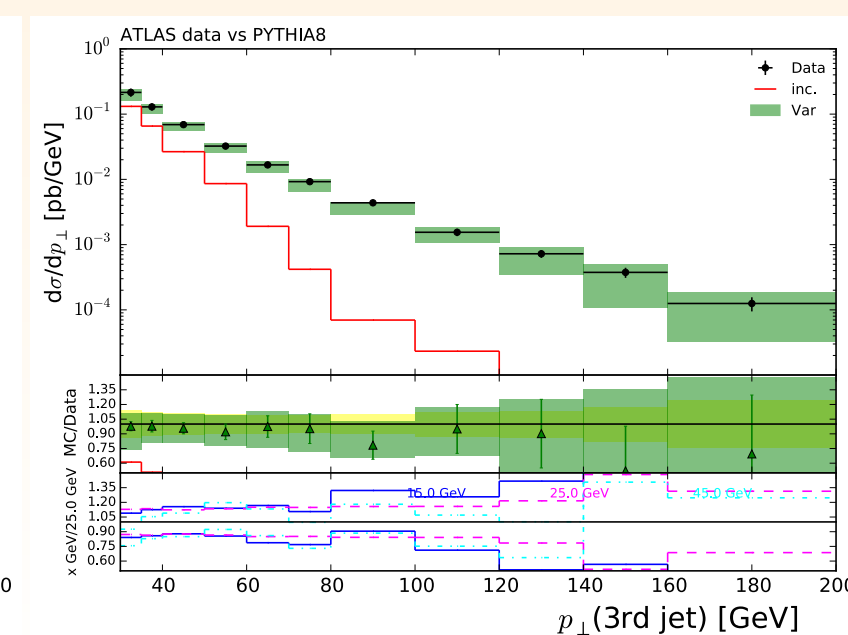
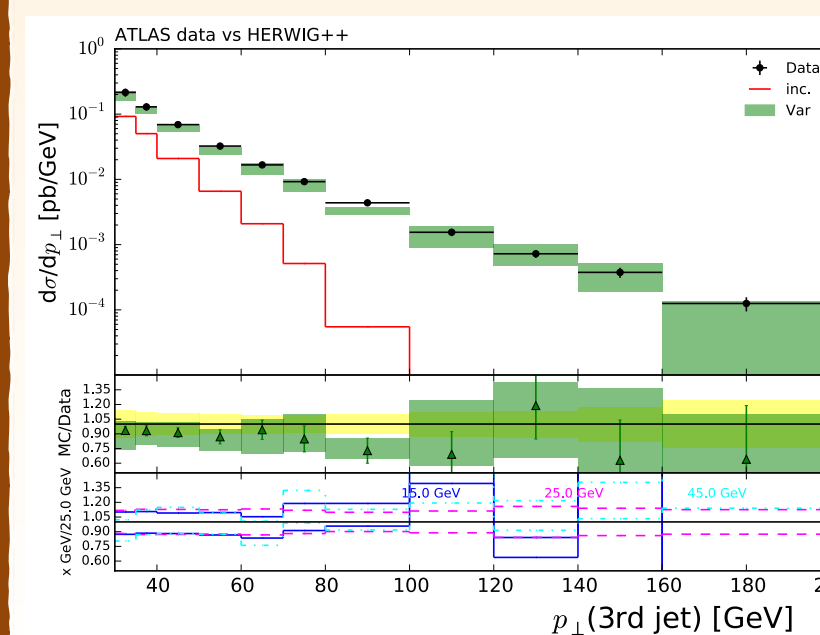
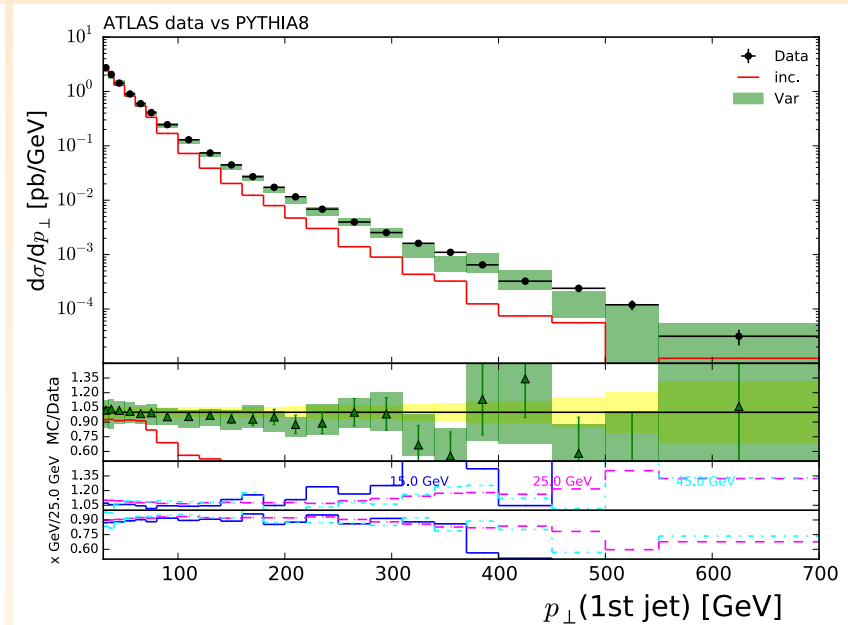
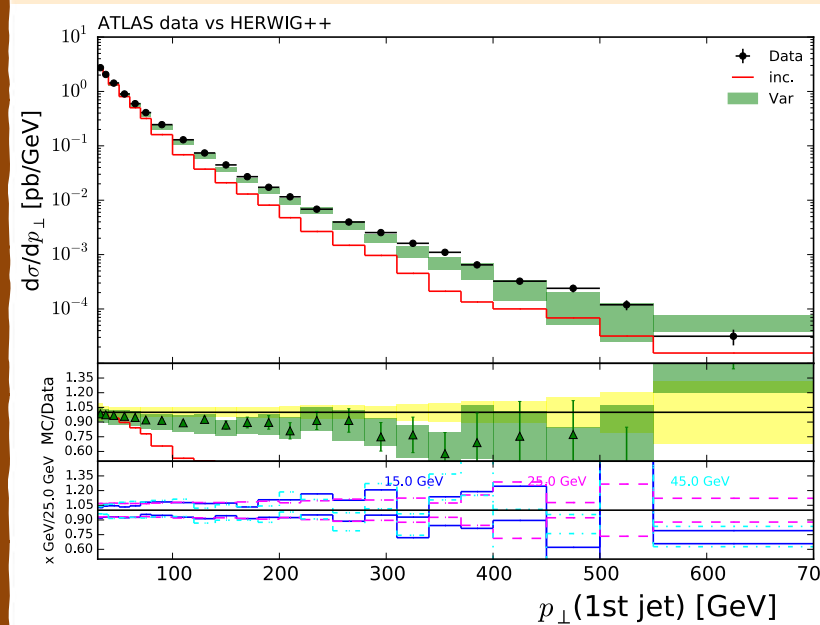
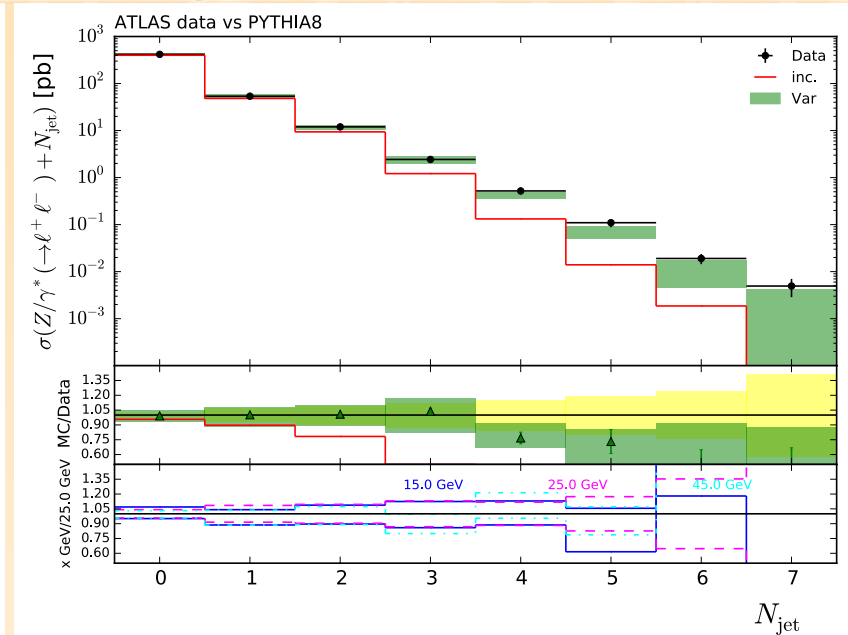
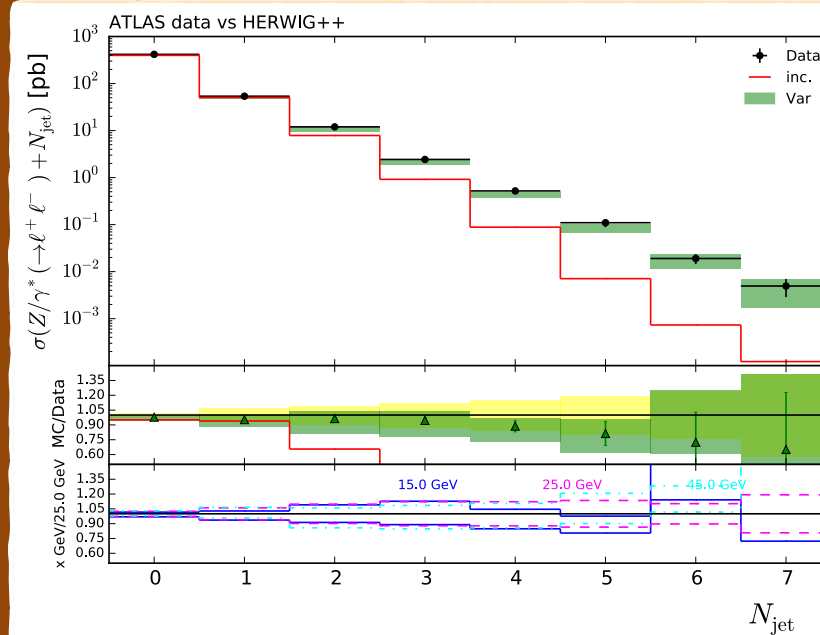
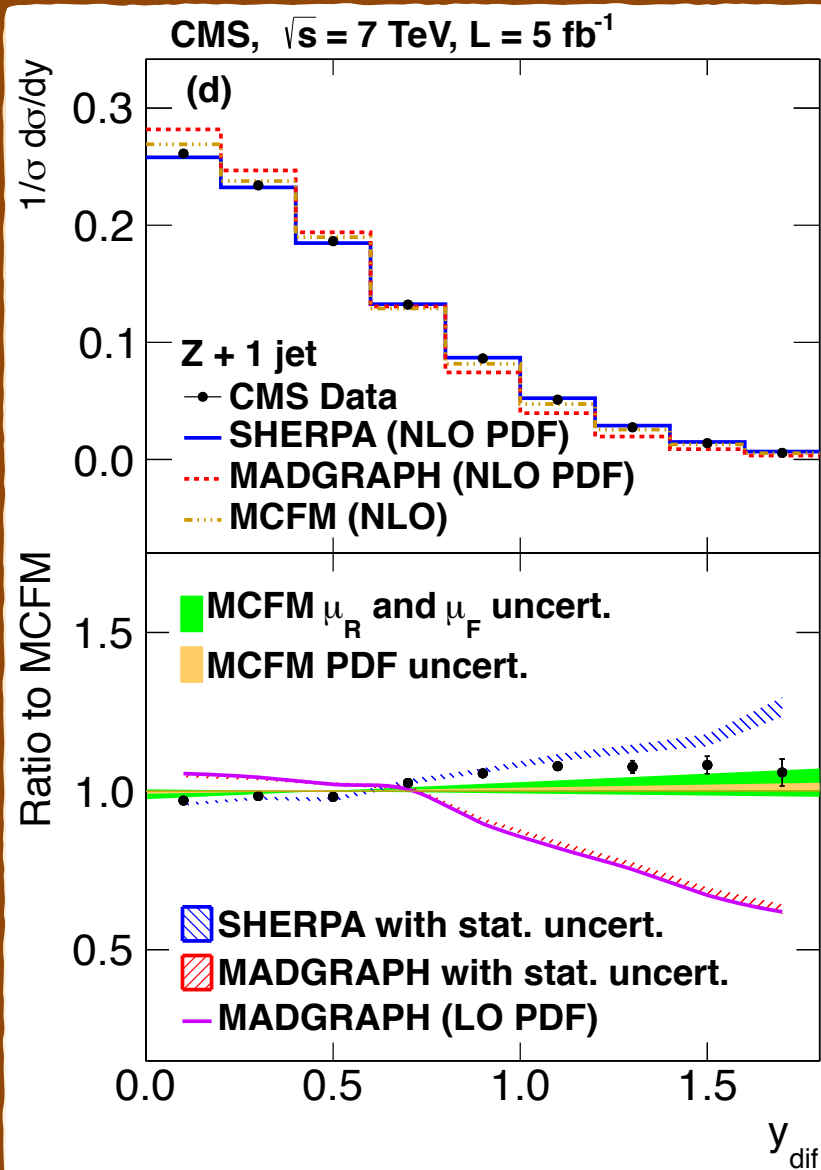
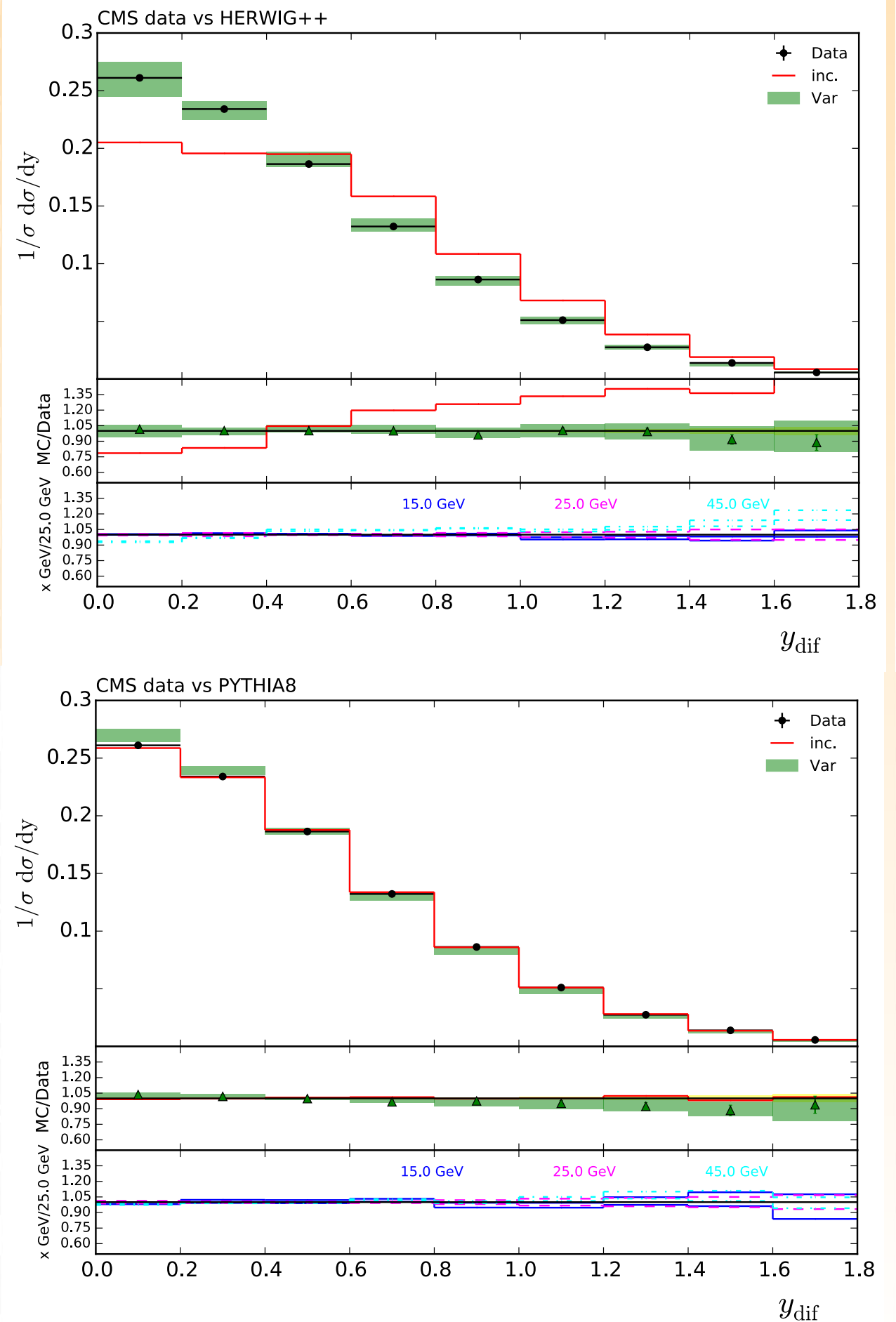


- ◆ Z+jets
- ◆ Exclusive jet multiplicity and hardest and 3rd hardest jet pT spectra
- ◆ Uncertainty band contains ren. & fac. scale, PDF & merging scale dependence
- ◆ Rather good agreement between data and theory



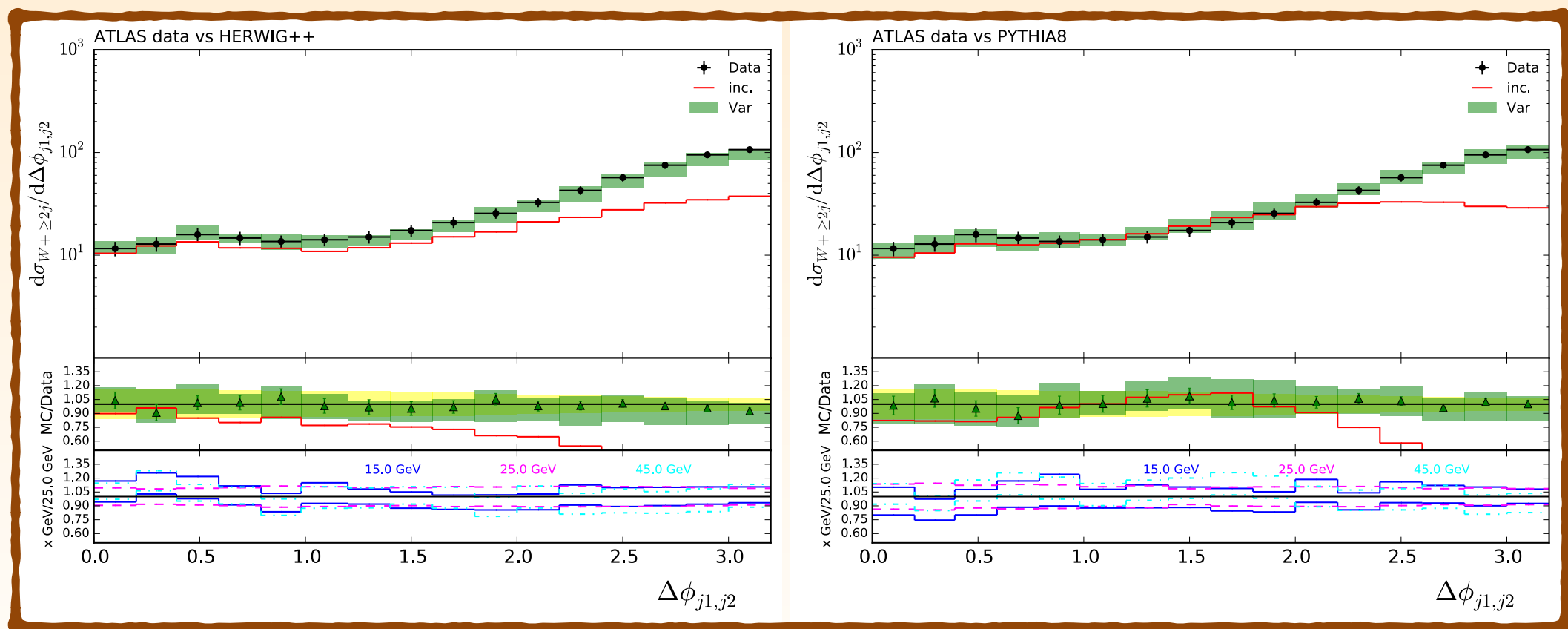


- ♦ Rapidity difference between Z-boson and hardest jet.
- ♦ Sensitive to higher multiplicity matrix elements
- ♦ LO predictions off (in particular MadGraph)
- ♦ No discrepancies at NLO

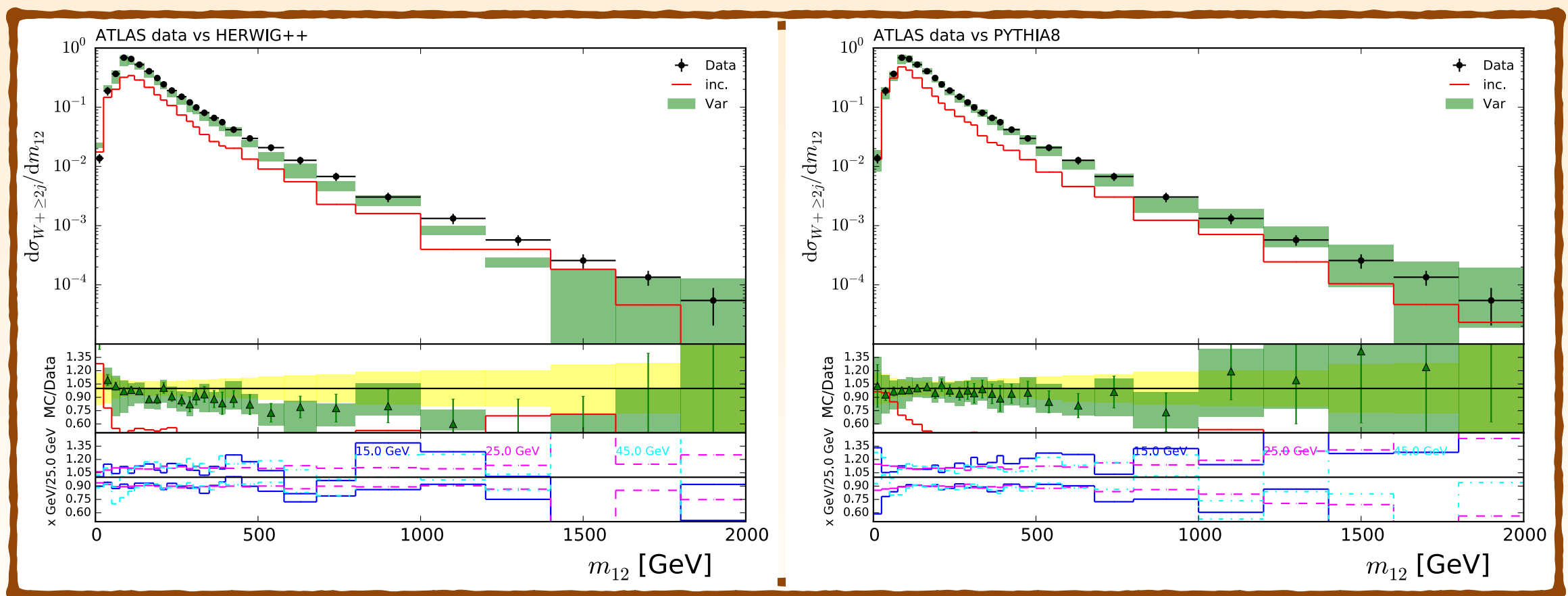


- ♦ $W_{\text{+jets}}$
- ♦ Agreement between FxFx merged results, matched to Herwig++ and Pythia8, and Atlas and CMS data is rather good
- ♦ Where data and theory differ, also differences between the results matched to HW++ and PY8 differ

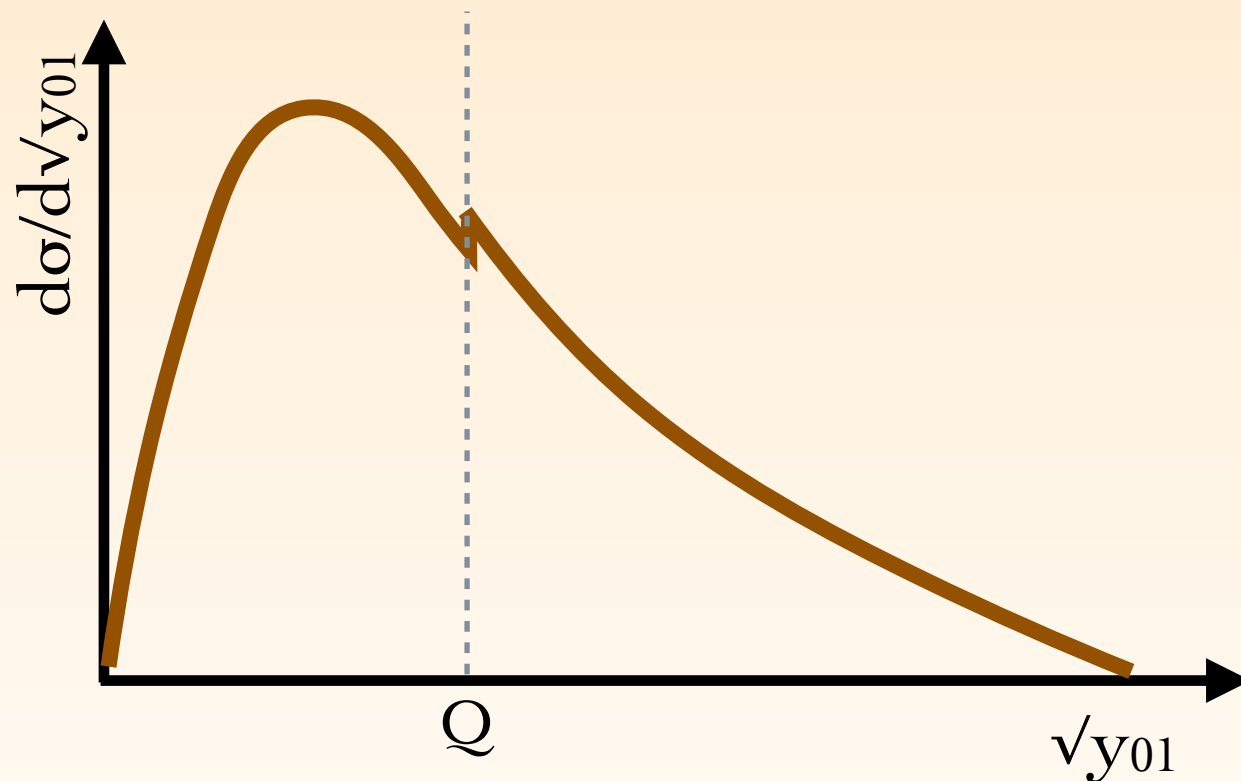
- ◆ W +jets
- ◆ Agreement between FxFx merged results, matched to Herwig++ and Pythia8, and Atlas and CMS data is rather good
- ◆ Where data and theory differ, also differences between the results matched to HW++ and PY8 differ



- ♦ W +jets
- ♦ Agreement between FxFx merged results, matched to Herwig++ and Pythia8, and Atlas and CMS data is rather good
- ♦ Where data and theory differ, also differences between the results matched to HW++ and PY8 differ



FXFX / MEPS@NLO: V & V+1J MERGING



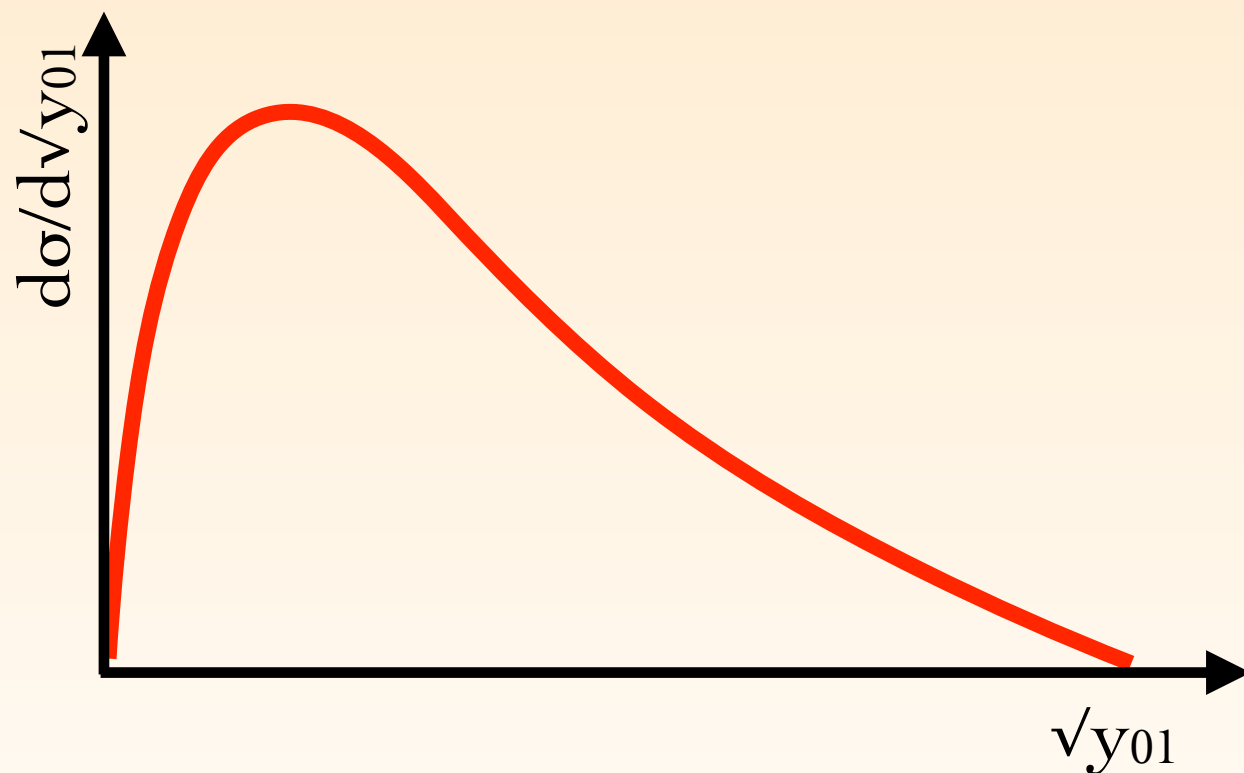
FXFX: [RF, Frixione (2012)]

MEPS@NLO: [Hoeche, Krauss, Schonherr, Siegert; +Gehrmann (2012)]

- ♦ Merge NLO+PS for V with Minlo for V+1j, at “merging scale” Q
- ♦ Above Q the tail is NLO accurate
- ♦ For not-too-small Q , integral is NLO accurate
- ♦ Used by ATLAS & CMS for LHC run II analyses

Physical curve	“Yes”
Tail	NLO
Integral	“NLO” (depending on Q)
Extendible to multi-jet	Yes

MINLO-REVISITED V+1J



[Hamilton, Nason, Oleari, Zanderighi (2012);
Hamilton, Nason, Re, Zanderighi (2013);
RF, Hamilton (2015)]

- ♦ Much simpler as Geneva
- ♦ Like Minlo V+1j, include Sudakov form factors to make distribution physical at low p_T
- ♦ **Modify the Sudakov form factors with subleading, process dependent terms such that total integral becomes NLO accurate**
- ♦ Can include NNLO corrections for V

Physical curve	Yes
Tail	NLO
Integral	(N)NLO
Extendible to multi-jet	Yes

MINLO ACCURACY FOR (INCLUSIVE) 0-JET OBSERVABLES

- ◆ An explicit comparison between the diff.-jet-rate-resummation formula (which integrates to the correct NLO 0-jet diff. cross section) and Minlo shows that they differ by terms of order
[Banfi, Salam, Zanderighi (2005);
Dokshitzer, Diakonov, Troian (1980)]

$$\frac{d\sigma_{\mathcal{MR}}}{d\Phi dL} = \frac{d\sigma_0}{d\Phi} \exp[-R(v)] \prod_{\ell=1}^{n_i} \frac{q^{(\ell)}(x_\ell, \mu_F^2 v)}{q^{(\ell)}(x_\ell, \mu_F^2)} \left[\bar{\alpha}_s^2 (K_R^2 y) \left[\tilde{R}_{21} L + \tilde{R}_{20} \right] + \bar{\alpha}_s^3 (K_R^2 y) L^2 \tilde{R}_{32} \right]$$

- ◆ After integration over the logarithm L (taking $R_{21}=0$, which is okay for the processes considered here) this results into terms of

$$\int dL' \frac{d\sigma_{\mathcal{MR}}}{d\Phi dL'} = -\frac{d\sigma_0}{d\Phi} \left[\tilde{R}_{20} - \bar{\beta}_0 \mathcal{H}_1(\mu_R^2) \right] \sqrt{\frac{\pi}{2}} \frac{1}{|2G_{12}|^{1/2}} \bar{\alpha}_s^{3/2} (1 + \mathcal{O}(\sqrt{\bar{\alpha}_s}))$$

- ◆ Hence, diff. NLO-0jet cross section not correct with NLO-1jet Minlo

[Hamilton, Nason, Oleari, Zanderighi (2012);
RF, Hamilton (2015)]

MINLO ACCURACY FOR (INCLUSIVE) 0-JET OBSERVABLES

- ♦ An explicit comparison between the diff.-jet-rate-resummation formula (which integrates to the correct NLO 0-jet diff. cross section) and Minlo shows that they differ by terms of order
[Banfi, Salam, Zanderighi (2005);
Dokshitzer, Diakonov, Troian (1980)]

$$\frac{d\sigma_{\mathcal{MR}}}{d\Phi dL} = \frac{d\sigma_0}{d\Phi} \exp[-R(v)] \prod_{\ell=1}^{n_i} \frac{q^{(\ell)}(x_\ell, \mu_F^2 v)}{q^{(\ell)}(x_\ell, \mu_F^2)} \left[\bar{\alpha}_s^2(K_R^2 y) \left[\tilde{R}_{21} L + \tilde{R}_{20} \right] + \bar{\alpha}_s^3(K_R^2 y) L^2 \tilde{R}_{32} \right]$$

- ♦ After integration over the logarithm L (taking $R_{21}=0$, which is okay for the processes considered here) this results into terms of

$$\int dL' \frac{d\sigma_{\mathcal{MR}}}{d\Phi dL'} = -\frac{d\sigma_0}{d\Phi} \left[\tilde{R}_{20} - \bar{\beta}_0 \mathcal{H}_1(\mu_R^2) \right] \sqrt{\frac{\pi}{2}} \frac{1}{|2G_{12}|^{1/2}} \bar{\alpha}_s^{3/2} (1 + \mathcal{O}(\sqrt{\bar{\alpha}_s}))$$

- ♦ Hence, diff. NLO-0jet cross section not correct with NLO-1jet Minlo

[Hamilton, Nason, Oleari, Zanderighi (2012);
RF, Hamilton (2015)]

MINLO ACCURACY FOR (INCLUSIVE) 0-JET OBSERVABLES

Explicitly compute and remove that term in the Minlo calculation such that the integral $\int dL \frac{d\sigma_{\mathcal{MR}}}{d\Phi dL}$ is zero up to NLO

It's process dependent and not a constant in phase-space

♦ An explicit comparison between the diff-jet rate-resummation formula (which integrates to the correct NLO 0-jet diff. cross section) and Minlo shows that they differ by terms of order α_s^2 [Banfi, Salam, Zanderighi (2005); Dokshitzer, Diakonov, Troian (1980)]

$$\frac{d\sigma_{\mathcal{MR}}}{d\Phi dL} = \frac{d\sigma_0}{d\Phi} \exp[-R(v)] \prod_{\ell=1}^{n_i} \frac{q^{(\ell)}(x_\ell, \mu_F^2 v)}{q^{(\ell)}(x_\ell, \mu_F^2)} \left[\bar{\alpha}_s^2(K_R^2 y) \left[\tilde{R}_{21} L + \tilde{R}_{20} \right] + \bar{\alpha}_s^3(K_R^2 y) L^2 \tilde{R}_{32} \right]$$

♦ After integration over the logarithm L (taking $R_{21}=0$, which is okay for the processes considered here) this results into terms of

$$\int dL' \frac{d\sigma_{\mathcal{MR}}}{d\Phi dL'} = -\frac{d\sigma_0}{d\Phi} \left[\tilde{R}_{20} - \bar{\beta}_0 \mathcal{H}_1(\mu_R^2) \right] \sqrt{\frac{\pi}{2}} \frac{1}{|2G_{12}|^{1/2}} \bar{\alpha}_s^{3/2} (1 + \mathcal{O}(\sqrt{\bar{\alpha}_s}))$$

♦ Hence, diff. NLO-0jet cross section not correct with NLO-1jet Minlo

[Hamilton, Nason, Oleari, Zanderighi (2012); RF, Hamilton (2015)]

MINLO ACCURACY FOR (INCLUSIVE) 0-JET OBSERVABLES

Explicitly compute and remove that term in the Minlo calculation such that the integral $\int dL d\Phi dL'$ is zero up to NLO

It's process dependent and not a constant in phase-space

[Banfi, Salam, Zanderighi (2005);
Dokshitzer, Diakonov, Troian (1980)]

$$\frac{d\sigma_{\mathcal{MR}}}{d\Phi dL} = \frac{d\sigma_0}{d\Phi} \exp[-R(v)] \prod_{\ell=1}^{n_i} \frac{q^{(\ell)}(x_\ell, \mu_F^2 v)}{q^{(\ell)}(x_\ell, \mu_F^2)} \left[\bar{\alpha}_s^2(K_R^2 y) \left[\tilde{R}_{21} L + \tilde{R}_{20} \right] + \bar{\alpha}_s^3(K_R^2 y) L^2 \tilde{R}_{32} \right]$$

- After integration over the logarithm L (taking $R_{21}=0$, which is okay for the processes considered here) this results into terms of

$$\int dL' \frac{d\sigma_{\mathcal{MR}}}{d\Phi dL'} = -\frac{d\sigma_0}{d\Phi} \left[\tilde{R}_{20} - \bar{\beta}_0 \mathcal{H}_1(\mu_R^2) \right]$$

Can either be done analytically
or numerically by enforcing
unitarity

- Hence, diff. NLO-0jet cross section not correct with NLO 1jet Minlo

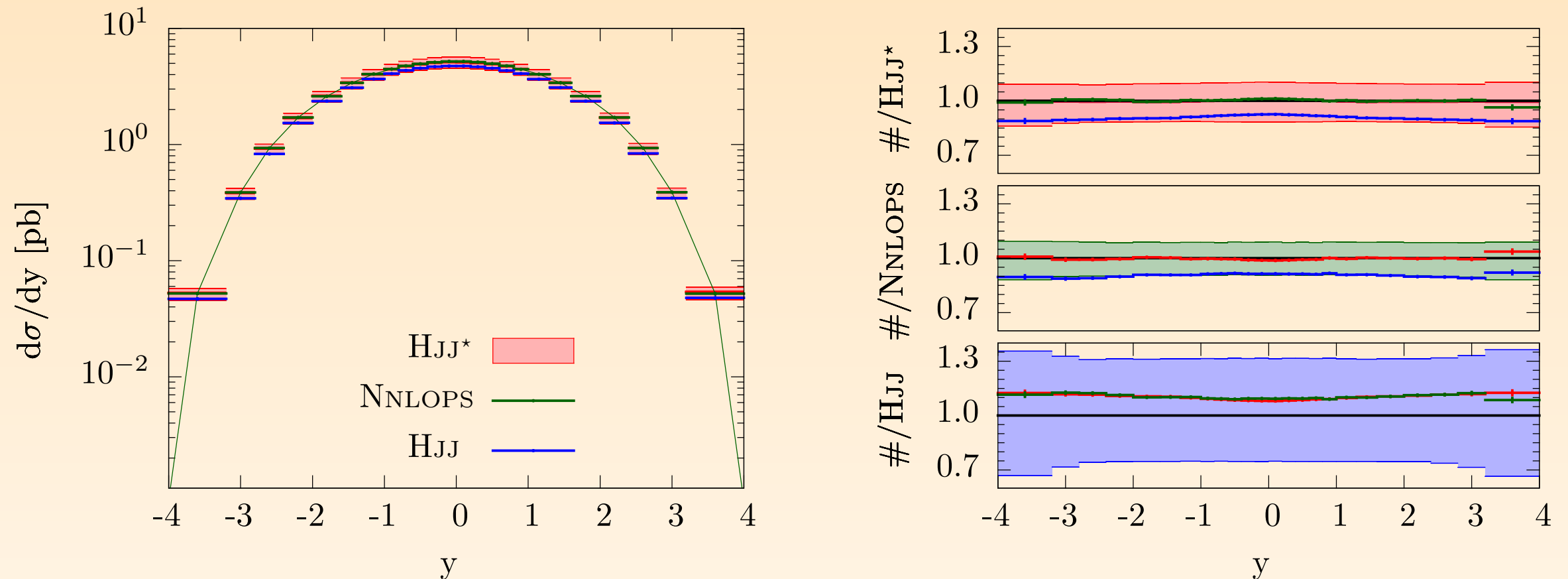
[Hamilton, Nason, Oleari, Zanderighi (2012);
RF, Hamilton (2015)]

PROOF-OF-CONCEPT

[RF, Hamilton (2015)]

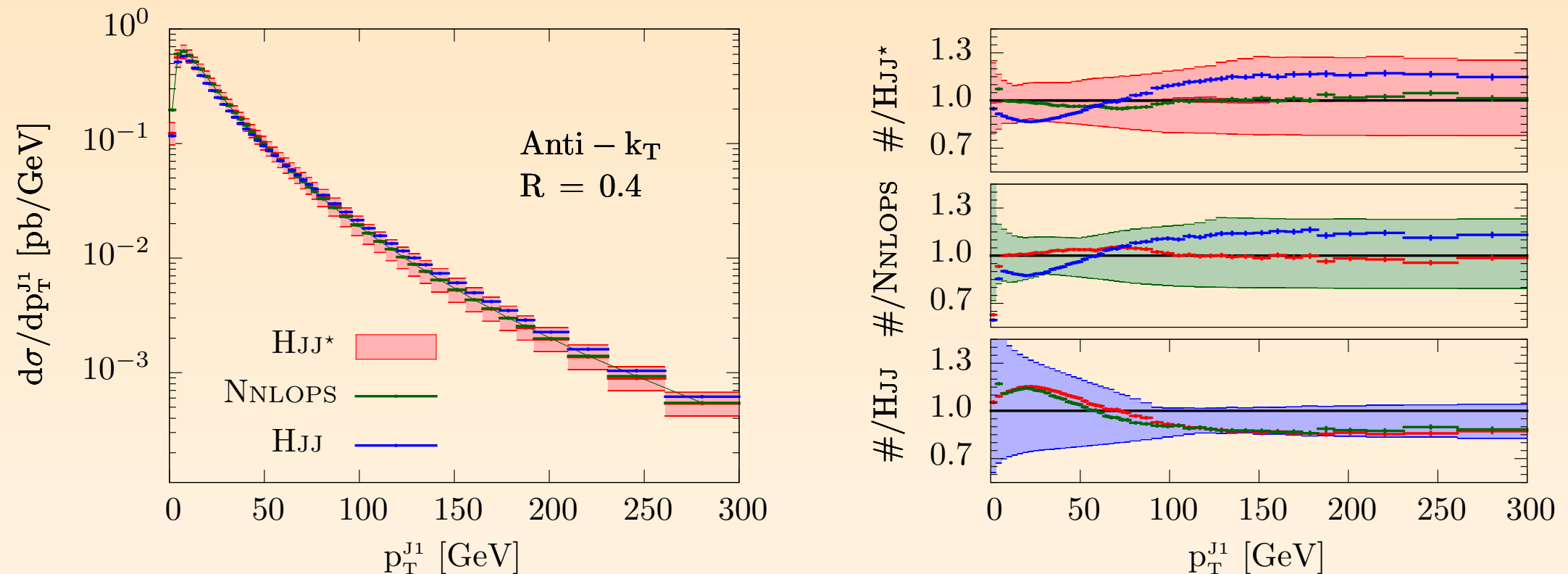
- ◆ Apply the method to **Higgs production by gluon fusion** in the infinite top quark limit (which is not a good approximation at high scales, but not a problem for a proof of concept)
- ◆ Start from H+J Minlo', corrected to include NNLO for H. Already available in the POWHEG BOX [Hamilton, Nason, Re, Zanderighi (2013)]
- ◆ Apply the extended Minlo' method to HJJ at NLO to get
 - NLO+PS predictions for inclusive HJJ observables
 - NLO+PS predictions for inclusive HJ observables
 - NNLO+PS predictions for inclusive H observables
- ◆ Study **renormalisation/factorisation scale dependence** and dependence on freezing parameter ρ (which we vary $\rho=\{1, 3, 9, 18, 27\}$)

RAPIDITY OF THE HIGGS BOSON



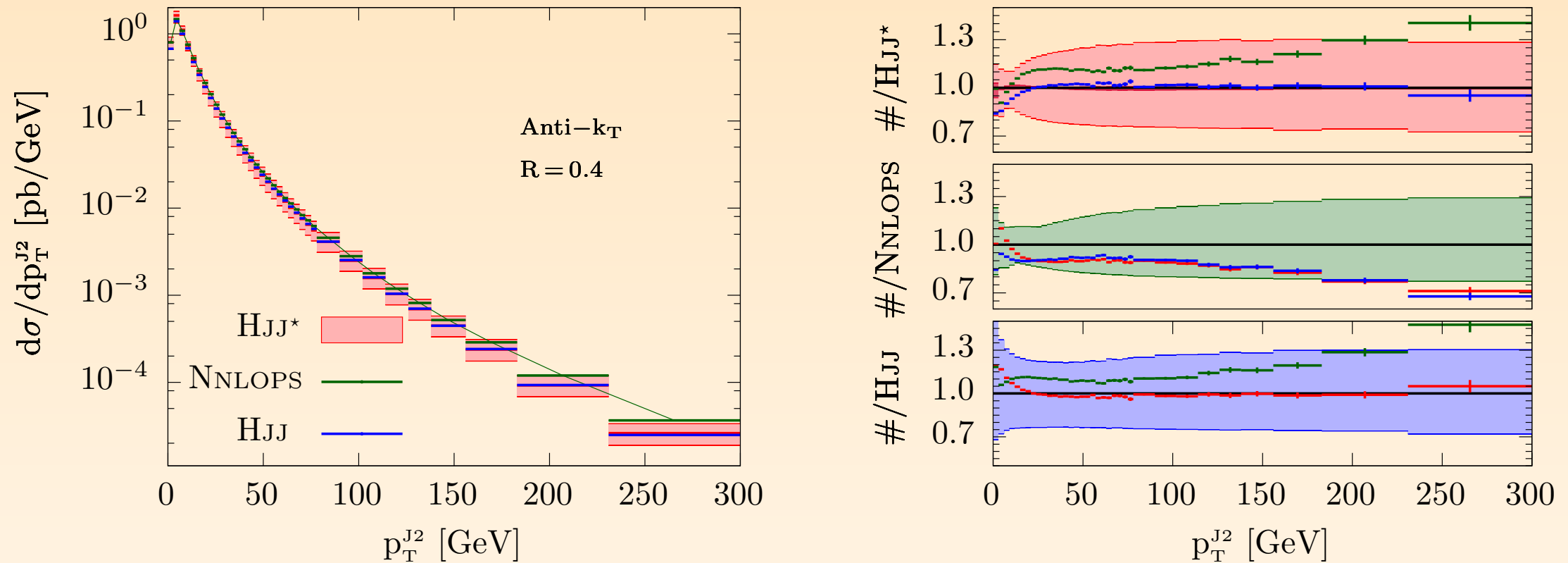
- ◆ Only observable truly NNLO correct
- ◆ Extended Minlo' method (**HJJ★**) agrees with **NNLOPS** by construction
- ◆ Normal **HJJ** Minlo shows larger uncertainty bands and different central value: it's only LO accurate for this observable

TRANSVERSE MOMENTUM OF THE LEADING JET



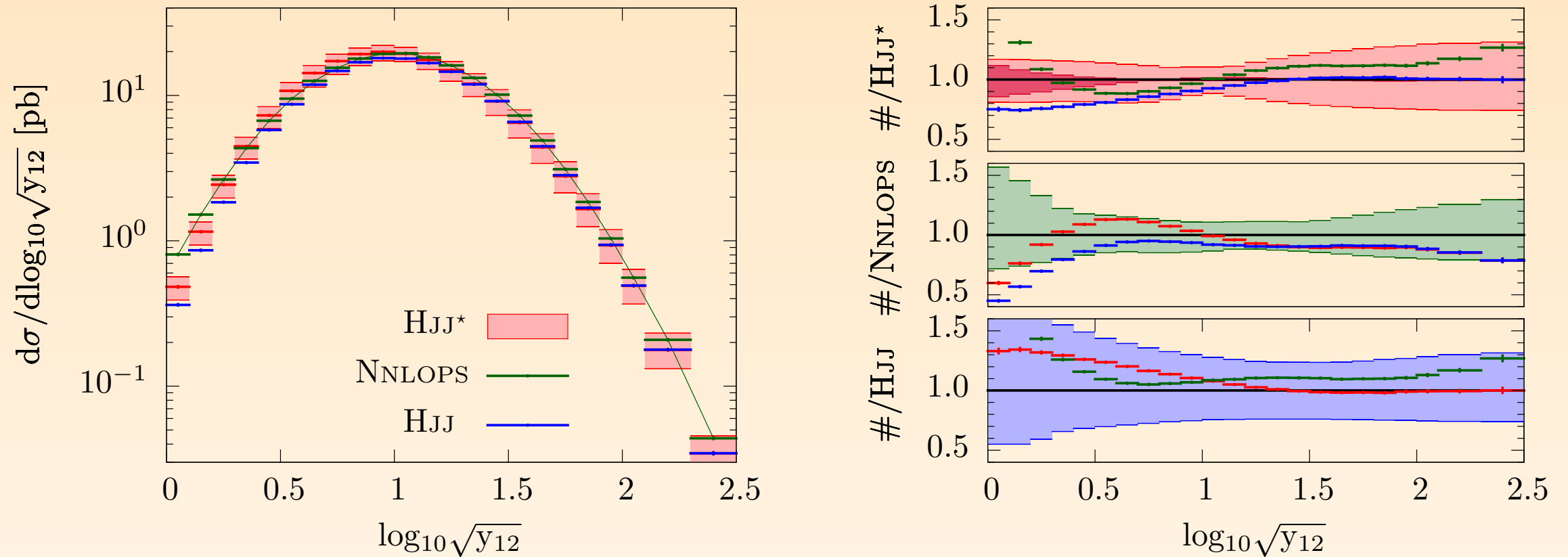
- ◆ Extended Minlo' method (**HJJ***) agrees with **NNLOPS** by construction.
 - apart from $p_T < 5$ GeV region: grid-granularity to compute δ not fine enough
 - Also region $60 < p_T < 80$ GeV shows 3-5% deviations: p_T derivative of the numerator of δ changes very rapidly
- ◆ Normal **HJJ** Minlo shows unphysical uncertainty band. Formally only LO for this observable

TRANSVERSE MOMENTUM OF THE SECOND JET



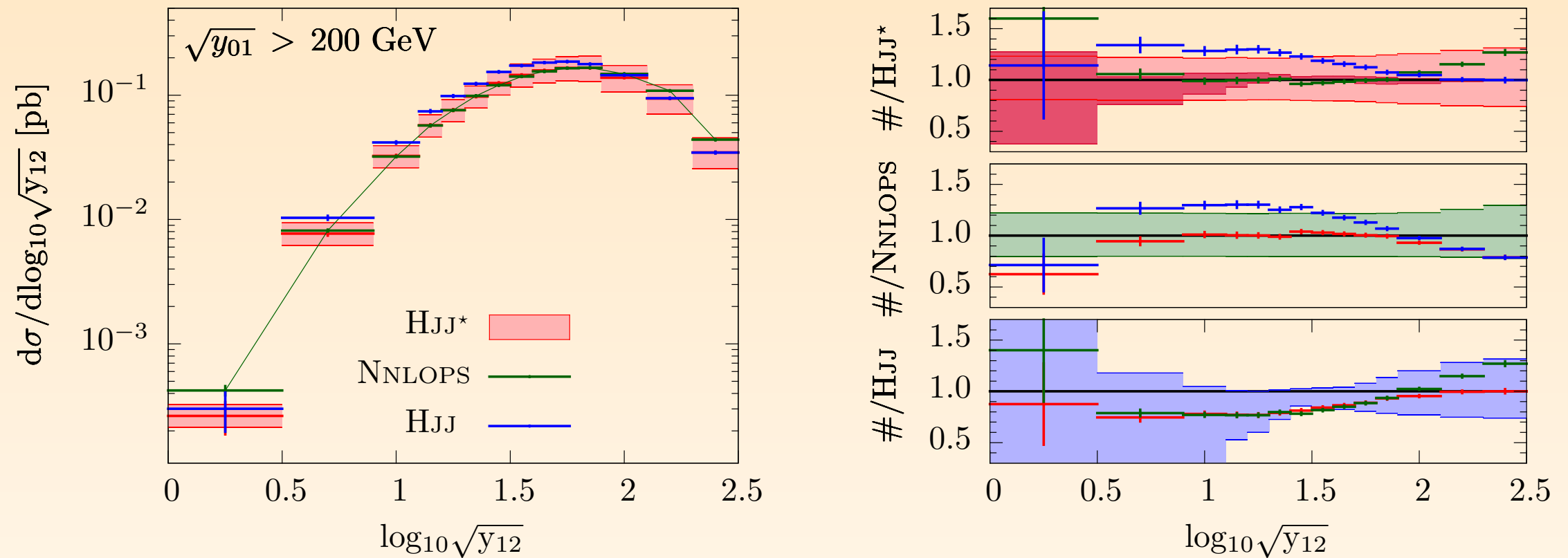
- ◆ Extended Minlo' method **HJJ*** agrees with Minlo **HJJ**, as expected
 - apart close to the Sudakov peak: the difference between **HJJ*** and **HJJ** is beyond LL/NNLL_σ accuracy, which is important close to the Sudakov peak
- ◆ **NNLOPS** only LO accurate for this observable: uncertainty band is too small (this is due to the POWHEG method)

Y_{12} RESOLUTION PARAMETER



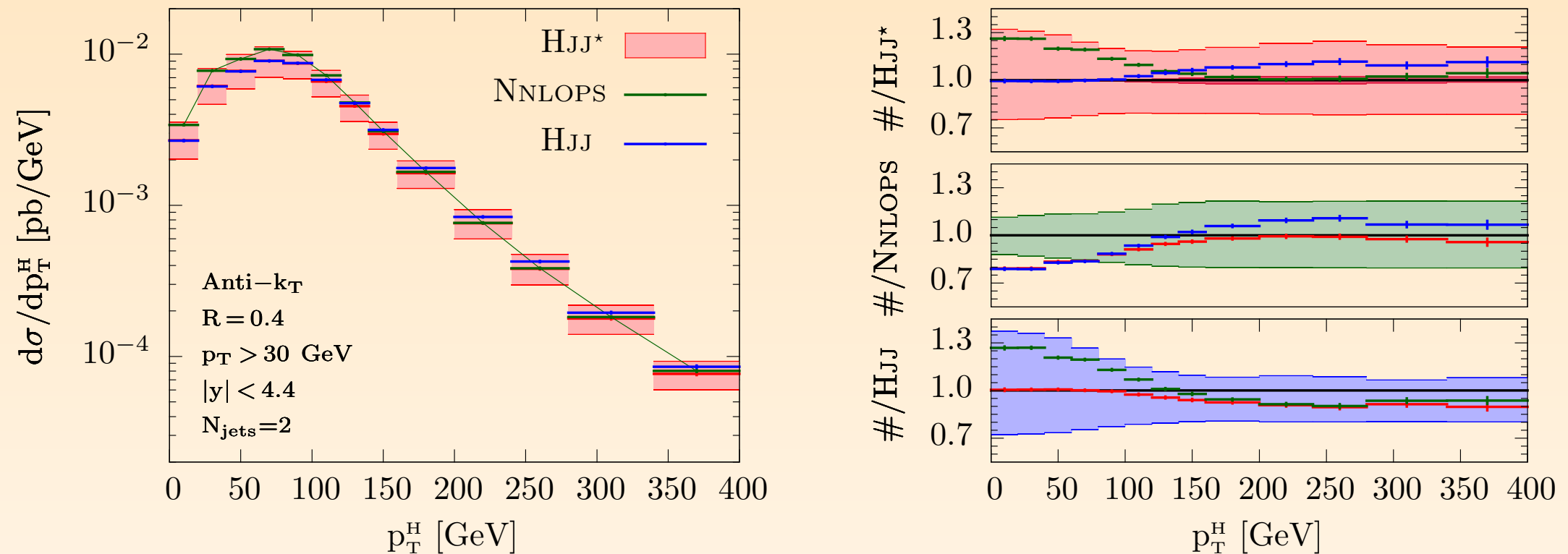
- ◆ Similar picture as for $p_T(j_2)$, but low p_T region easier to see due to logarithmic x-axis
- ◆ First observable where we see some non-zero dependence on the freezing parameter ρ (red solid). Well below the Sudakov peak where higher-logarithmic corrections are large as well as non-perturbative corrections

Y₁₂ RESOLUTION PARAMETER WITH $\sqrt{y_{01}} > 200$ GeV



- ♦ At very large y_{12} , all scales are large and of the same order \rightarrow the Minlo method switches off: **HJJ*** agrees with **HJJ**
- ♦ When $y_{12} \ll y_{01}$, large logarithms build up, and the extended Minlo' method brings the **HJJ*** to the **NNLOPS**

HIGGS BOSON p_T IN EVENTS WITH EXACTLY 2 JETS



- ◆ At small p_T , all scales are of the same order. The Minlo method does not do much: **HJJ*** agrees with **HJJ**
- ◆ At large p_T , **HJJ*** agrees with **NNLOPS** dominated by events with one hard jet ($p_T(j_1) \sim p_T(H)$) and one soft jet: a 30 GeV jet comes basically for free
 - The $p_T(H)$ spectrum with $N_{\text{jets}}=2$ becomes essentially $N_{\text{jets}} \geq 1$ $p_T(H)$ distribution

CONCLUSIONS

- ♦ In the last couple of years the accuracy of event generation has greatly improved, and full automation has been achieved at NLO accuracy
 - FxFx Merging is one of the methods to combine NLO matrix elements of various multiplicities with the parton shower
 - NLO accuracy in multiple regions of phase-space, separated by a merging scale
- ♦ A lot of freedom in tuning has been replaced by accurate theory descriptions:
 - More predictive power
 - Better control on uncertainties in predictions
 - Greater trust in the measurements
- ♦ One of the latest developments, 'Minlo revisited', allows for similar accuracy as FxFx in multi-jets, but without the introduction of a merging scale and with the possibility to include NNLO. Only proof-of-concept so far.