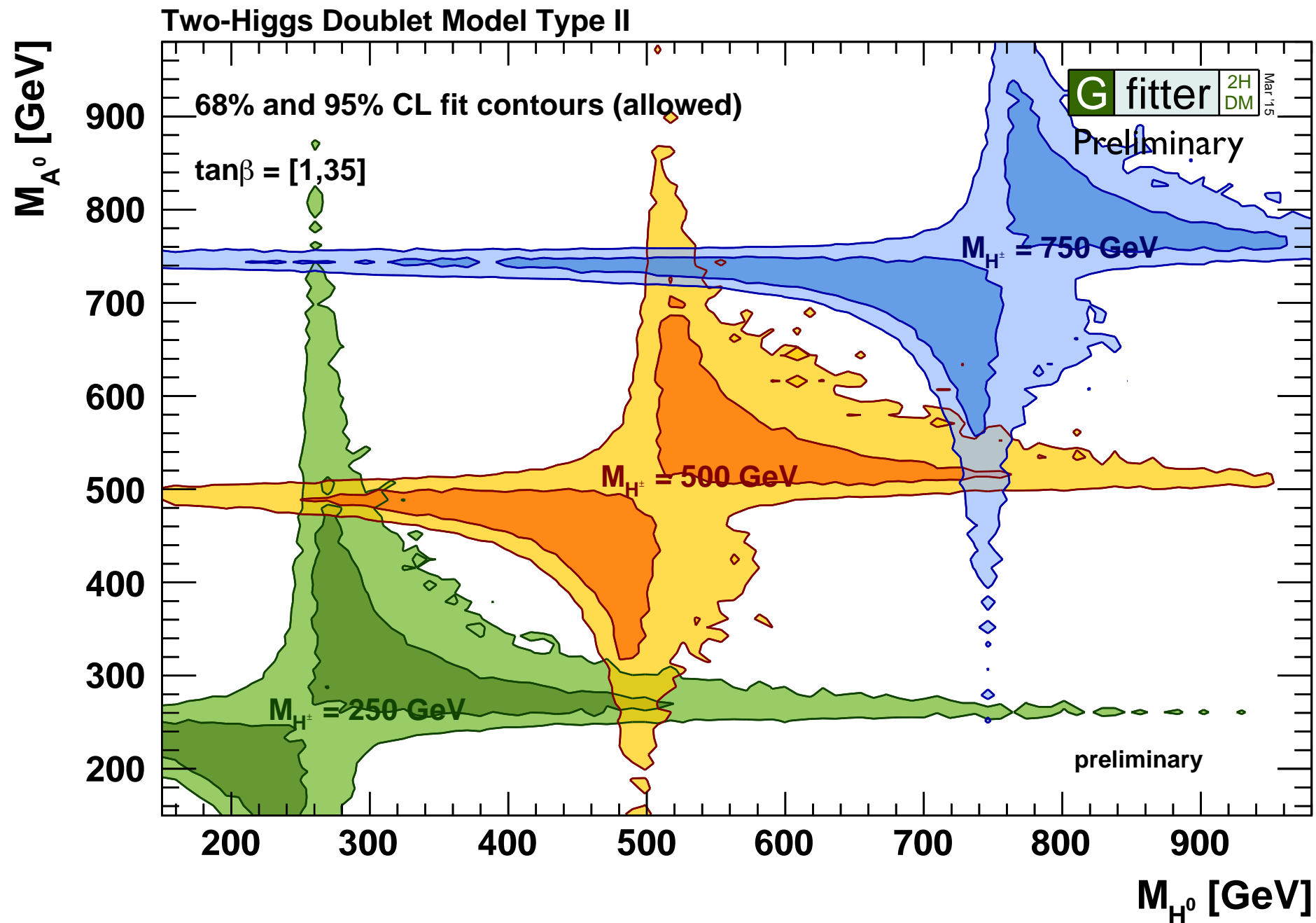
2HDM and H Coupling Measurements

- ▶ coupling measurements place important constraints on 2HDMs
- ▶ predictions of BRs using 2HDMC [D. Eriksson et al., CPC 181, 189 (2010)]
- ▶ 7 additional, unconstrained parameters (4 masses, 2 angles, soft breaking scale): importance sampling with MultiNest [F. Feroz et al., arXiv:1306.2144]



- ▶ additional constraints from flavour data
 - $B \rightarrow X_s \gamma$: $\tan \beta > 1$
 - $B_s \rightarrow \mu\mu$: constraints depending on M_H and M_{H^\pm}

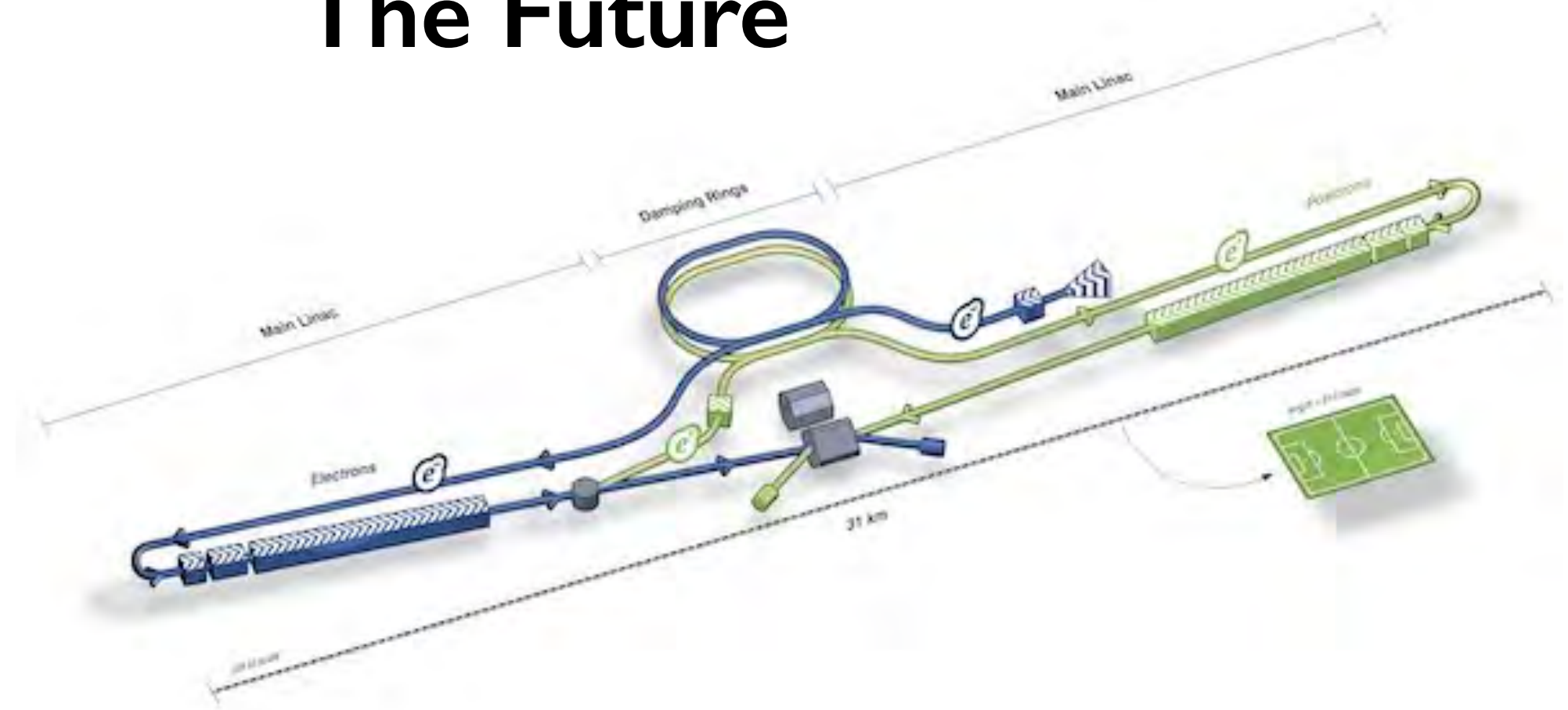
Global Fit to 2HDM of Type-2



- ▶ for given M_{H^\pm} tight constraints from H coupling measurements and EWPD
- ▶ expect improvement from direct searches at the LHC

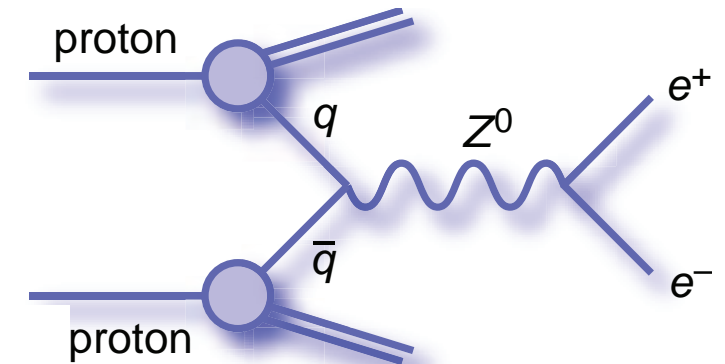


The Future

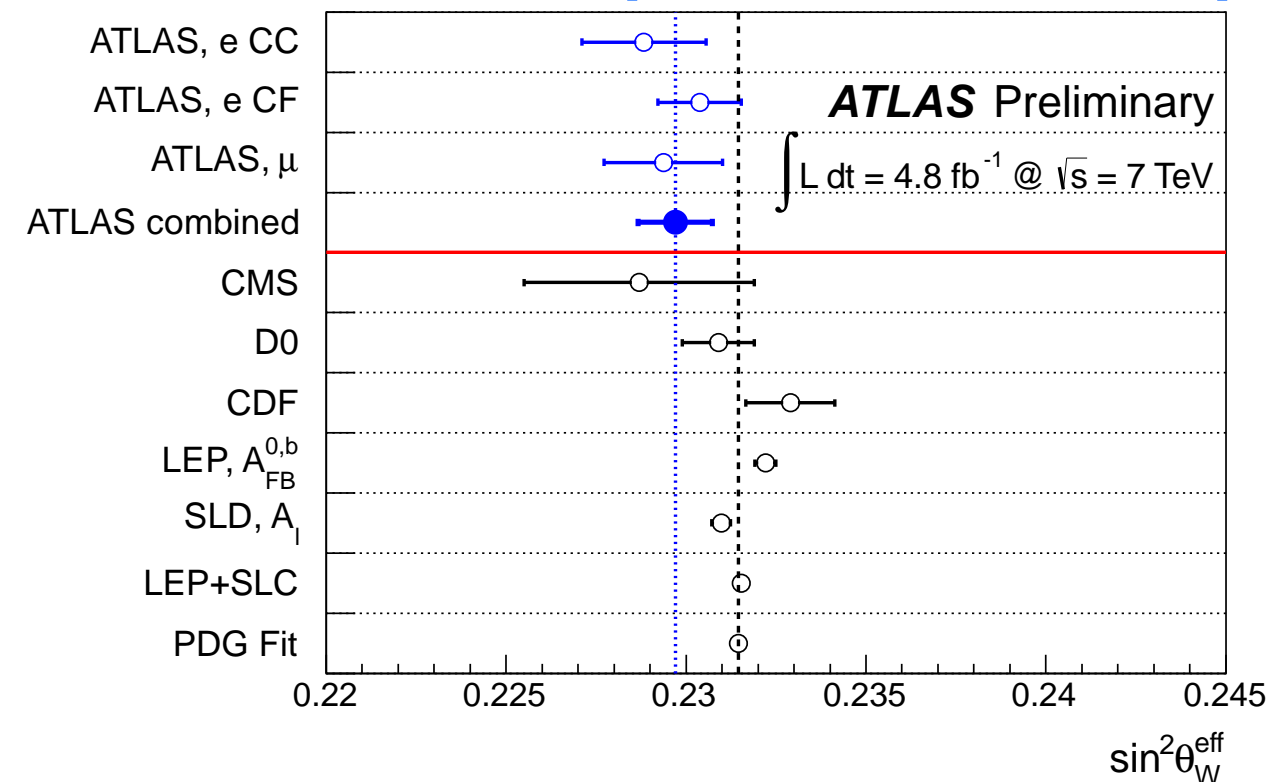


$\sin^2\theta_{\text{eff}}^l$ measurements at the LHC

- ▶ Drell-Yan: A_{FB} sensitive to distribution of polar angle of lepton w.r.t. *quark* direction
 - LHC: quark direction unknown!
- ▶ assume: dilepton boost is quark direction
 - often: interaction of valence quark with sea antiquark
 - important: reach in $|y_{\parallel}|$, ie. $|\eta_{\parallel}|$
- ▶ ambiguity due to PDFs dilution of A_{FB}
- ▶ $\sin^2\theta_{\text{eff}}^l$ from MC templates
 - accuracy of 9.8×10^{-4}
 - consistent with LEP/SLD result (accuracy 1.6×10^{-4})
- ▶ prediction for LHC 14/300
 - accuracy of 3.6×10^{-4} [[arXiv:1310.6708](https://arxiv.org/abs/1310.6708)]



[ATLAS-CONF-2013-043]



$$\sin^2\theta_{\text{eff}}^l(\text{exp}) = 0.23153 \pm 0.00016$$

$$\sin^2\theta_{\text{eff}}^l(\text{fit}) = 0.23149 \pm 0.00007$$

substantial contribution from LHC difficult

Future improvements

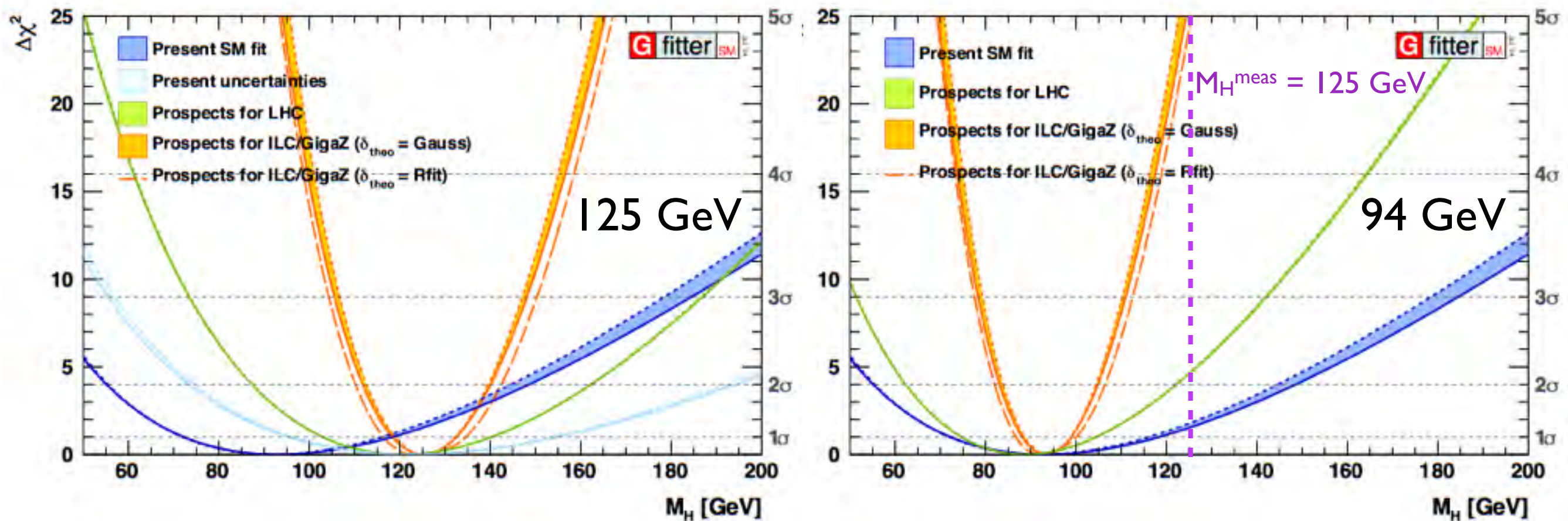
Parameter	Present	LHC	ILC/GigaZ	
M_H [GeV]	0.2	$\rightarrow < 0.1$	< 0.1	
M_W [MeV]	15	$\rightarrow 8$	$\rightarrow 5$	WW threshold
M_Z [MeV]	2.1	2.1	2.1	
m_t [GeV]	0.8	$\rightarrow 0.6$	$\rightarrow 0.1$	$t\bar{t}$ threshold scan
$\sin^2\theta_{\text{eff}}^\ell$ [10^{-5}]	16	16	$\rightarrow 1.3$	$\delta A^{0,f}_{LR}: 10^{-3} \rightarrow 10^{-4}$
$\Delta\alpha_{\text{had}}^5(M_Z^2)$ [10^{-5}]	10	$\rightarrow 4.7$	4.7	low energy data, better α_s
R_l^0 [10^{-3}]	25	25	$\rightarrow 4$	high statistics on Z-pole
κ_V ($\lambda = 3 \text{ TeV}$)	0.05	$\rightarrow 0.03$	$\rightarrow 0.01$	direct measurement of BRs

LHC = LHC with 300 fb^{-1}
 ILC/GigaZ = future e^+e^- collider, option to run on Z-pole (w polarized beams)

- ▶ theoretical uncertainties reduced by a factor of 4 (esp. M_W and $\sin^2\theta_{\text{eff}}^\ell$)
 - implies three-loop calculations!
 - exception: $\delta_{\text{theo}} m_t (\text{LHC}) = 0.25 \text{ GeV}$ (factor 2)
- ▶ central values of input measurements adjusted to $M_H = 125 \text{ GeV}$

[Baak et al, arXiv:1310.6708]

Higgs mass

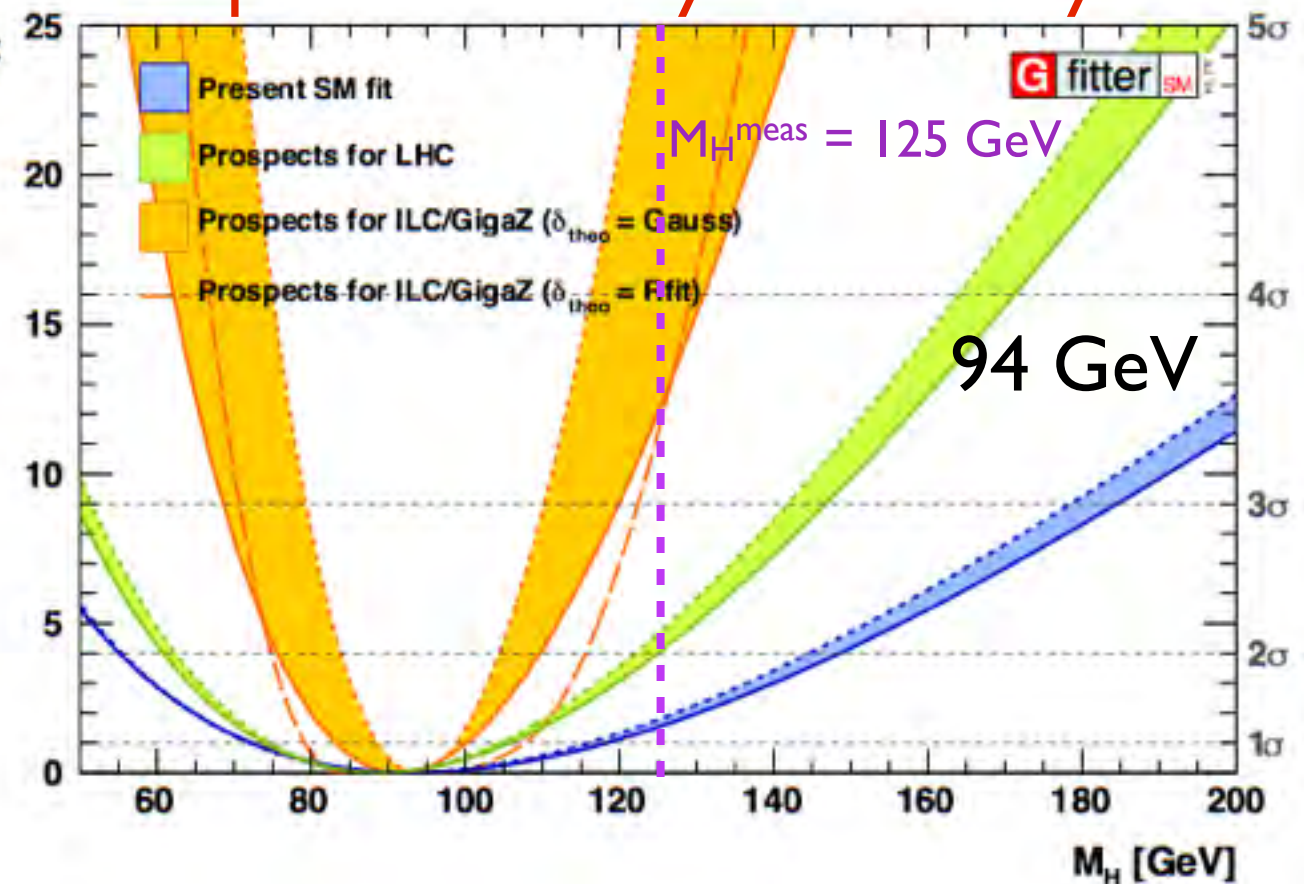
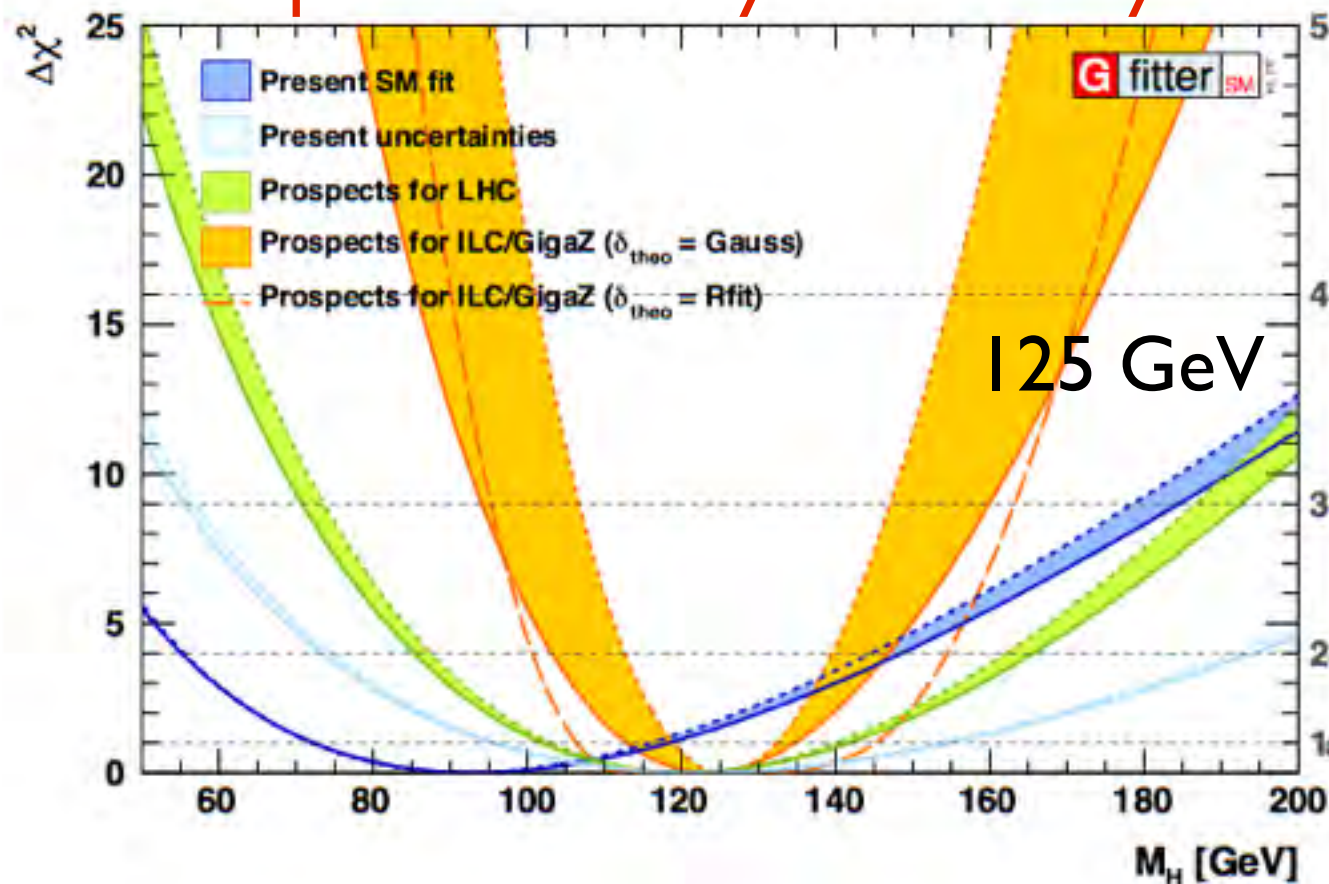


- ▶ Logarithmic dependency on $M_H \rightarrow$ cannot compete with direct M_H meas.
 - no theory uncertainty: $M_H = 125 \pm 7 \text{ GeV}$
 - future theory uncertainty (Rfit): $M_H = 125^{+10}_{-9} \text{ GeV}$
 - present day theory uncertainty: $M_H = 125^{+20}_{-17} \text{ GeV}$
- ▶ If EWPO central values unchanged (94 GeV), $\sim 5\sigma$ discrepancy with measured Higgs mass

Higgs mass

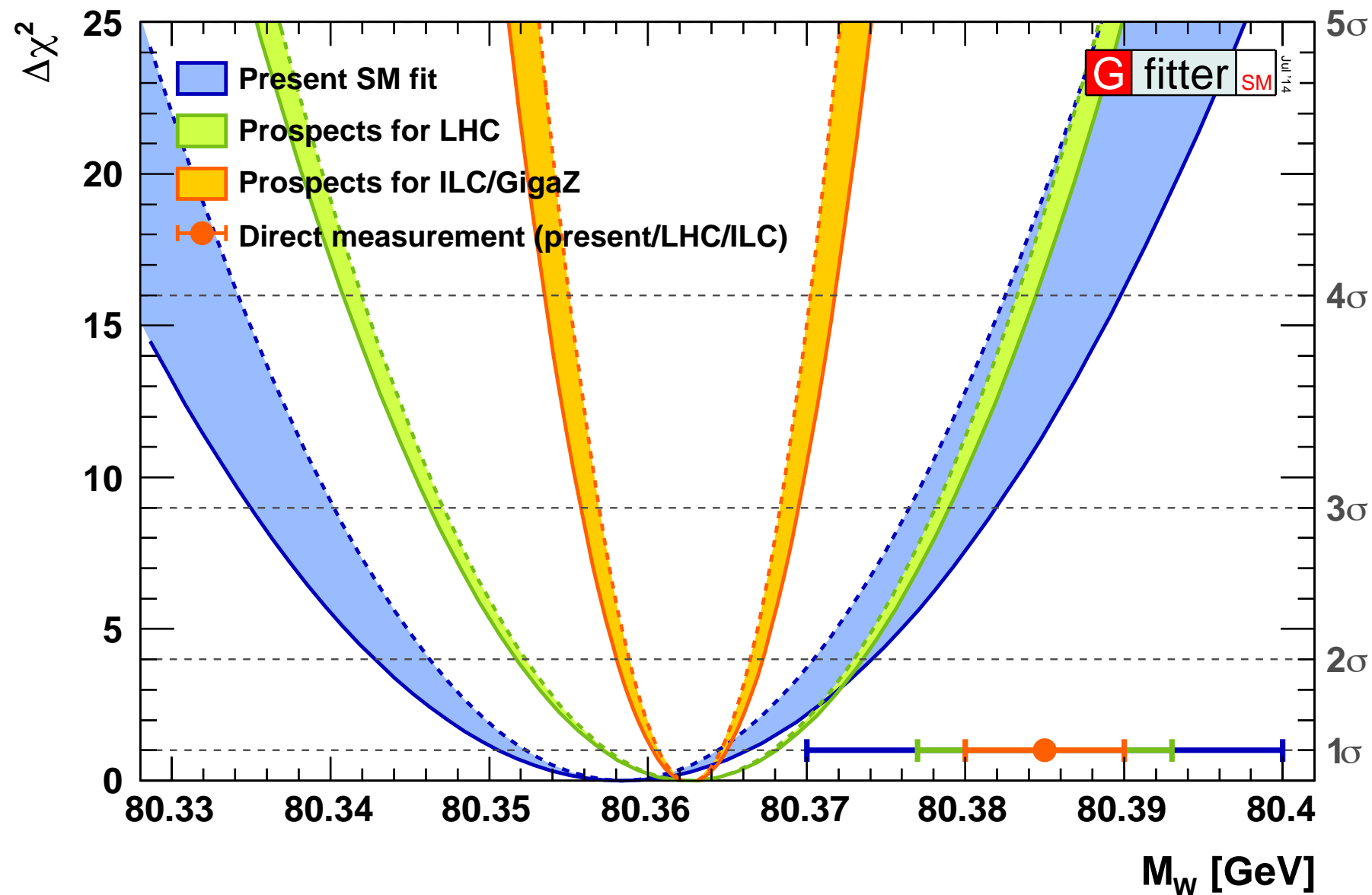
present theory uncertainty

present theory uncertainty



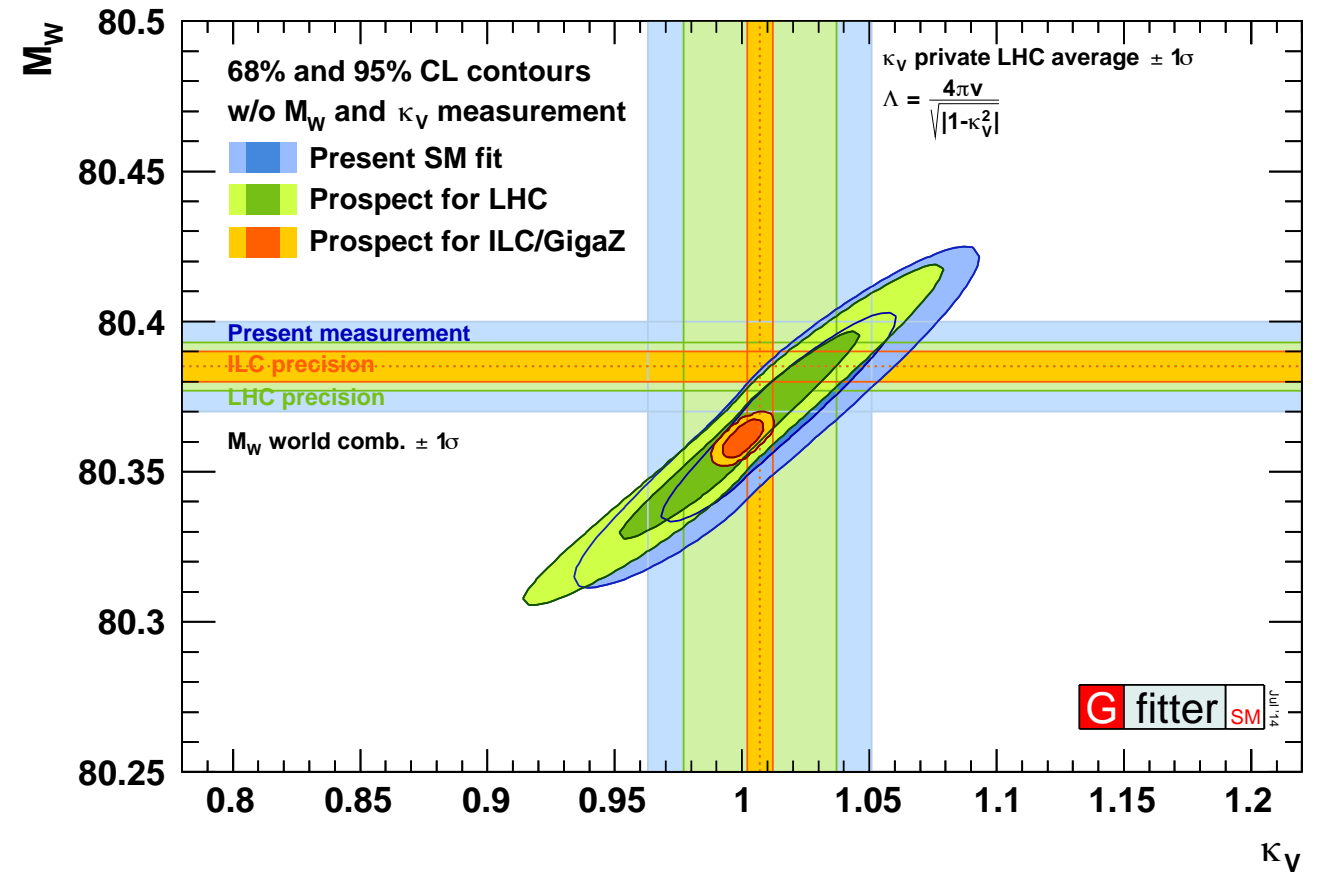
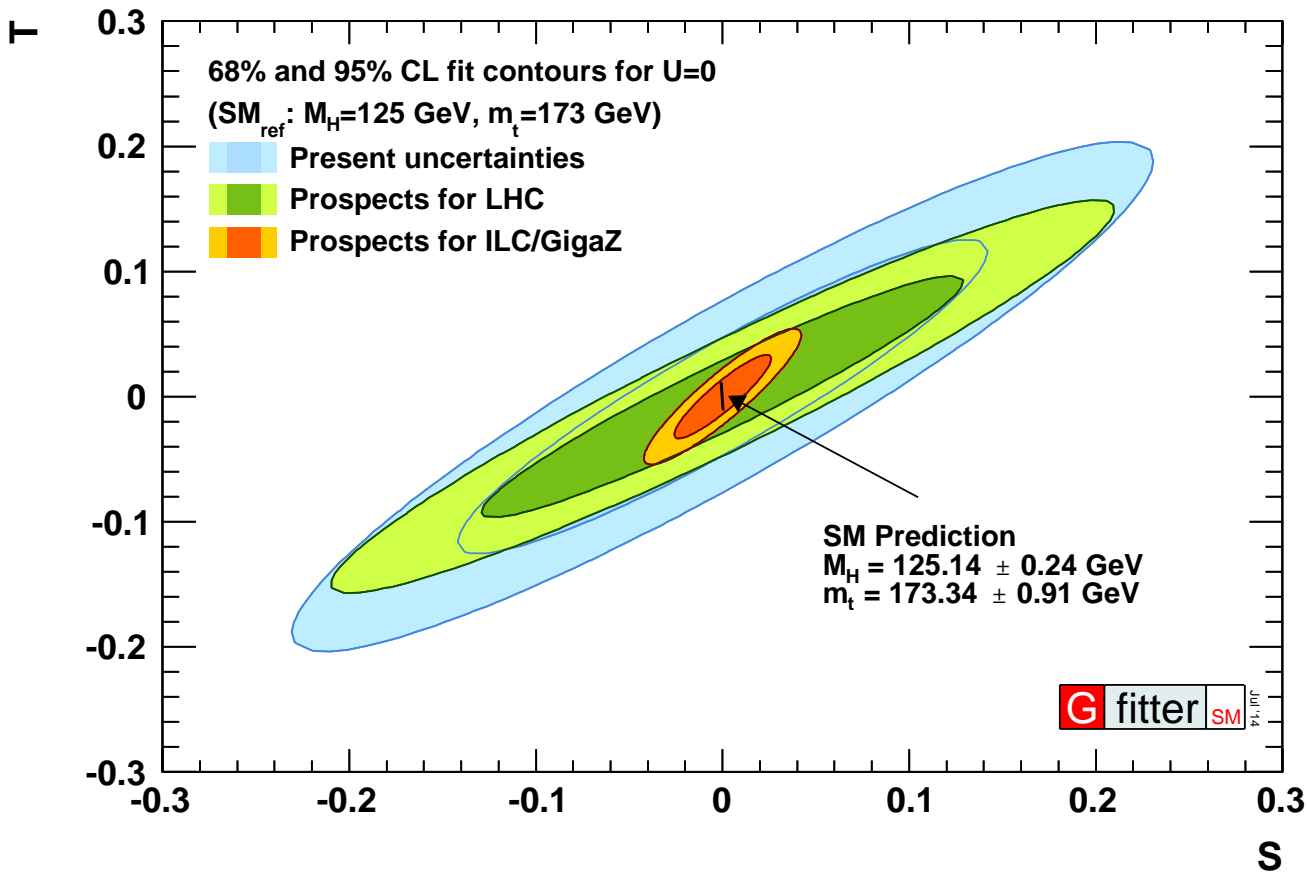
- ▶ Logarithmic dependency on $M_H \rightarrow$ cannot compete with direct M_H meas.
 - no theory uncertainty: $M_H = 125 \pm 7$ GeV
 - future theory uncertainty (Rfit): $M_H = 125^{+10}_{-9}$ GeV
 - present day theory uncertainty: $M_H = 125^{+20}_{-17}$ GeV
- ▶ If EWPO central values unchanged (94 GeV), $\sim 5\sigma$ discrepancy with measured Higgs mass **compromised by present theory uncertainty!**

Prospects for M_W



- ▶ improvement of a factor of 3 with the ILC (similar to measurement)
- ▶ stringent test of internal consistency of SM
- ▶ moderate improvement with LHC (~30%)
 - nevertheless, if at present values, theory uncertainties already important

BSM Prospects of EW fit



- ▶ for STU parameters, improvement of factor of >3 is possible at ILC
- ▶ again, at ILC a deviation between the SM predictions and direct measurements would be prominently visible.
- ▶ competitive results between EW fit and Higgs coupling measurements!
 - precision of about 1%

Summary of indirect predictions

Parameter	Experimental input [$\pm 1\sigma_{\text{exp}}$]				Indirect determination [$\pm 1\sigma_{\text{exp}}, \pm 1\sigma_{\text{theo}}$]		
	Present	LHC	ILC/GigaZ		Present	LHC	ILC/GigaZ
M_H [GeV]	0.2	< 0.1	< 0.1		$^{+31}_{-26}, ^{+10}_{-8}$	$^{+20}_{-18}, ^{+3.9}_{-3.2}$	$^{+6.8}_{-6.5}, ^{+2.5}_{-2.4}$
M_W [MeV]	15	8	5		6.0, 5.0	5.2, 1.8	1.9, 1.3
M_Z [MeV]	2.1	2.1	2.1		11, 4	7.0, 1.4	2.5, 1.0
m_t [GeV]	0.8	0.6	0.1		2.4, 0.6	1.5, 0.2	0.7, 0.2
$\sin^2\theta_{\text{eff}}^\ell$ [10^{-5}]	16	16	1.3		4.5, 4.9	2.8, 1.1	2.0, 1.0
$\Delta\alpha_{\text{had}}^5(M_Z^2)$ [10^{-5}]	10	4.7	4.7		42, 13	36, 6	5.6, 3.0
R_l^0 [10^{-3}]	25	25	4		—	—	—
$\alpha_s(M_Z^2)$ [10^{-4}]	—	—	—		40, 10	39, 7	6.4, 6.9
$S _{U=0}$	—	—	—		0.094, 0.027	0.086, 0.006	0.017, 0.006
$T _{U=0}$	—	—	—		0.083, 0.023	0.064, 0.005	0.022, 0.005
κ_V ($\lambda = 3 \text{ TeV}$)	0.05	0.03	0.01		0.02	0.02	0.01

Summary of indirect predictions

Parameter	Experimental input [$\pm 1\sigma_{\text{exp}}$]				Indirect determination [$\pm 1\sigma_{\text{exp}}, \pm 1\sigma_{\text{theo}}$]		
	Present	LHC	ILC/GigaZ		Present	LHC	ILC/GigaZ
M_H [GeV]	0.2	< 0.1	< 0.1		$^{+31}_{-26}, ^{+10}_{-8}$	$^{+20}_{-18}, ^{+3.9}_{-3.2}$	$^{+6.8}_{-6.5}, ^{+2.5}_{-2.4}$
M_W [MeV]	15	8	5		6.0, 5.0	5.2, 1.8	1.9, 1.3
M_Z [MeV]	2.1	2.1	2.1		11, 4	7.0, 1.4	2.5, 1.0
m_t [GeV]	0.8	0.6	0.1		2.4, 0.6	1.5, 0.2	0.7, 0.2
$\sin^2\theta_{\text{eff}}^\ell$ [10^{-5}]	16	16	1.3		4.5, 4.9	2.8, 1.1	2.0, 1.0
$\Delta\alpha_{\text{had}}^5(M_Z^2)$ [10^{-5}]	10	4.7	4.7		42, 13	36, 6	5.6, 3.0
R_l^0 [10^{-3}]	25	25	4		—	—	—
$\alpha_s(M_Z^2)$ [10^{-4}]	—	—	—		40, 10	39, 7	6.4, 6.9
$S _{U=0}$	—	—	—		0.094, 0.027	0.086, 0.006	0.017, 0.006
$T _{U=0}$	—	—	—		0.083, 0.023	0.064, 0.005	0.022, 0.005
κ_V ($\lambda = 3 \text{ TeV}$)	0.05	0.03	0.01		0.02	0.02	0.01

- theory uncertainty needs to be reduced if we want to achieve the ultimate precision with the LHC!
- ILC/GigaZ offers fantastic possibilities to test the SM and constrain NP

M_W : Impact of Uncertainties

Today

$$\delta_{\text{meas}} = 15 \text{ MeV}$$

$$\delta_{\text{fit}} = 8 \text{ MeV}$$

LHC-300

$$\delta_{\text{meas}} = 8 \text{ MeV}$$

$$\delta_{\text{fit}} = 6 \text{ MeV}$$

ILC/GigaZ

$$\delta_{\text{meas}} = 5 \text{ MeV}$$

$$\delta_{\text{fit}} = 2 \text{ MeV}$$

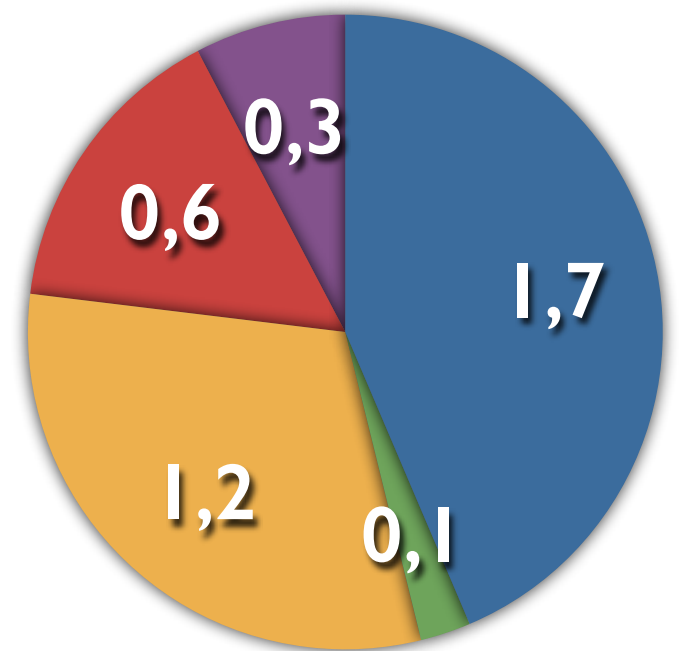
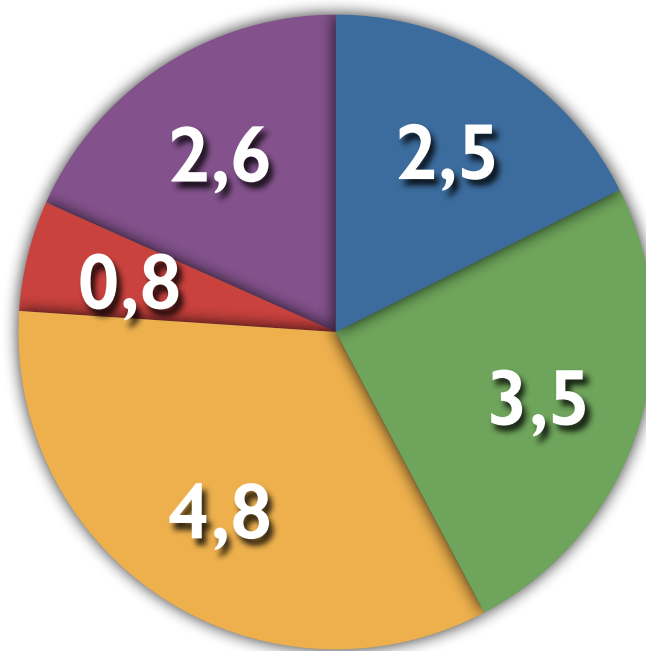
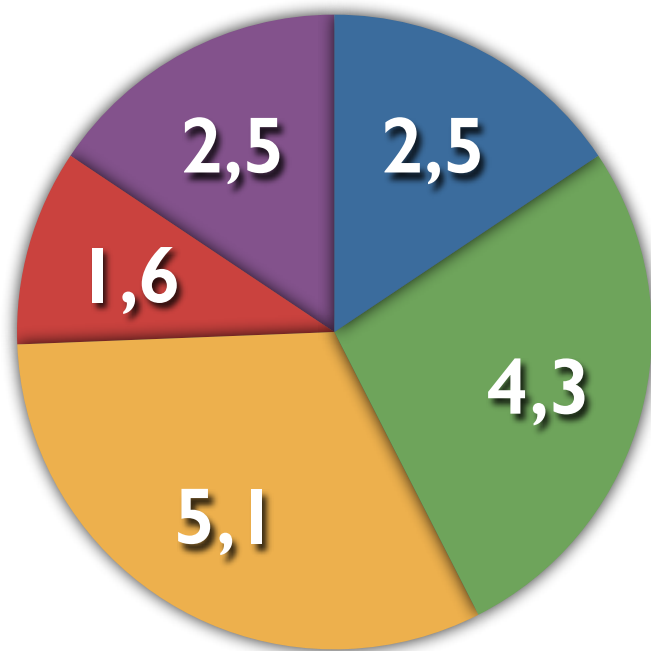
● δM_Z

● δm_{top}

● $\delta \sin^2(\theta_{\text{eff}}^l)$

● $\delta \Delta \alpha_{\text{had}}$

● $\delta \alpha_s$



Impact of individual uncertainties on δM_W in fit (numbers in MeV)

► ILC/GigaZ: impact δM_Z of will become important again!

Summary

Paradigm change

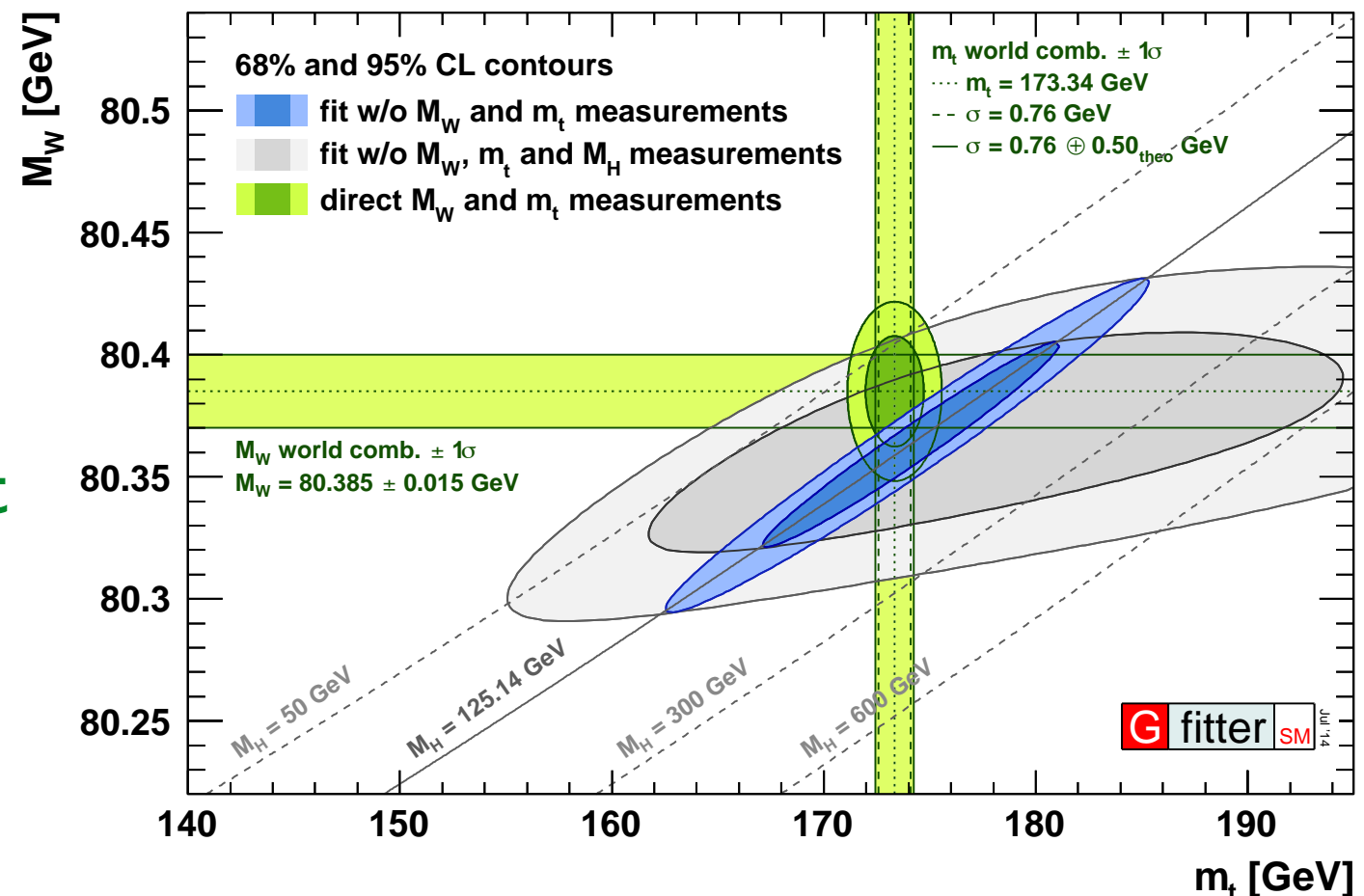
- ▶ from the discovery of the Higgs boson to a **probe of new physics**
- ▶ knowledge of M_H and two-loop calculations
unprecedented precision of EW fit
- ▶ cannot know M_W and $\sin^2\theta_{\text{eff}}^l$ precise enough

LHC 14/300

- ▶ ΔM_W (indirect) = 5.5 MeV
 ΔM_W (exp) = 8 MeV

ILC with GigaZ

- ▶ Δm_t (exp) = 100 MeV \rightarrow ΔM_W (indirect) = 2 MeV
measurement of M_Z will become important again ($\Delta\alpha_{\text{had}}$ as well)
- ▶ indirect determinations of M_Z and $\Delta\alpha_{\text{had}}$ will match exp. precision



More information and latest results:

www.cern.ch/gfitter

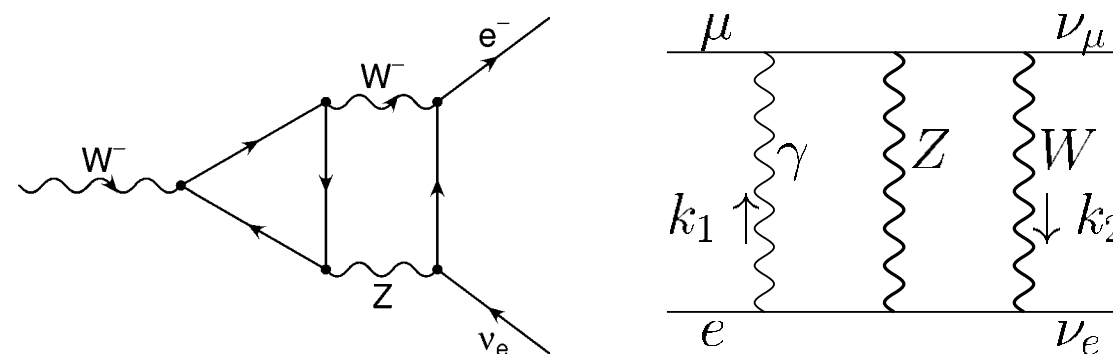
Additional Material

Calculation of M_W

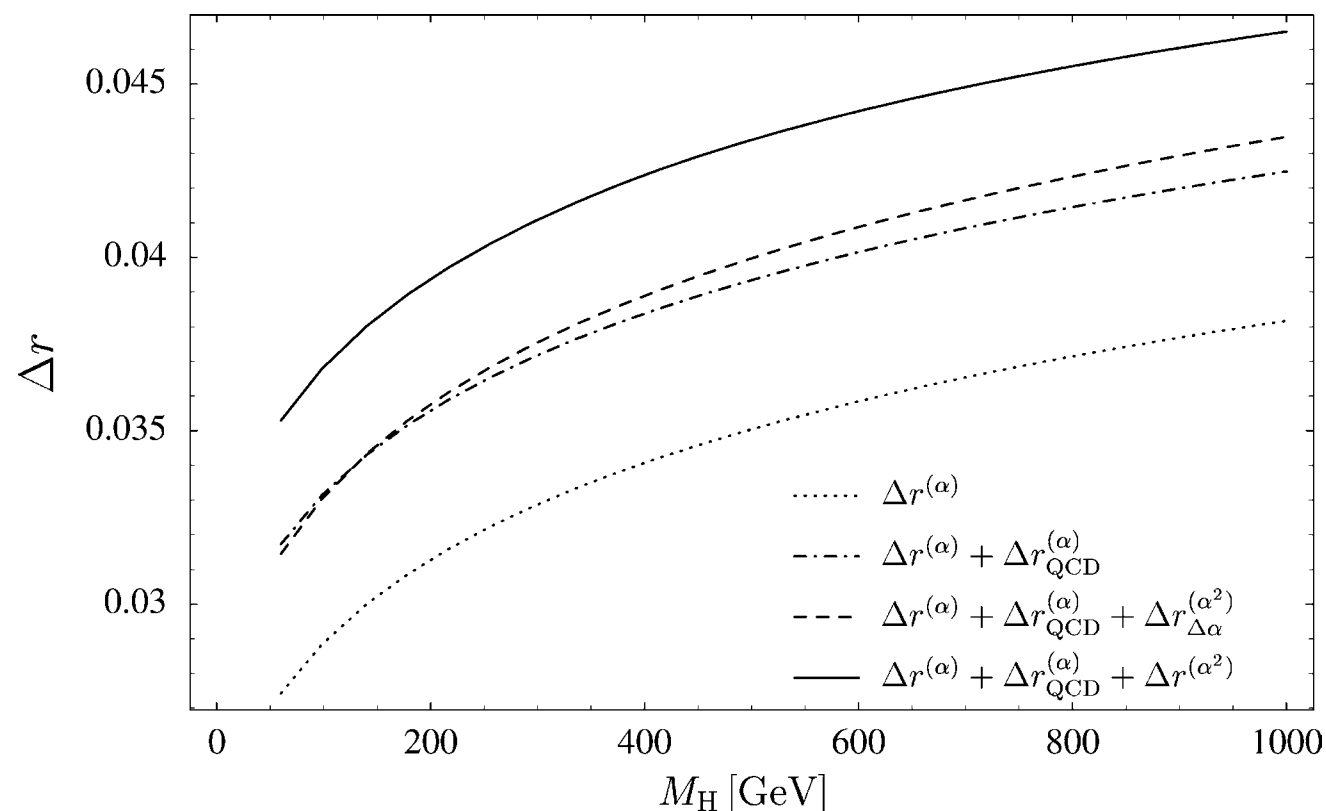
- ▶ Full **EW** one- and two-loop calculation of fermionic and bosonic contributions
- ▶ One- and two-loop **QCD** corrections and leading terms of higher order corrections
- ▶ **Results** for Δr include terms of order $O(\alpha)$, $O(\alpha\alpha_s)$, $O(\alpha\alpha_s^2)$, $O(\alpha^2_{\text{ferm}})$, $O(\alpha^2_{\text{bos}})$, $O(\alpha^2\alpha_s m_t^4)$, $O(\alpha^3 m_t^6)$
- ▶ Uncertainty estimate:
 - missing terms of order $O(\alpha^2\alpha_s)$: about 3 MeV (from $O(\alpha^2\alpha_s m_t^4)$)
 - electroweak three-loop correction $O(\alpha^3)$: < 2 MeV
 - three-loop QCD corrections $O(\alpha\alpha_s^3)$: < 2 MeV
 - **Total: $\delta M_W \approx 4 \text{ MeV}$**

[M Awramik et al., Phys. Rev. D69, 053006 (2004)]

[M Awramik et al., Phys. Rev. Lett. 89, 241801 (2002)]



A Freitas et al., Phys. Lett. B495, 338 (2000)]



Calculation of $\sin^2(\theta_{\text{eff}}^l)$

- Effective mixing angle:

$$\sin^2 \theta_{\text{eff}}^{\text{lept}} = (1 - M_W^2/M_Z^2) (1 + \Delta\kappa)$$

- Two-loop EW and QCD correction to $\Delta\kappa$ known, leading terms of higher order QCD corrections

- fermionic two-loop correction about 10^{-3} , whereas bosonic one 10^{-5}

- **Uncertainty** estimate obtained with different methods, geometric progression:

$$\mathcal{O}(\alpha^2\alpha_s) = \frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha\alpha_s).$$

$$\mathcal{O}(\alpha^2\alpha_s) \text{ beyond leading } m_t^4 \quad 3.3 \dots 2.8 \times 10^{-5}$$

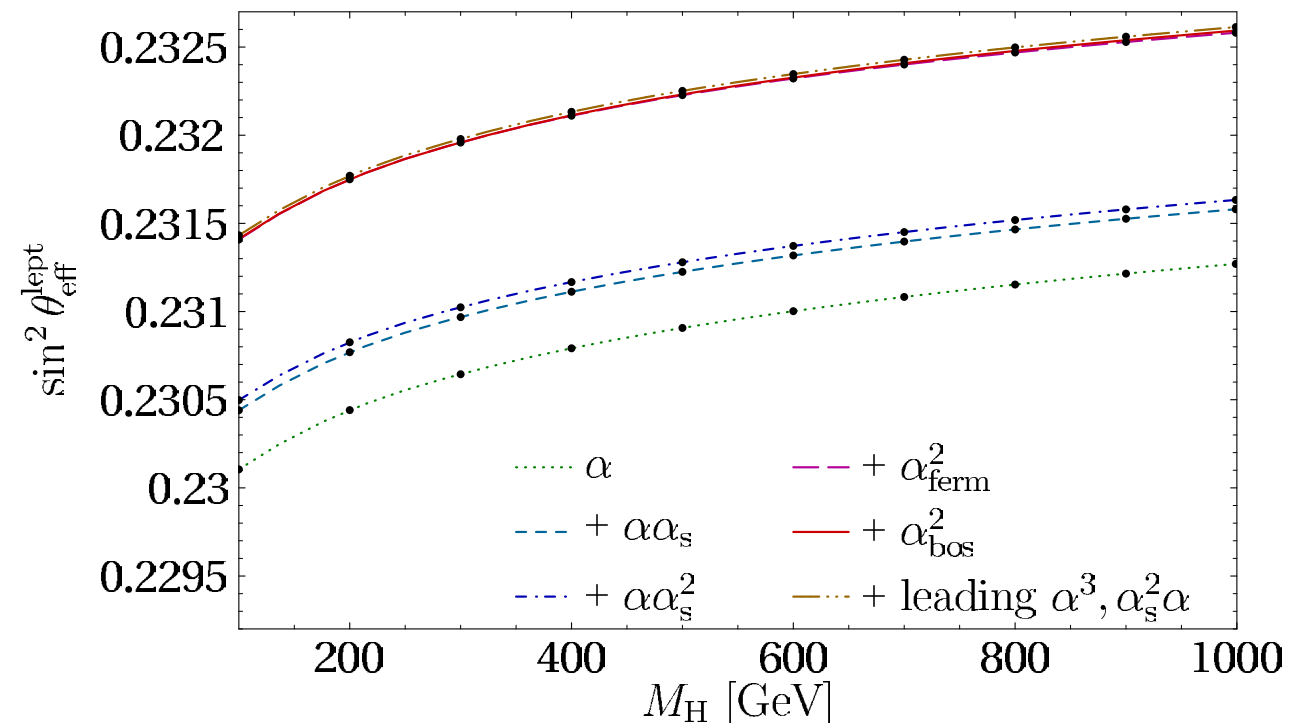
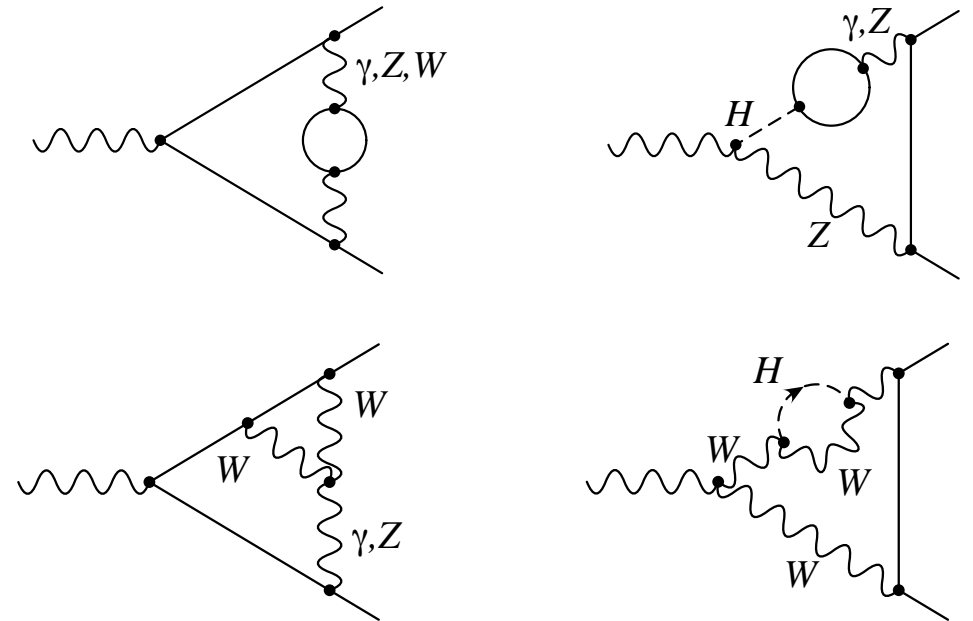
$$\mathcal{O}(\alpha\alpha_s^3) \quad 1.5 \dots 1.4$$

$$\mathcal{O}(\alpha^3) \text{ beyond leading } m_t^6 \quad 2.5 \dots 3.5$$

$$\text{Total: } \delta\sin^2\theta_{\text{eff}}^l \approx 4.7 \cdot 10^{-5}$$

[M Awramik et al, Phys. Rev. Lett. 93, 201805 (2004)]

[M Awramik et al., JHEP 11, 048 (2006)]

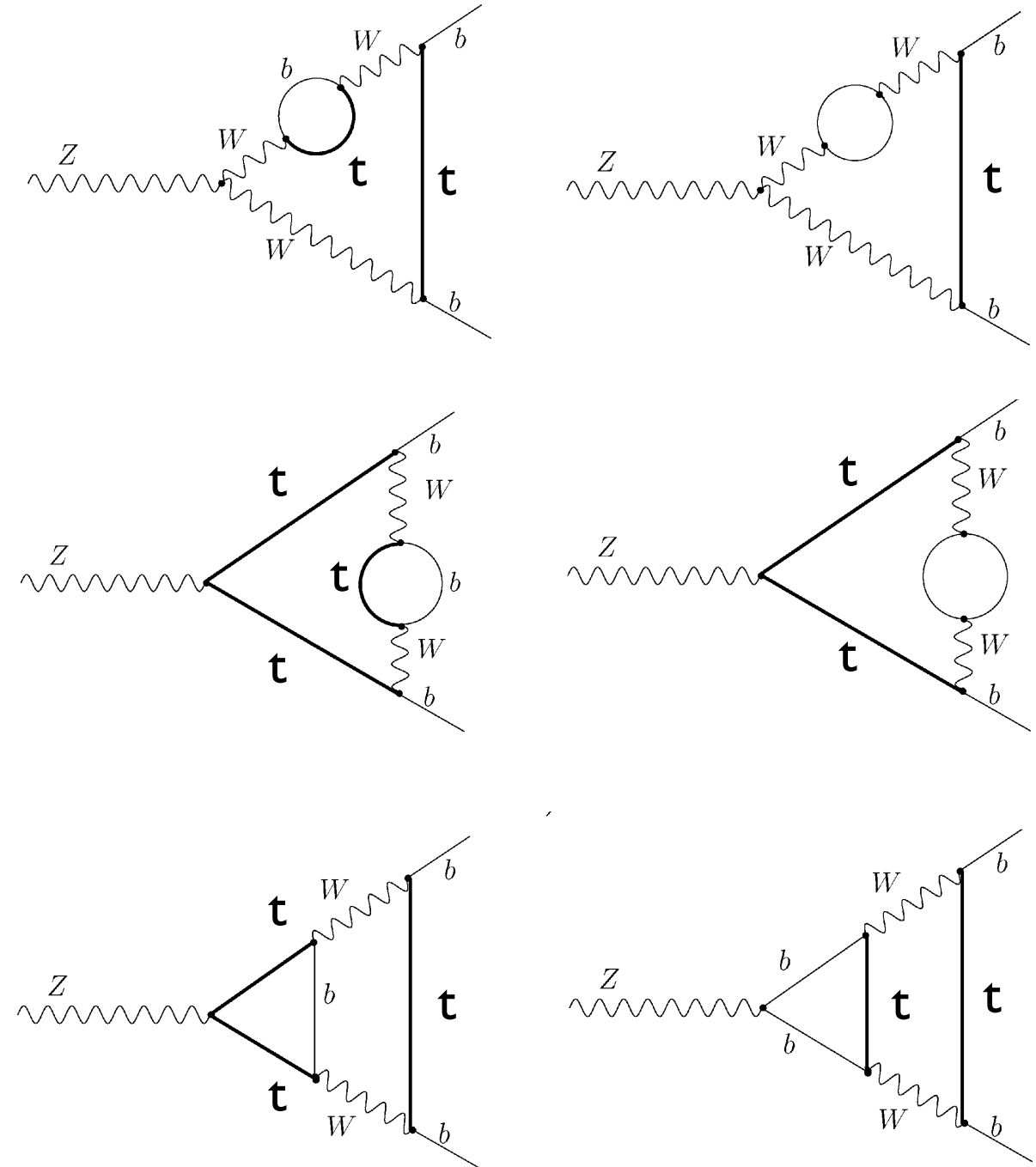


Calculation of $\sin^2(\theta_{\text{eff}}^{bb})$

[M Awramik et al, Nucl. Phys. B813, 174 (2009)]

- ▶ Calculation of $\sin^2\theta_{\text{eff}}$ for **b-quarks** more involved, because of top quark propagators in the $Z \rightarrow b\bar{b}$ vertex
- ▶ Investigation of known discrepancy between $\sin^2\theta_{\text{eff}}$ from leptonic and hadronic asymmetry measurements
- ▶ Two-loop **EW** correction only recently completed, effect of $O(10^{-4})$
- ▶ Now $\sin^2\theta_{\text{eff}}^{bb}$ known at the same order as $\sin^2\theta_{\text{eff}}$ for leptons and light quarks
- ▶ Uncertainty assumed to be of same size as for $\sin^2\theta_{\text{eff}}$:

$$\delta\sin^2\theta_{\text{eff}}^{bb} \approx 4.7 \cdot 10^{-5}$$



Calculation of R_b^0

Full two-loop calculation of $Z \rightarrow b\bar{b}$

— [A. Freitas et al., JHEP 1208, 050 (2012)
Erratum ibid. 1305 (2013) 074]

- ▶ The branching ratio R_b^0 : partial decay width of $Z \rightarrow b\bar{b}$ and $Z \rightarrow q\bar{q}$

$$R_b \equiv \frac{\Gamma_b}{\Gamma_{\text{had}}} = \frac{\Gamma_b}{\Gamma_d + \Gamma_u + \Gamma_s + \Gamma_c + \Gamma_b} = \frac{1}{1 + 2(\Gamma_d + \Gamma_u)/\Gamma_b}$$

- ▶ Contribution of same terms as in the calculation of $\sin^2\theta_{\text{eff}}^{bb}$
→ cross-check the two results, found good agreement
- ▶ Two-loop corrections small compared to experimental uncertainty ($6.6 \cdot 10^{-4}$)

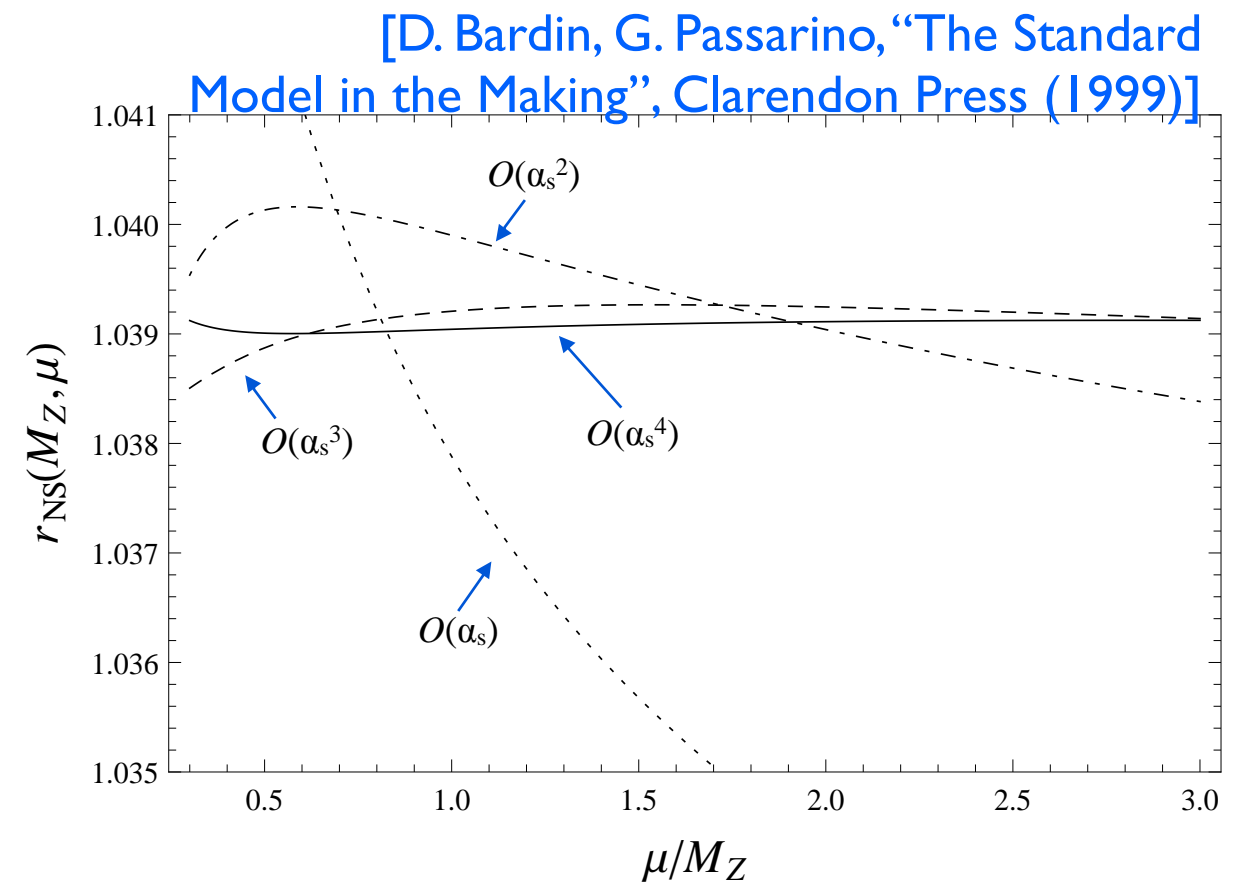
	I-loop EW and QCD correction to FSR	2-loop EW correction	2-loop EW and 2+3-loop QCD correction to FSR	1+2-loop QCD correction to gauge boson selfenergies
M_H [GeV]	$\mathcal{O}(\alpha) + \text{FSR}_{\alpha, \alpha_s, \alpha_s^2}$ [10^{-4}]	$\mathcal{O}(\alpha_{\text{ferm}}^2)$ [10^{-4}]	$\mathcal{O}(\alpha_{\text{ferm}}^2) + \text{FSR}_{\alpha_s^3, \alpha\alpha_s, m_b^2\alpha_s, m_b^4}$ [10^{-4}]	$\mathcal{O}(\alpha\alpha_s, \alpha\alpha_s^2)$ [10^{-4}]
100	−35.66	−0.856	−2.496	−0.407
200	−35.85	−0.851	−2.488	−0.407
400	−36.09	−0.846	−2.479	−0.406

Radiator Functions

- ▶ Partial widths are defined inclusively: they contain QCD and QED contributions
- ▶ Corrections can be expressed as radiator functions $R_{A,f}$ and $R_{V,f}$

$$\Gamma_{f\bar{f}} = N_c^f \frac{G_F M_Z^3}{6\sqrt{2}\pi} \left(|g_{A,f}|^2 R_{A,f} + |g_{V,f}|^2 R_{V,f} \right)^2$$

- ▶ High sensitivity to the strong coupling α_s
- ▶ Full four-loop calculation of QCD Adler function available (**N³LO**)
- ▶ Much reduced scale dependence
- ▶ Theoretical uncertainty of 0.1 MeV, compare to experimental uncertainty of 2.0 MeV



[P. Baikov et al., Phys. Rev. Lett. 108, 222003 (2012)]
 [P. Baikov et al Phys. Rev. Lett. 104, 132004 (2010)]

Modified Higgs Couplings

Study of potential deviations of Higgs couplings from SM

- ▶ BSM modelled as extension of SM through effective Lagrangian
 - Leading corrections only
- ▶ Benchmark model:
 - Scaling of Higgs-vector boson (κ_V) and Higgs-fermion couplings (κ_F)
 - **No additional loops** in the production or decay of the Higgs, **no invisible Higgs decays and undetectable width**
- ▶ Main effect on EWPO due to modified Higgs coupling to gauge bosons (κ_V)
 - Involving the longitudinal d.o.f.
- ▶ Most BSM models: $\kappa_V < 1$
- ▶ Additional Higgses typically give positive contribution to M_W

$$L_V = \frac{h}{v} \left(2\kappa_V m_W^2 W_\mu W^\mu + \kappa_V m_Z^2 Z_\mu Z^\mu \right)$$

$$L_F = -\frac{h}{v} \left(\kappa_F m_t \bar{t}t + \kappa_F m_b \bar{b}b + \kappa_F m_\tau \bar{\tau}\tau \right)$$

