

LABORATORIUM FÜR HOCHENERGIEPHYSIK

LHEP

UNIVERSITÄT BERN



Sterile Neutrinos

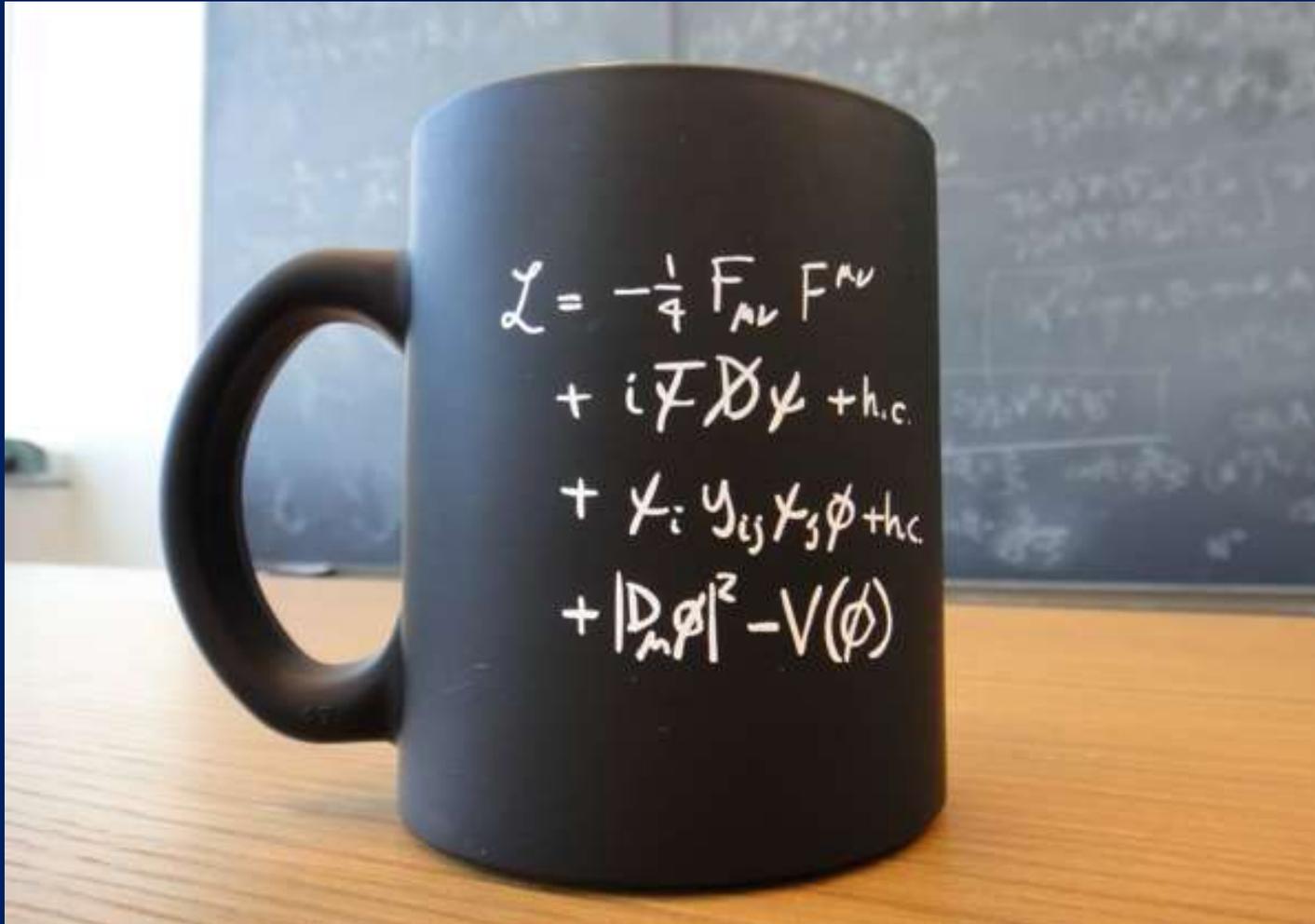
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LHEP-AEC, Uni Bern



Seminar, Freiburg (D), June 2019

Standard Model of particle physics



$$\begin{aligned}
 & -\frac{1}{2}\partial_\mu g_\nu^\rho \partial_\mu g_\nu^\rho - g_\nu^{\rho\sigma} \partial_\mu g_\nu^\rho g_\mu^\sigma - \frac{1}{2}g_\nu^{\rho\sigma} g_\mu^\rho g_\nu^\sigma g_\mu^\sigma + \\
 & \frac{1}{2}ig^2(\bar{\psi}\gamma^\mu \psi)g_\mu^\rho + G^a \partial^\mu G^a + g_s f^{abc} \partial_\mu G^a G^b G^c - \partial_\mu W_\nu^+ \partial_\mu W_\nu^- - \\
 & M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\mu Z_\nu^0 \partial_\mu Z_\nu^0 - \frac{1}{2}M^2 Z_\nu^0 Z_\nu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \\
 & \frac{1}{2}m_H^2 H^2 - \partial_\mu \phi^\dagger \partial_\mu \phi - M^2 \phi^\dagger \phi - \frac{1}{2}\partial_\mu \phi^\dagger \partial_\mu \phi - \frac{1}{2}M^2 \phi^\dagger \phi - \beta_h \left[\frac{2M^2}{\Lambda^2} + \right. \\
 & \left. \frac{2M}{\Lambda} H + \frac{1}{2}(H^2 + \phi^\dagger \phi + 2\phi^\dagger \phi) \right] + \frac{2M^2}{\Lambda^2} \alpha_h - ig_{\text{CW}} [\partial_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
 & W_\mu^- W_\nu^+) - Z_\nu^0 (W_\mu^+ \partial_\mu W_\nu^- - W_\mu^- \partial_\mu W_\nu^+) + Z_\nu^0 (W_\mu^+ \partial_\mu W_\nu^- - \\
 & W_\mu^- \partial_\mu W_\nu^+) - ig_{\text{SW}} [\partial_\mu A_\nu (W_\mu^+ W_\nu^- - W_\mu^- W_\nu^+) - A_\nu (W_\mu^+ \partial_\mu W_\nu^- - \\
 & W_\mu^- \partial_\mu W_\nu^+) + A_\nu (W_\mu^+ \partial_\mu W_\nu^- - W_\mu^- \partial_\mu W_\nu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + \\
 & \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 Z_\nu^0 (Z_\mu^0 W_\nu^+ Z_\mu^0 W_\nu^- - Z_\nu^0 Z_\mu^0 W_\nu^+ W_\nu^-) + \\
 & g^2 Z_\nu^0 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_{\text{CW}} [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
 & W_\mu^- W_\nu^+) - 2A_\mu Z_\nu^0 W_\mu^+ W_\nu^-] - g\alpha [H^2 + H \phi^\dagger \phi + 2H \phi^\dagger \phi] - \\
 & \frac{1}{2}g^2 \alpha_h [H^4 + (\phi^\dagger)^4 + 4(\phi^\dagger)^2 \phi^2 + 4(\phi^\dagger)^2 \phi^\dagger \phi + 4H^2 \phi^\dagger \phi + 2(\phi^\dagger)^2 H^2] - \\
 & gM W_\mu^+ W_\nu^- H - \frac{1}{2}g \frac{M}{\Lambda^2} Z_\nu^0 Z_\mu^0 H - \frac{1}{2}ig [W_\mu^+ (\partial^\mu \phi^\dagger \phi - \phi^\dagger \partial_\mu \phi) - \\
 & W_\mu^- (\partial^\mu \phi^\dagger \phi - \phi^\dagger \partial_\mu \phi)] + \frac{1}{2}g [W_\mu^+ (H \partial_\mu \phi^\dagger - \phi^\dagger \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^\dagger - \\
 & \phi^\dagger \partial_\mu H)] + \frac{1}{2}g \frac{1}{\Lambda^2} (Z_\nu^0 (H \partial_\mu \phi^\dagger - \phi^\dagger \partial_\mu H) - ig \frac{M}{\Lambda^2} M Z_\nu^0 (W_\mu^+ \phi^\dagger - W_\mu^- \phi^\dagger) + \\
 & ig_{\text{SW}} M A_\mu (W_\mu^+ \phi^\dagger - W_\mu^- \phi^\dagger) - ig \frac{1}{\Lambda^2} Z_\nu^0 (\phi^\dagger \partial_\mu \phi - \phi^\dagger \partial_\mu \phi) + \\
 & ig_{\text{SW}} A_\mu (\phi^\dagger \partial_\mu \phi - \phi^\dagger \partial_\mu \phi) - \frac{1}{2}g^2 W_\mu^+ W_\nu^- [H^2 + (\phi^\dagger)^2 + 2\phi^\dagger \phi] - \\
 & \frac{1}{2}g^2 \frac{1}{\Lambda^2} Z_\nu^0 Z_\mu^0 [H^2 + (\phi^\dagger)^2 + 2(2c_{\text{W}}^2 - 1)\phi^\dagger \phi] - \frac{1}{2}g^2 \frac{1}{\Lambda^2} Z_\nu^0 \phi^\dagger (W_\mu^+ \phi^\dagger + \\
 & W_\mu^- \phi^\dagger) - \frac{1}{2}ig^2 \frac{1}{\Lambda^2} Z_\nu^0 H (W_\mu^+ \phi^\dagger - W_\mu^- \phi^\dagger) + \frac{1}{2}g^2 s_{\text{W}} A_\mu \phi^\dagger (W_\mu^+ \phi^\dagger - \\
 & W_\mu^- \phi^\dagger) + \frac{1}{2}ig^2 s_{\text{W}} A_\mu H (W_\mu^+ \phi^\dagger - W_\mu^- \phi^\dagger) - g^2 \frac{1}{\Lambda^2} (2c_{\text{W}}^2 - 1) Z_\nu^0 A_\mu \phi^\dagger \phi - \\
 & g^4 s_{\text{W}}^2 A_\mu A_\nu \phi^\dagger \phi - e^4 (\gamma \partial + m_e^2) e^\lambda - e^4 \gamma \partial e^\lambda - m_e^2 (\gamma \partial + m_e^2) u_e^\lambda - m_e^2 (\gamma \partial + \\
 & m_e^2) d_e^\lambda + ig_{\text{SW}} A_\mu [-(e^\lambda \gamma e^\lambda) + \frac{2}{3}(\bar{u}_e^\lambda \gamma u_e^\lambda) - \frac{1}{3}(\bar{d}_e^\lambda \gamma d_e^\lambda)] + \frac{1}{2}g_{\text{W}} Z_\nu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \\
 & \gamma^5) \nu^\lambda) + (e^\lambda \gamma^\mu (4c_{\text{W}}^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_e^\lambda \gamma^\mu (\frac{1}{3} s_{\text{W}}^2 - 1 - \gamma^5) u_e^\lambda) + \\
 & (\bar{d}_e^\lambda \gamma^\mu (1 - \frac{2}{3} s_{\text{W}}^2 - \gamma^5) d_e^\lambda)] + \frac{1}{2}g_{\text{W}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{u}_e^\lambda \gamma^\mu (1 + \gamma^5) u_e^\lambda) + \\
 & (\bar{d}_e^\lambda \gamma^\mu C_{\text{W}} d_e^\lambda)] + \frac{1}{2}g_{\text{W}} W_\mu^- [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_e^\lambda \gamma^\mu C_{\text{W}} d_e^\lambda) + (\bar{u}_e^\lambda \gamma^\mu (1 + \gamma^5) u_e^\lambda)] + \\
 & \frac{1}{2}g_{\text{W}} \frac{m_e^2}{\Lambda^2} [-(\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^\dagger (e^\lambda (1 + \gamma^5) \nu^\lambda) - \frac{2}{3} \frac{m_e^2}{\Lambda^2} [H (e^\lambda e^\lambda) + \\
 & i\phi^\dagger (e^\lambda \gamma^5 e^\lambda)] + \frac{1}{2M} \frac{m_e^2}{\Lambda^2} \phi^\dagger [-m_e^2 (\bar{u}_e^\lambda C_{\text{W}} (1 - \gamma^5) d_e^\lambda) + m_e^2 (\bar{u}_e^\lambda C_{\text{W}} (1 + \\
 & \gamma^5) d_e^\lambda) + \frac{1}{2M} \frac{m_e^2}{\Lambda^2} \phi^\dagger [m_e^2 (\bar{d}_e^\lambda C_{\text{W}} (1 + \gamma^5) u_e^\lambda) - m_e^2 (\bar{d}_e^\lambda C_{\text{W}} (1 - \gamma^5) u_e^\lambda) - \\
 & \frac{g}{2} \frac{m_e^2}{M} H (\bar{u}_e^\lambda u_e^\lambda) - \frac{g}{2} \frac{m_e^2}{M} H (\bar{d}_e^\lambda d_e^\lambda) + \frac{1}{2} \frac{m_e^2}{M} \phi^\dagger (\bar{u}_e^\lambda \gamma^5 u_e^\lambda) - \frac{1}{2} \frac{m_e^2}{M} \phi^\dagger (\bar{d}_e^\lambda \gamma^5 d_e^\lambda) + \\
 & X^+ (\partial^2 - M^2) X^+ + X^- (\partial^2 - M^2) X^- + X^0 (\partial^2 - \frac{M^2}{2}) X^0 + Y \partial^2 Y + \\
 & ig_{\text{CW}} W_\mu^+ (\partial_\mu X^0 X^- - \partial_\mu X^+ X^0) + ig_{\text{CW}} W_\mu^+ (\partial_\mu Y X^- - \partial_\mu X^+ Y) + \\
 & ig_{\text{CW}} W_\mu^- (\partial_\mu X^- X^0 - \partial_\mu X^0 X^+) + ig_{\text{CW}} W_\mu^- (\partial_\mu X^- Y - \partial_\mu Y X^+) + \\
 & ig_{\text{CW}} Z_\nu^0 (\partial_\mu X^+ X^- - \partial_\mu X^- X^+) + ig_{\text{CW}} A_\mu (\partial_\mu X^+ X^- - \partial_\mu X^- X^+) - \\
 & \frac{1}{2}g M [X^+ X^+ H + X^- X^- H + \frac{1}{2} X^0 X^0 H] + \frac{1}{2} \frac{g}{\Lambda^2} ig M [X^+ X^0 \phi^\dagger - \\
 & X^- X^0 \phi^\dagger] + \frac{1}{2} ig M [X^0 X^- \phi^\dagger - X^0 X^+ \phi^\dagger] + ig M s_{\text{W}} [X^0 X^- \phi^\dagger - \\
 & X^0 X^+ \phi^\dagger] + \frac{1}{2} ig M [X^+ X^+ \phi^\dagger - X^- X^- \phi^\dagger]
 \end{aligned}$$

Are we done? No !

- Gravity is not included
- Matter-antimatter asymmetry?
- Large tuning needed to obtain the measured low Higgs mass (natural ?)
- Unification of electroweak and strong forces?
- Dark matter and dark energy?
- Origin of neutrino masses?
- Additional neutrino states ?
- ...



Many scientists and projects addressing these questions

I will discuss neutrinos today, focusing on sterile neutrinos



Wolfgang Pauli

$$i\bar{\psi}\gamma^{\mu}\psi$$

$$\frac{ig}{2\sqrt{2}}W_{\mu}^{-}[(\bar{e}^{\lambda}\gamma^{\mu}(1+\gamma^5)\nu^{\lambda})]$$

Coupling = small



Enrico Fermi

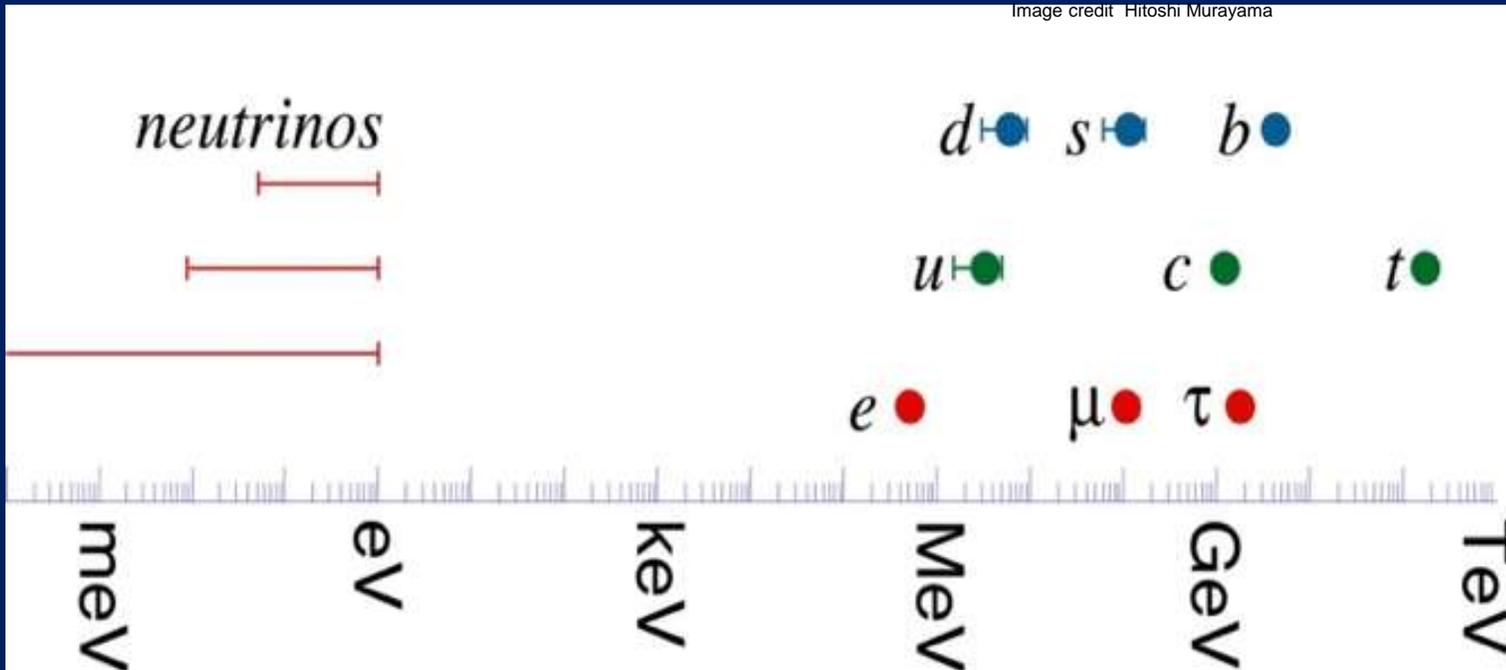
Neutrino Energy
few MeV

$$\sigma = 10^{-44} \text{ cm}^2$$

Interaction
Probability = 10^{-11}

$$10^{-39} \text{ cm}^2 = 1 \text{ fb}$$

Image credit Hitoshi Murayama



W und Z
Bosons



Photon



Gluons

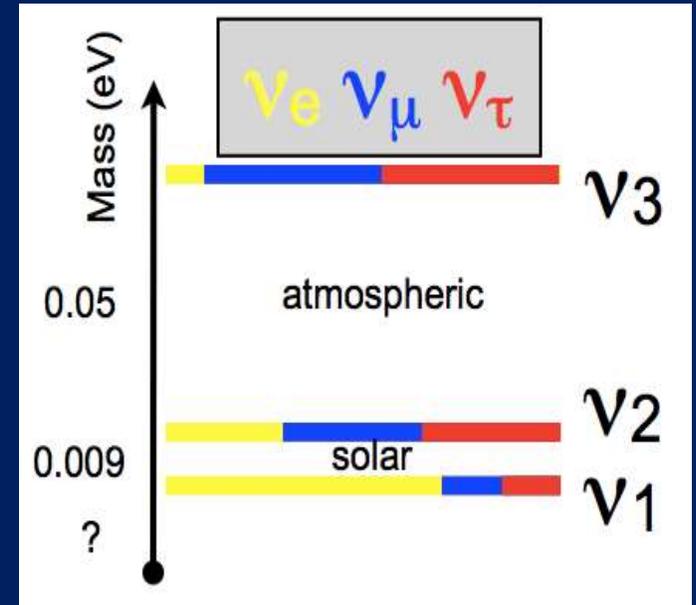


Higgs Boson

Neutrino oscillations

$$\begin{array}{c} \text{flavor} \\ \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \end{array} = \begin{array}{c} \text{atmospheric} \\ \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \end{array} \begin{array}{c} \text{cross-mixing} \\ \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \end{array} \begin{array}{c} \text{solar} \\ \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \end{array} \begin{array}{c} \text{mass} \\ \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \end{array}$$

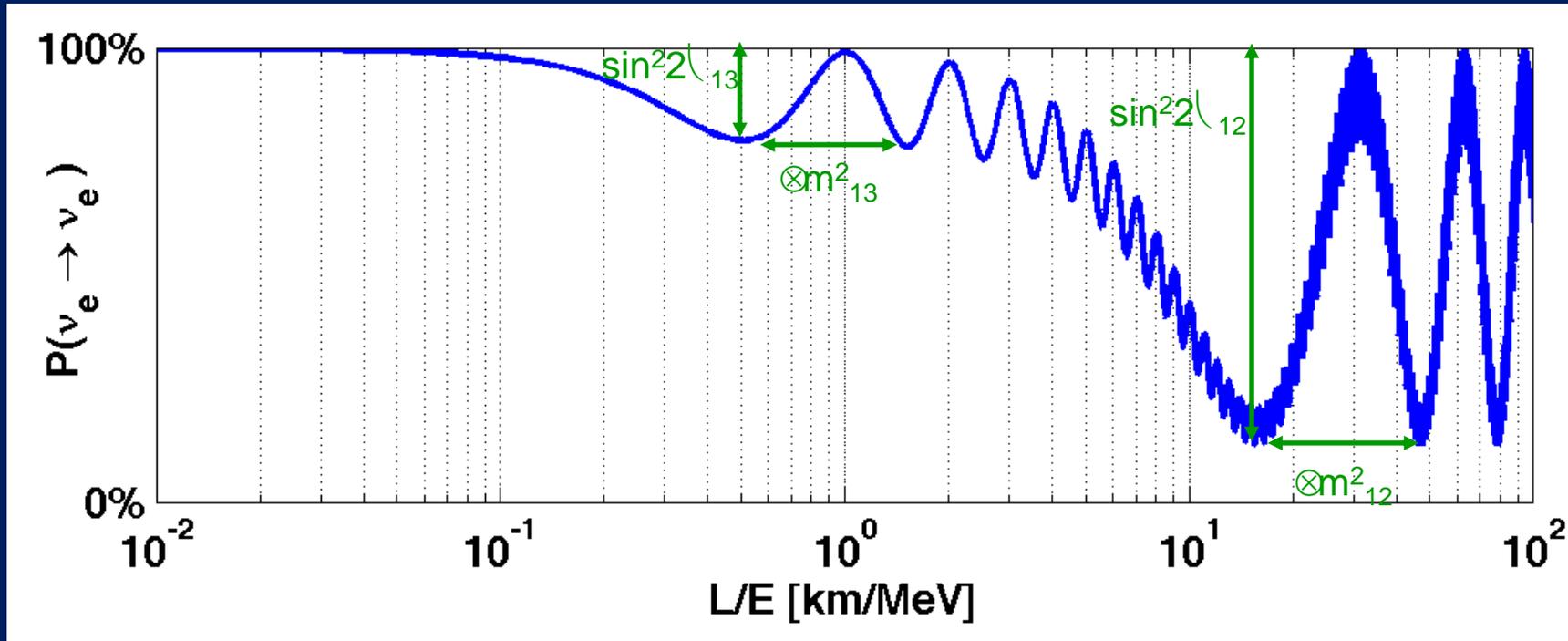
The PMNS matrix



Two-neutrino approximation

$$P(n_a \rightarrow n_b) = \sin^2 2\theta_{ij} \cdot \sin^2 \left(1.27 D m_{ij}^2 \frac{L}{E} \right)$$

- Three angles θ_{12} , θ_{23} and θ_{13}
- One phase δ_{CP}
- Two mass-squared differences Δm_{21}^2 and Δm_{23}^2



$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 1 - \sin^2 2\theta_{13} \sin^2 \left(\frac{\Delta m^2_{13} L}{4E} \right) - \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \left(\frac{\Delta m^2_{12} L}{4E} \right),$$

A global experimental effort

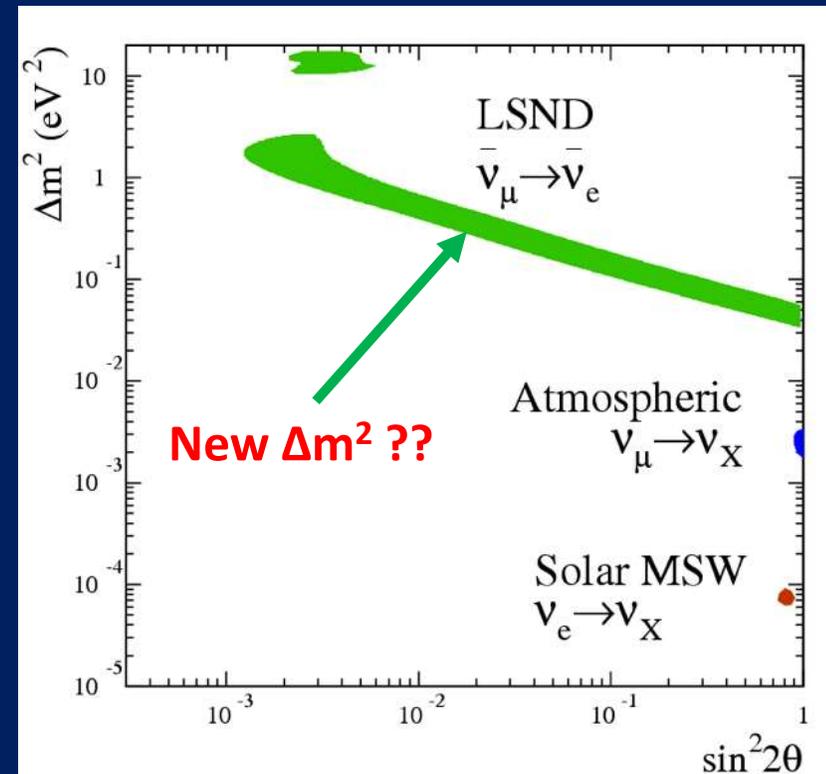
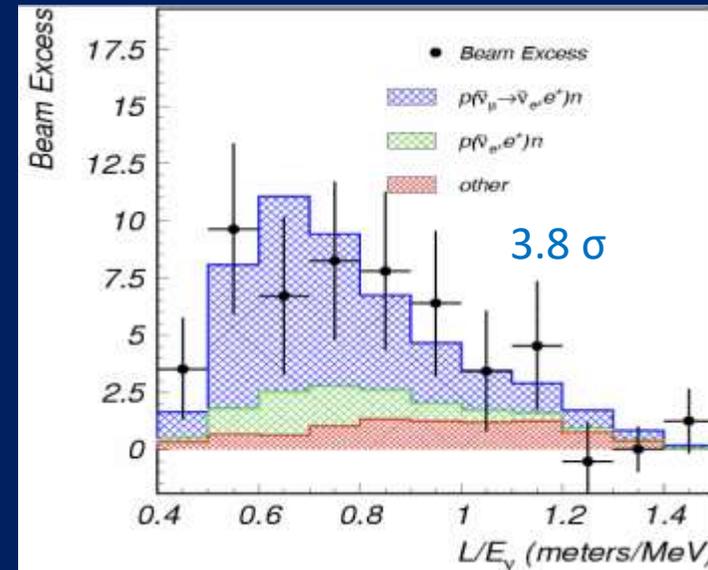
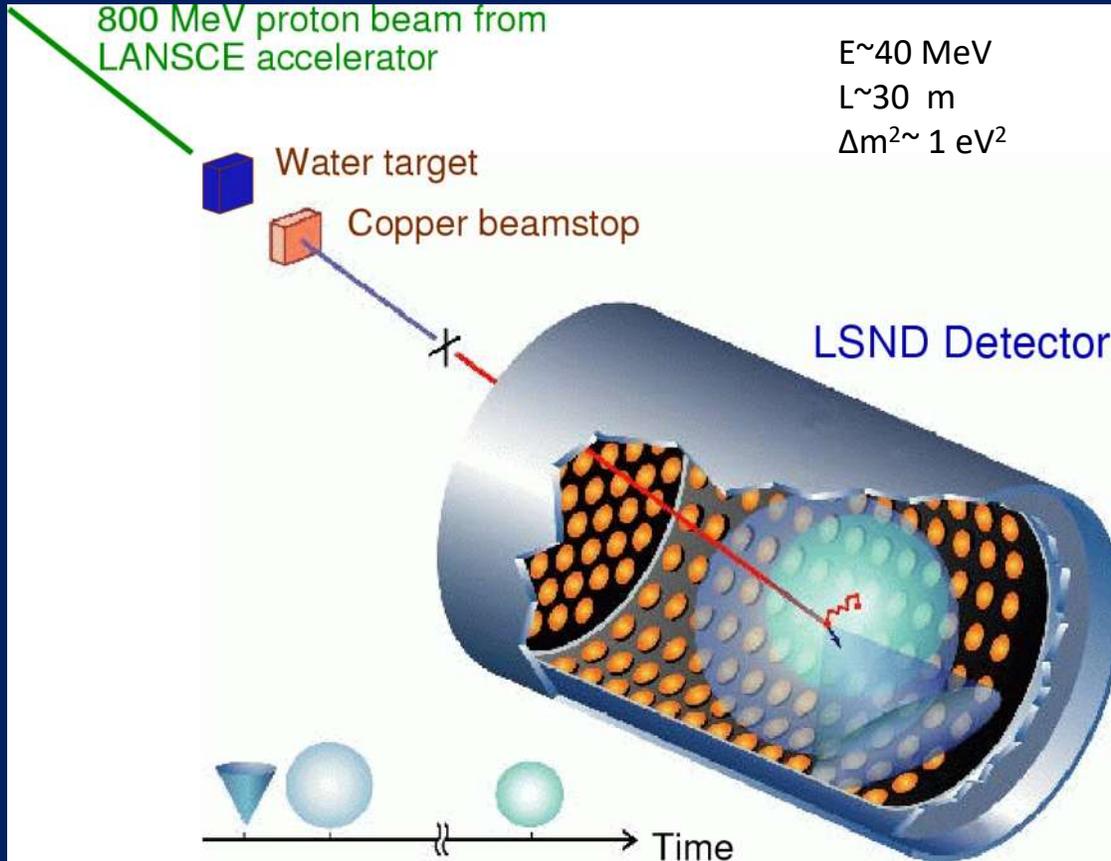
KamLAND, Double Chooz, Daya Bay, RENO, Gösgen, SNO, SK, T2K, MINOS, NOvA, OPERA, ICARUS, IceCube, ANTARES, ...

Parameter	best-fit	3σ
$\Delta m_{21}^2 [10^{-5} \text{ eV}^2]$	7.37	6.93 – 7.96
$\Delta m_{31(23)}^2 [10^{-3} \text{ eV}^2]$	2.56 (2.54)	2.45 – 2.69 (2.42 – 2.66)
$\sin^2 \theta_{12}$	0.297	0.250 – 0.354
$\sin^2 \theta_{23}, \Delta m_{31(32)}^2 > 0$	0.425	0.381 – 0.615
$\sin^2 \theta_{23}, \Delta m_{32(31)}^2 < 0$	0.589	0.384 – 0.636
$\sin^2 \theta_{13}, \Delta m_{31(32)}^2 > 0$	0.0215	0.0190 – 0.0240
$\sin^2 \theta_{13}, \Delta m_{32(31)}^2 < 0$	0.0216	0.0190 – 0.0242
δ/π	1.38 (1.31)	2σ : (1.0 - 1.9) (2σ : (0.92-1.88))

PDG update Dec 2017

LSND

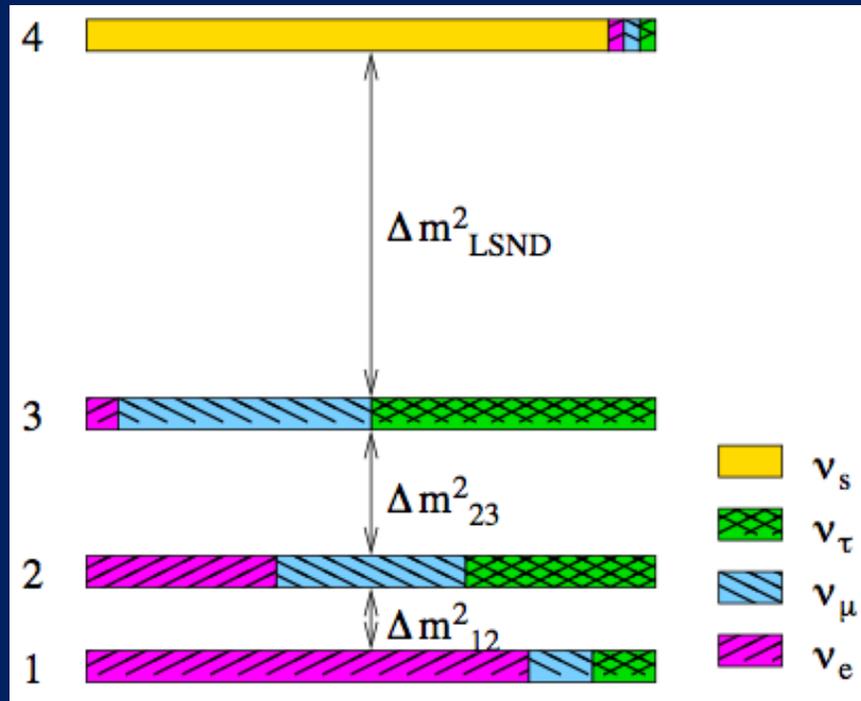
The Liquid Scintillator Neutrino Detector
at Los Alamos National Lab
1993 -- 1998



Two-neutrino approximation

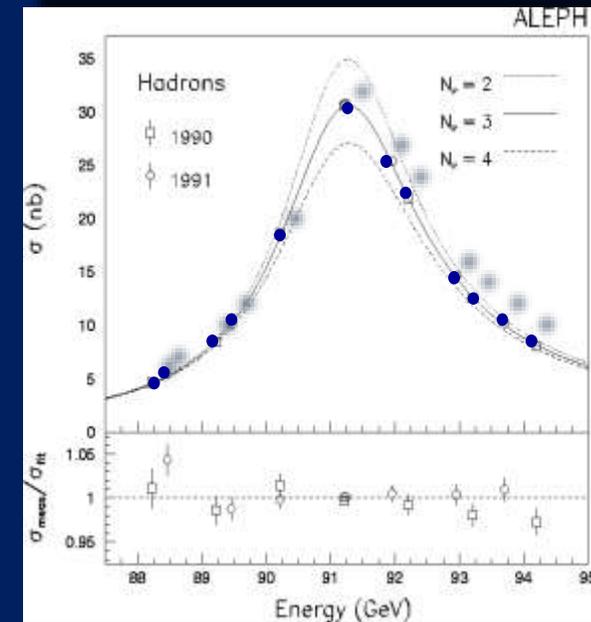
$$P(n_a \rightarrow n_b) = \sin^2 2q_{ij} \cdot \sin^2 \left(1.27 D m_{ij}^2 \frac{L}{E} \right)$$

4th neutrino mass state



Constraints

$N_\nu=3$ from LEP

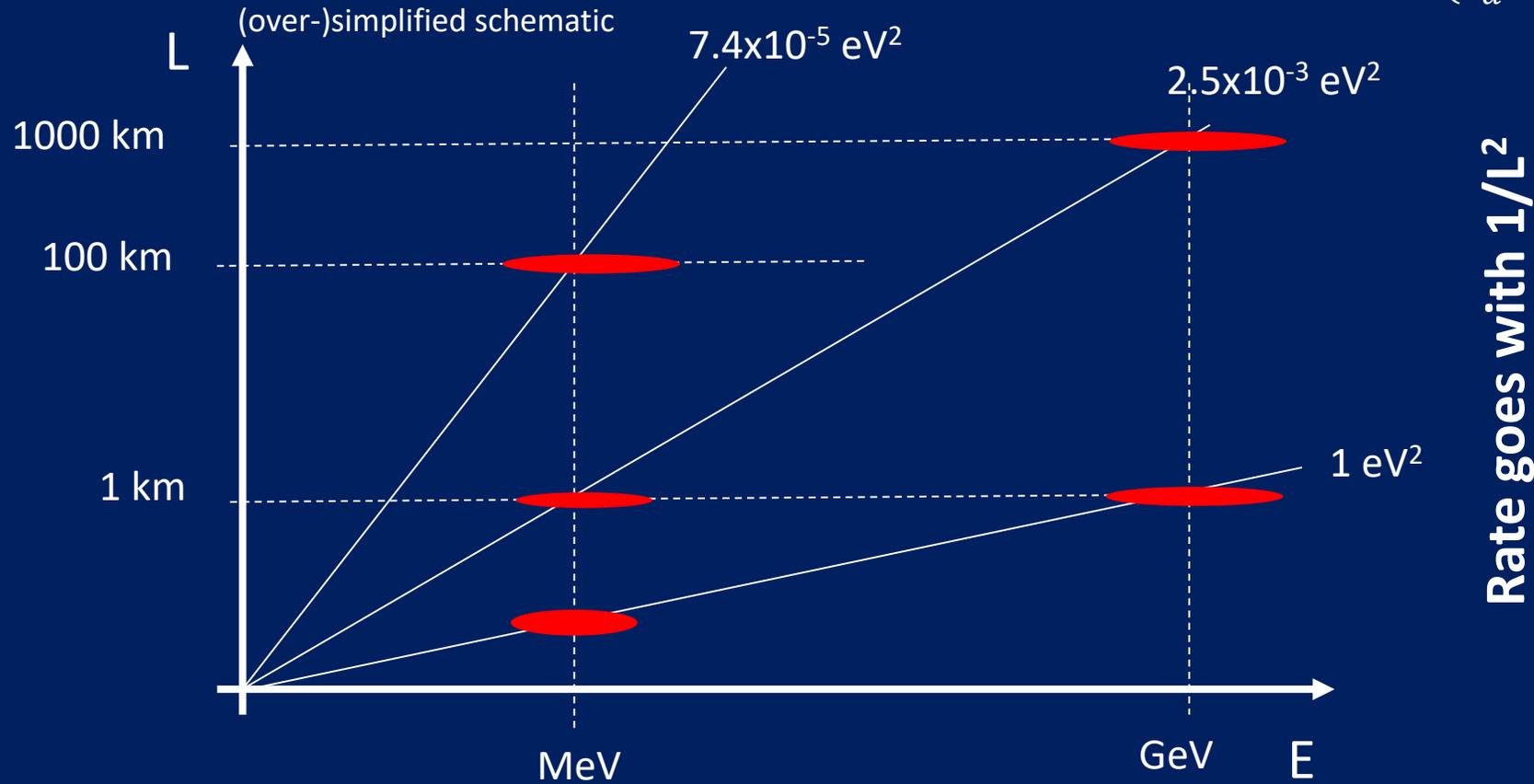


Cosmology

Λ CDM is sensitive to N_ν
(Large scale structures (BAO) and light nuclei abundance.)
Best fit is also consistent with $N_\nu = 3$ active neutrinos.

→ Sterile

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta_{ij} \cdot \sin^2 \left(1.27 \Delta m_{ij}^2 \frac{L}{E} \right)$$



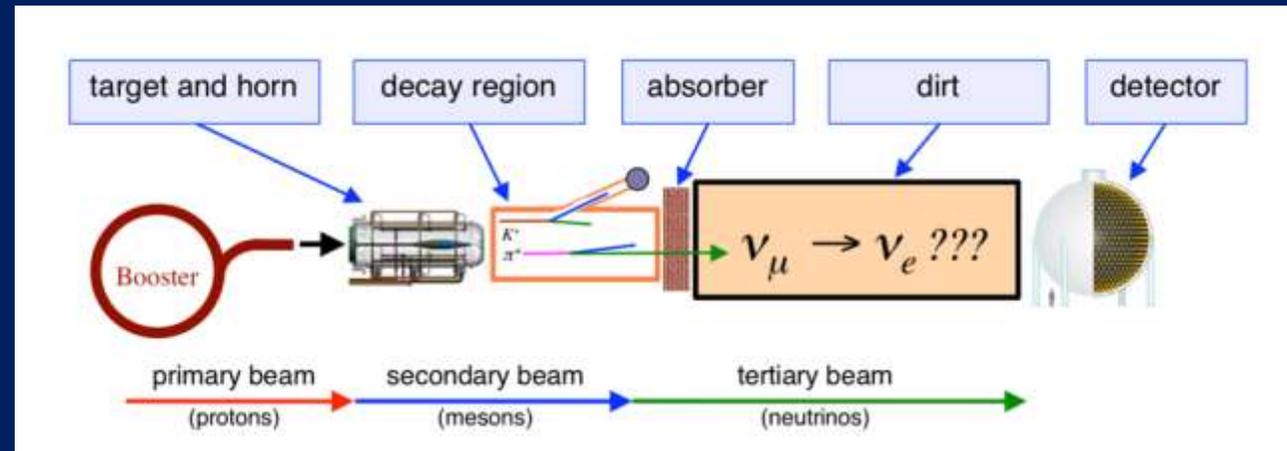
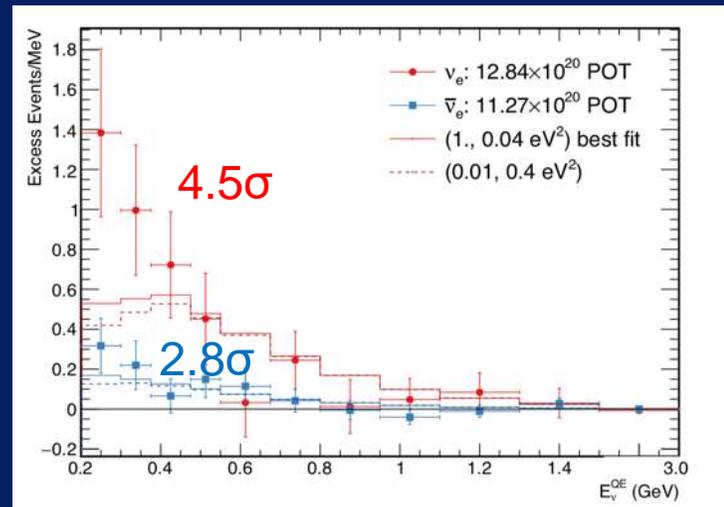
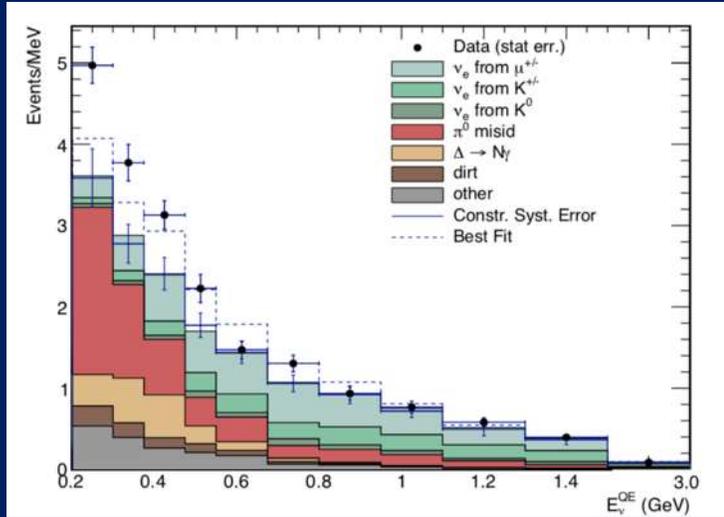
- **Short Base-Line:**

- Sensitivity to high Δm^2
- High rates, statistics to measure cross sections

- Three (four) neutrinos
- Neutrino and Antineutrino
- Appearance and disappearance
- Mixing angles

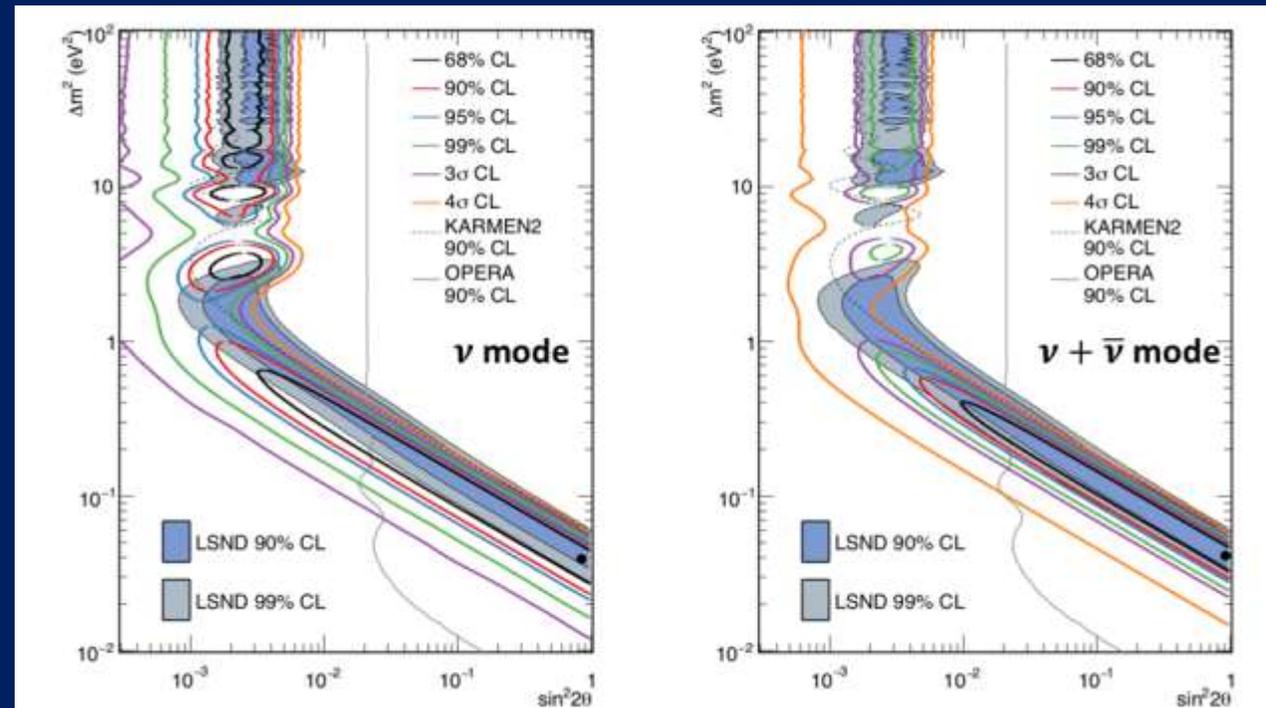
MiniBooNE

@Fermilab (2002--)



$L/E \sim 1 \text{ MeV/m}$
 $E \sim 800 \text{ MeV}$
 $L \sim 600 \text{ m}$

$\Delta m^2 \sim 1 \text{ eV}^2$



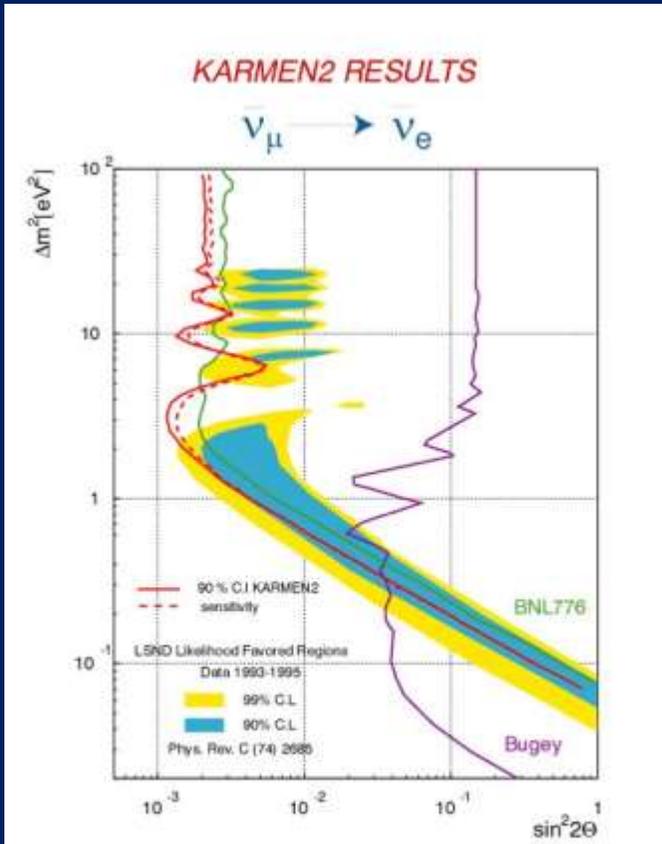
A.A. Aguilar-Arevalo et al. [MiniBooNE Collaboration], arXiv:1805.12028

More hints, experiments, searches

- Investigations, also with new technologies (liquid argon TPC)
- "Reactor anomaly"
- Source experiments
- Disappearance analyses at long baseline experiments
- Dedicated reactors neutrino experiments

Investigating LSND

(Antineutrino, Pion Decay At Rest DAR)



large Δm^2 : $\sin^2(2\theta) < 2.1 \times 10^{-3}$ 90% C.I. Unified approach

sensitivity $\sin^2(2\theta) < 2.4 \times 10^{-3}$

JSNS²
J-PARC MLF

MLF building (bird's view) arXiv:1705.08629

Detector @ 3rd floor (24m from target)

Hg target = Neutron and Neutrino source

50t of Liquid scintillator (including 17t Gd loaded) (4.6m diameter x 4.0m height) 193PMTs

3GeV pulsed proton beam

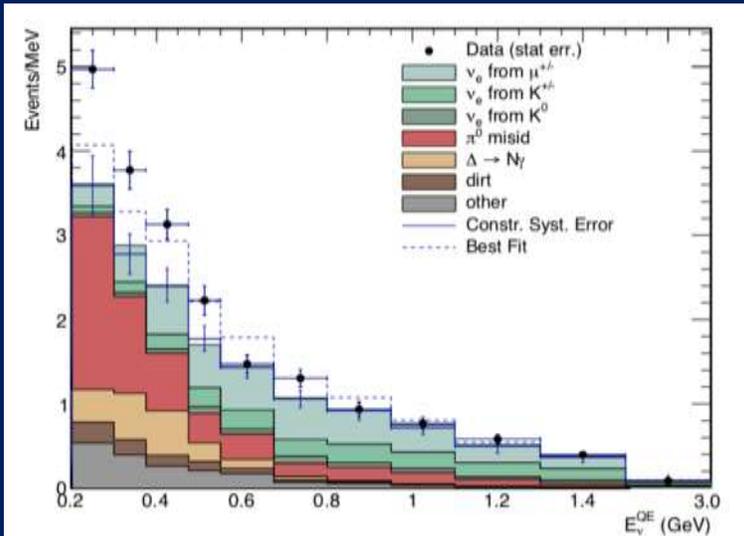
Searching for neutrino oscillation: $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ with baseline of 24m. no new beamline, no new buildings are needed → quick start-up

Carsten Rott - JSNS2 8 ICHEP 2018 July 4-11, 2018

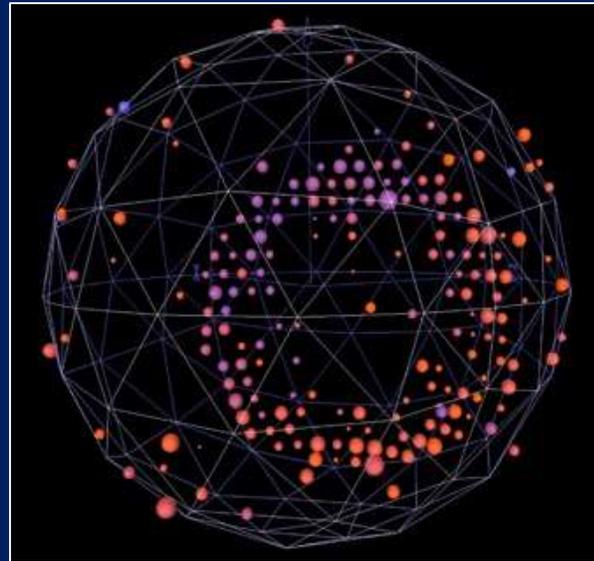
Also DAE δ ALUS and IsoDAR

Operation start in 2019

Explore the MiniBooNE excess of events with a high resolution liquid Argon TPC



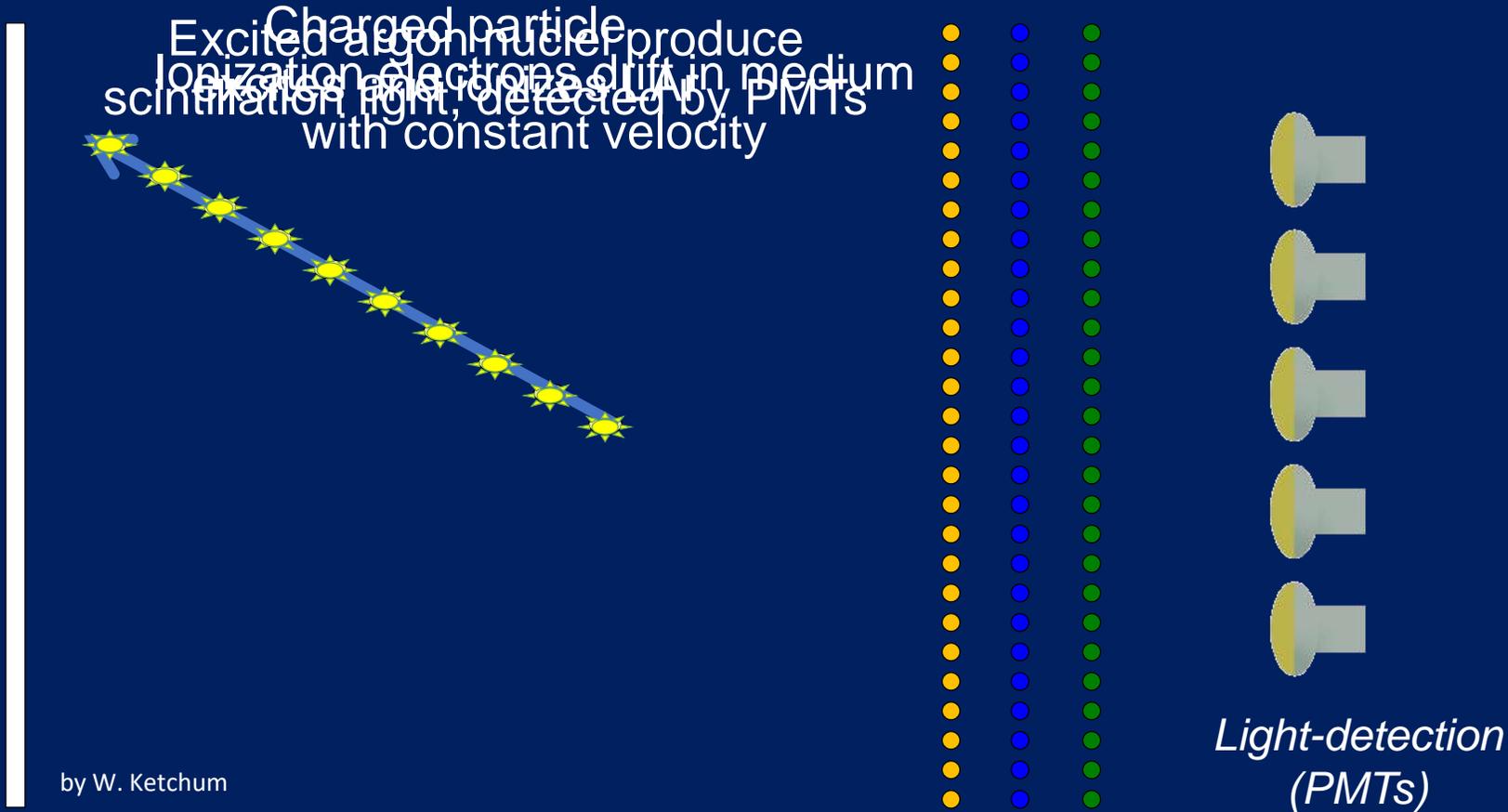
MiniBooNE event



LArTPC

Cathode

Anode wire planes



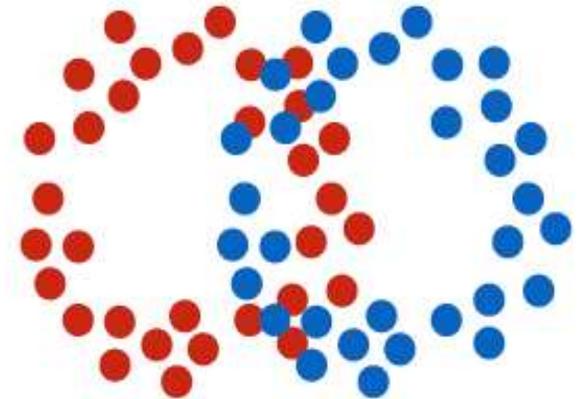
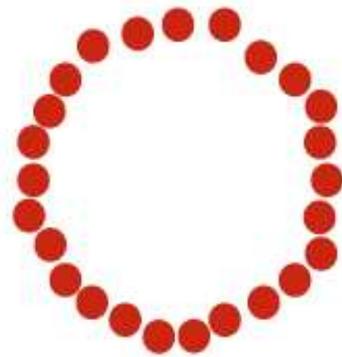
Electron,
Photon

Muon

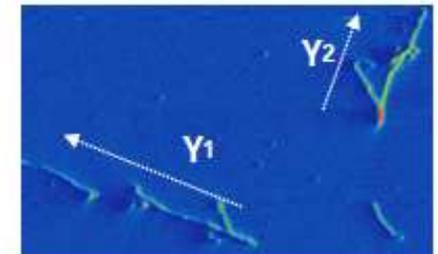
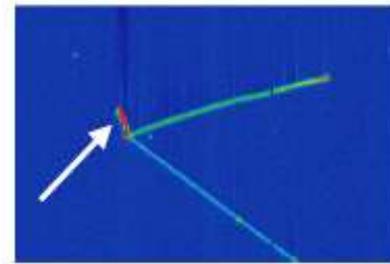
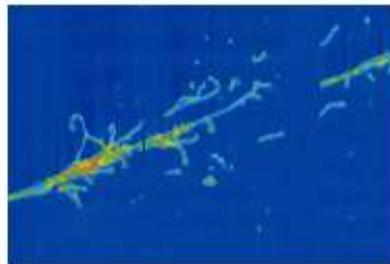
Proton

$\pi^0 \rightarrow \gamma\gamma$

Cherenkov
(MiniBooNE)

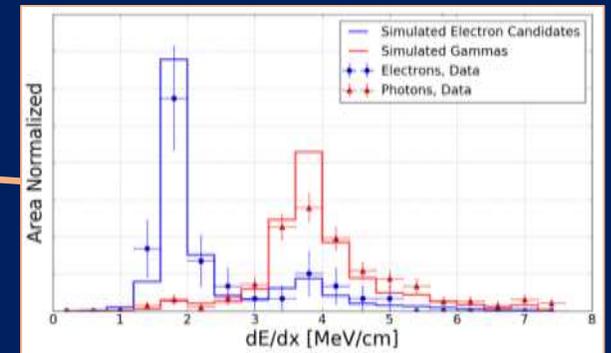
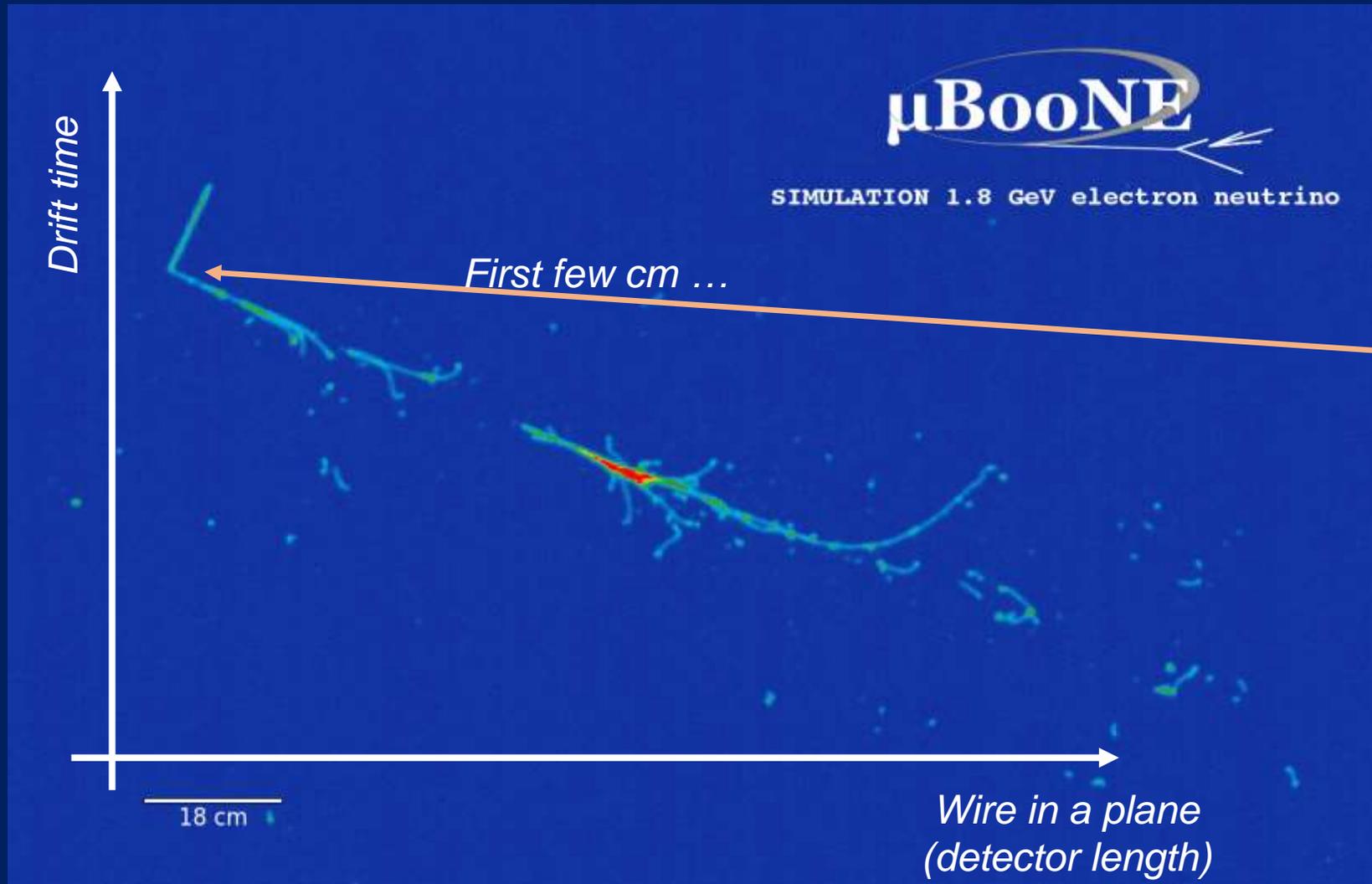


LArTPC
(MicroBooNE)

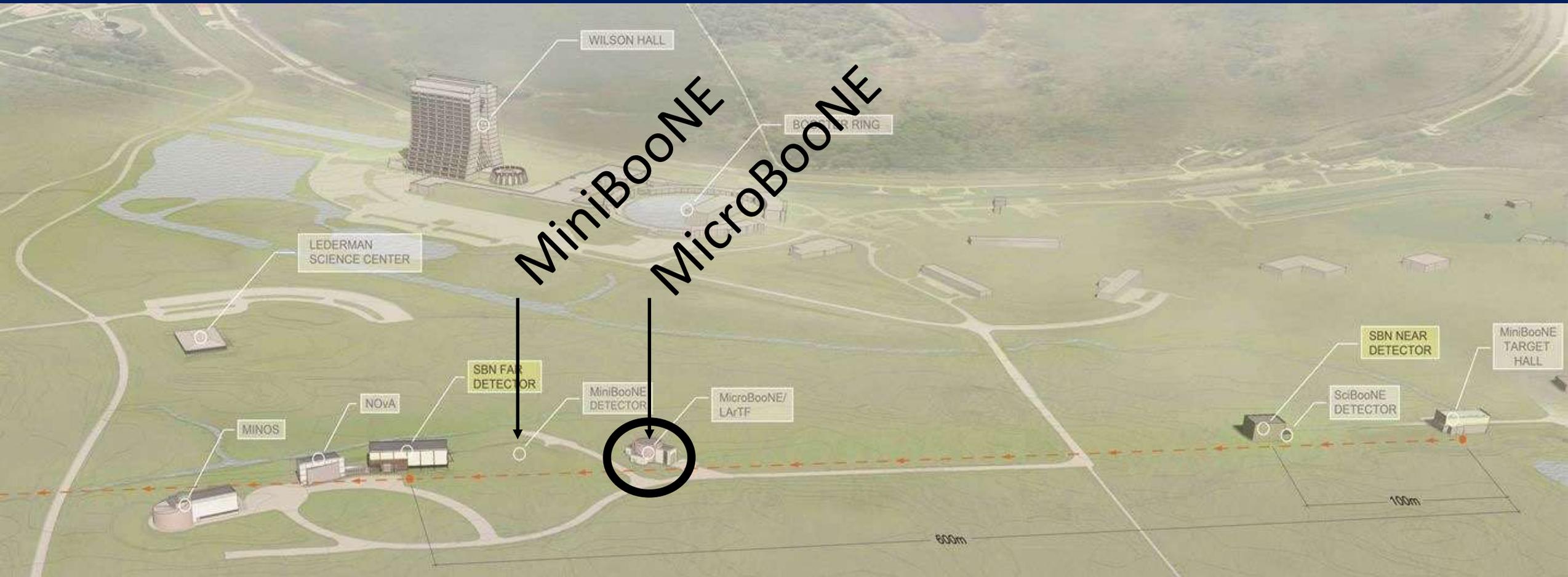


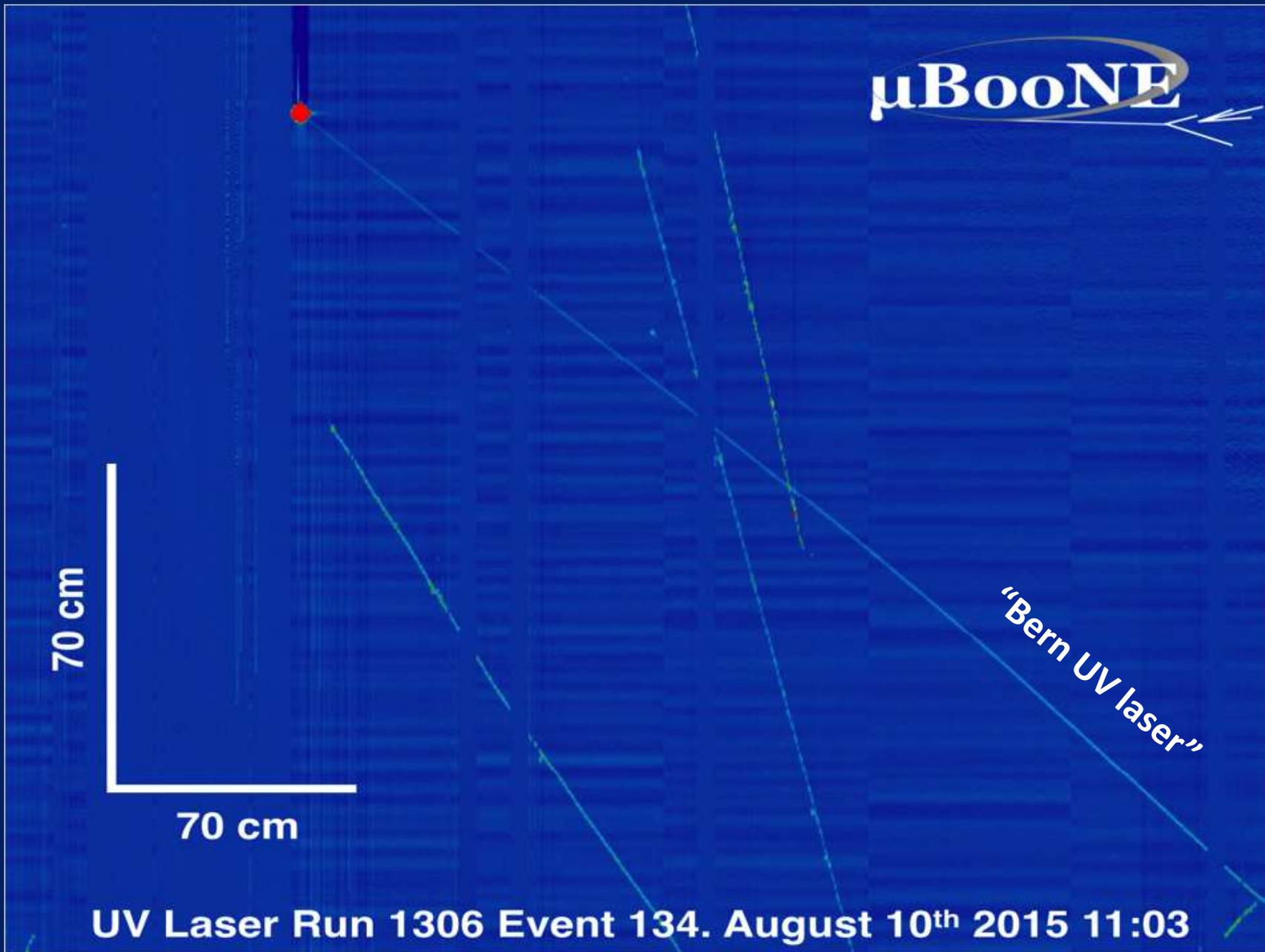
M. Del Tutto

High-Definition 3D imaging detector and calorimeter



ArgoNeut ν_e CC search
Phys. Rev. D 95, 072005 (2017)

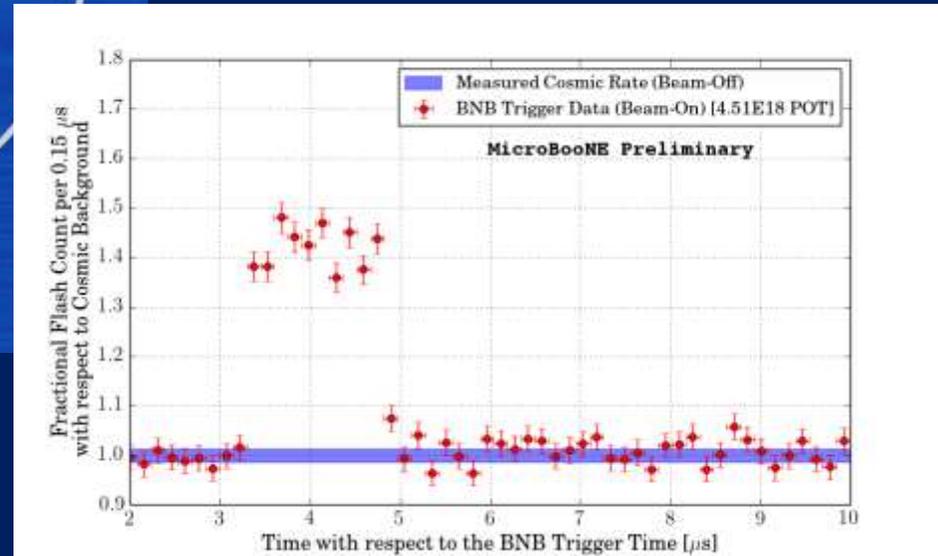
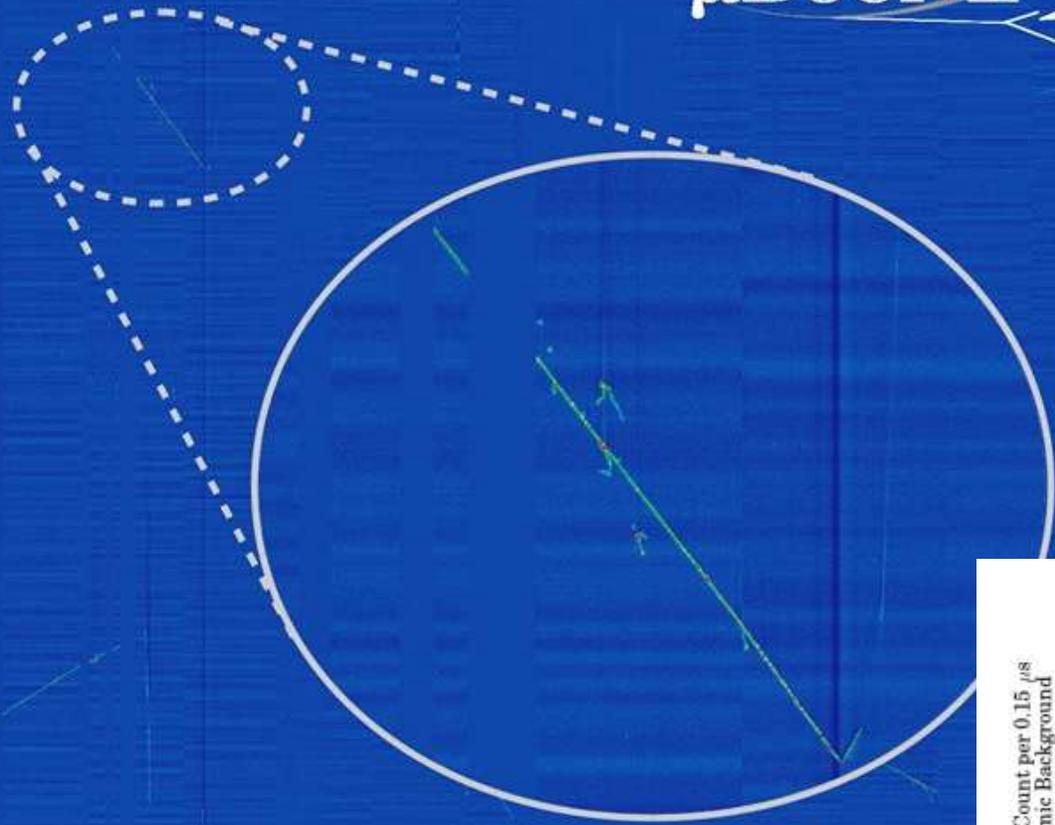




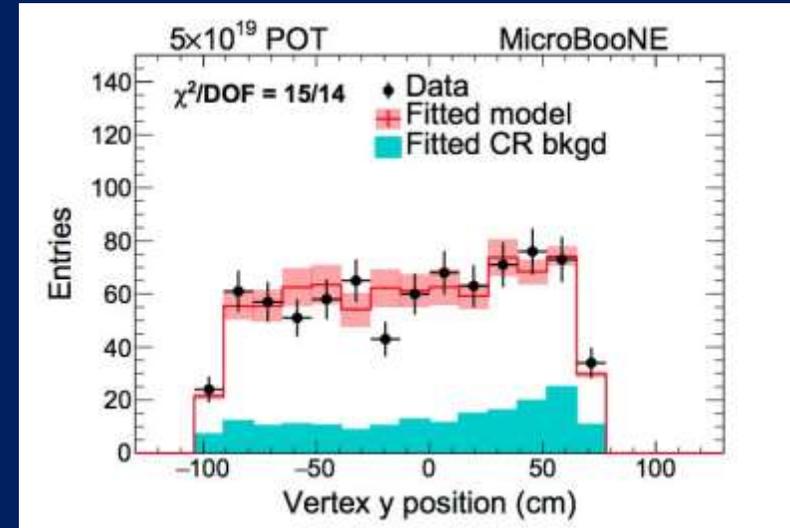
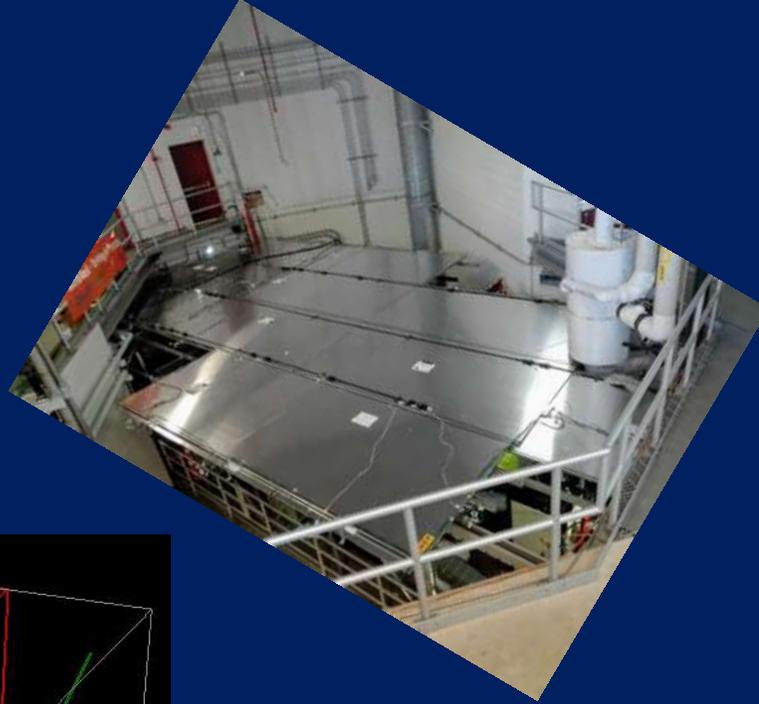
Run 1147 Event 0. August 6th 2015 16:59

μ BooNE

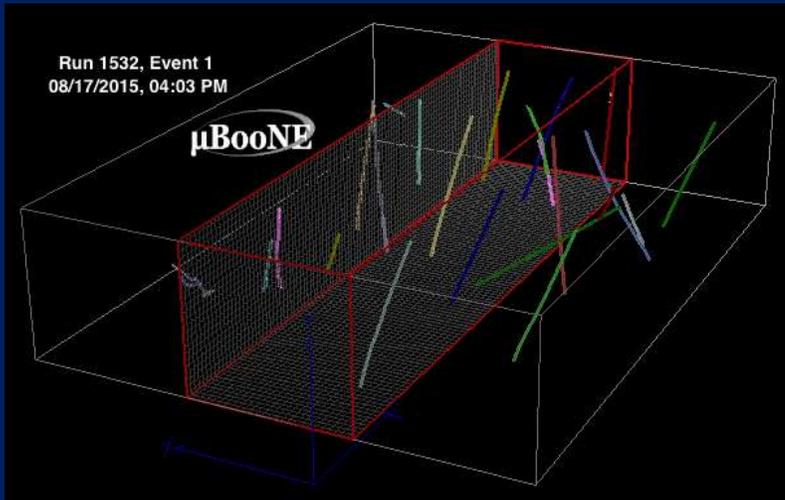
200 cm
200 cm



Towards the low energy oscillation result

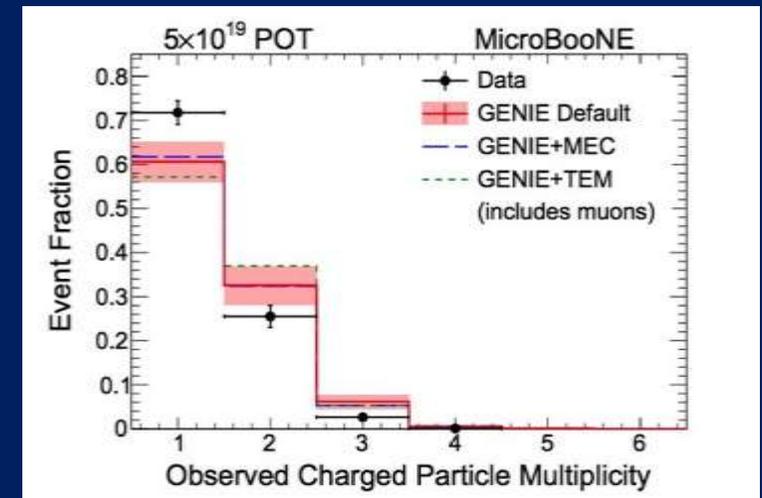


[Eur. Phys. J. C79, 248 \(2019\)](#)



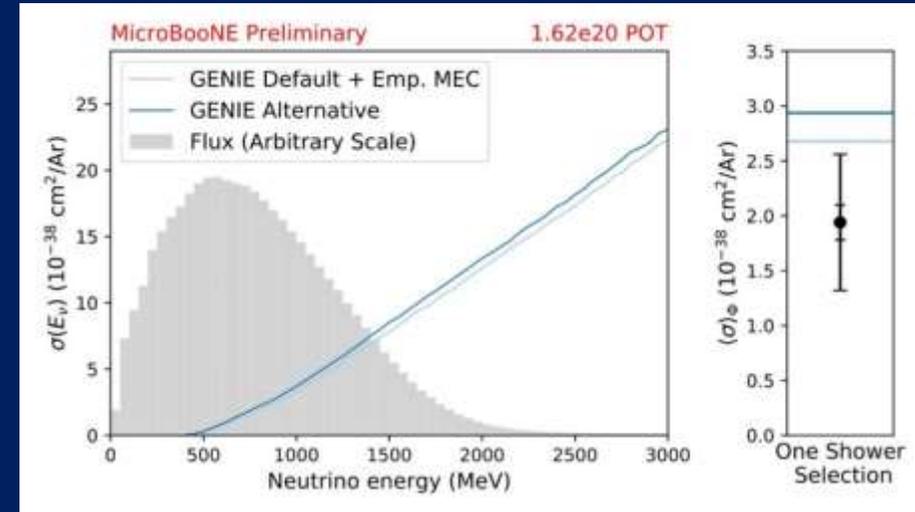
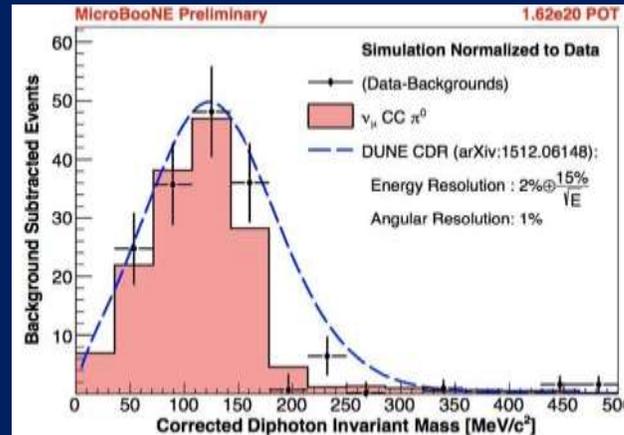
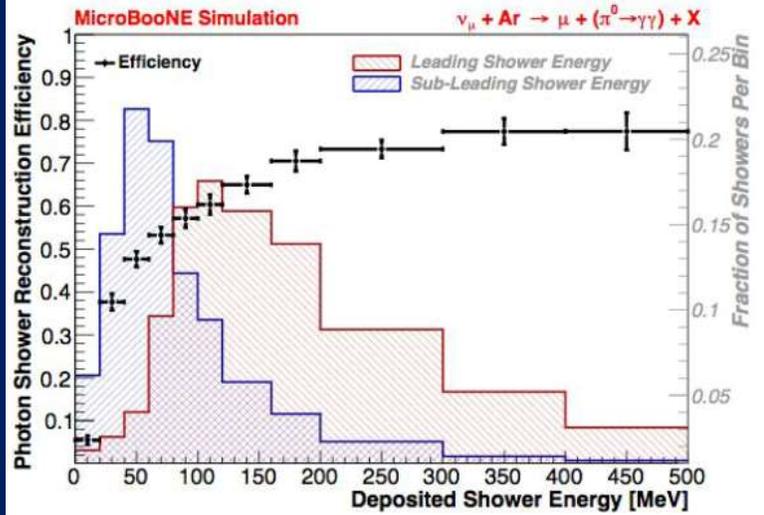
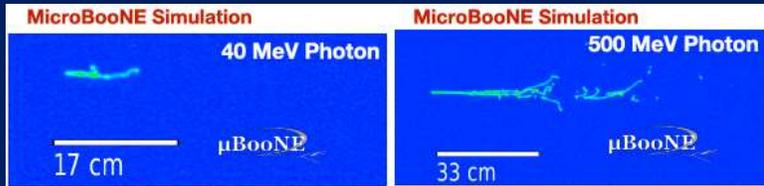
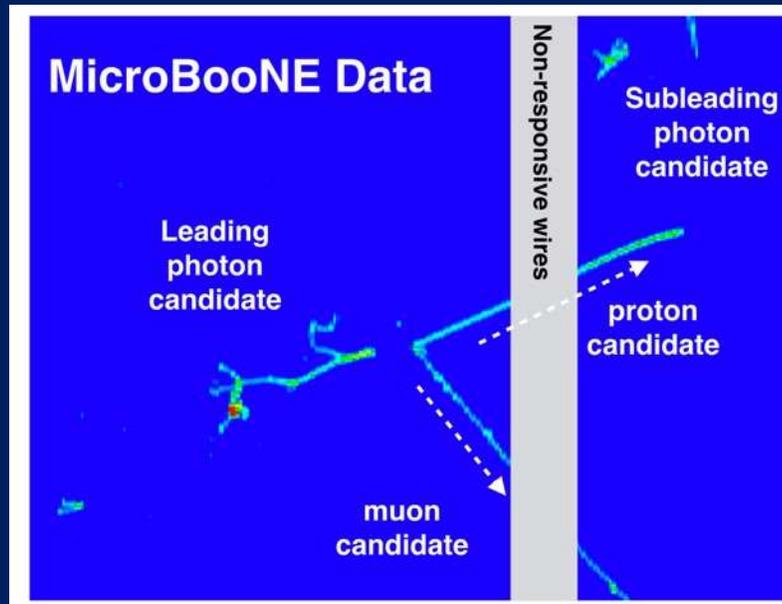
Charged current ν_μ
interactions

Count tracks attached to vertex



e.m. showers

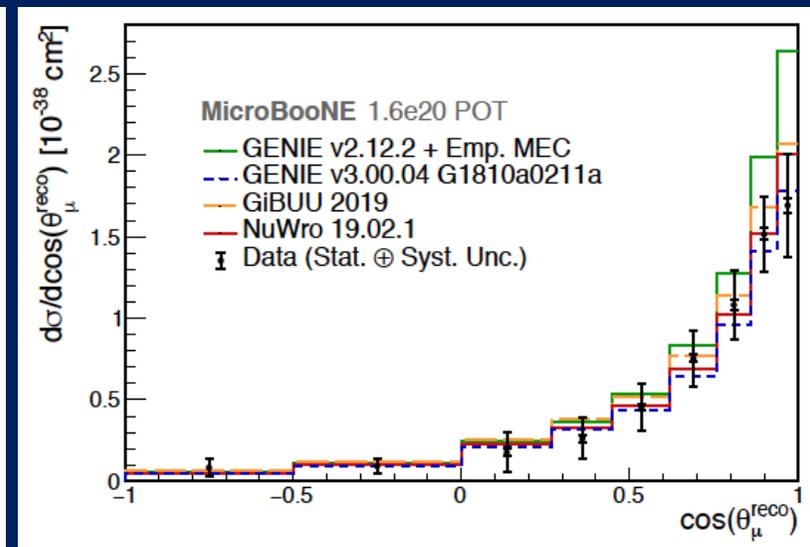
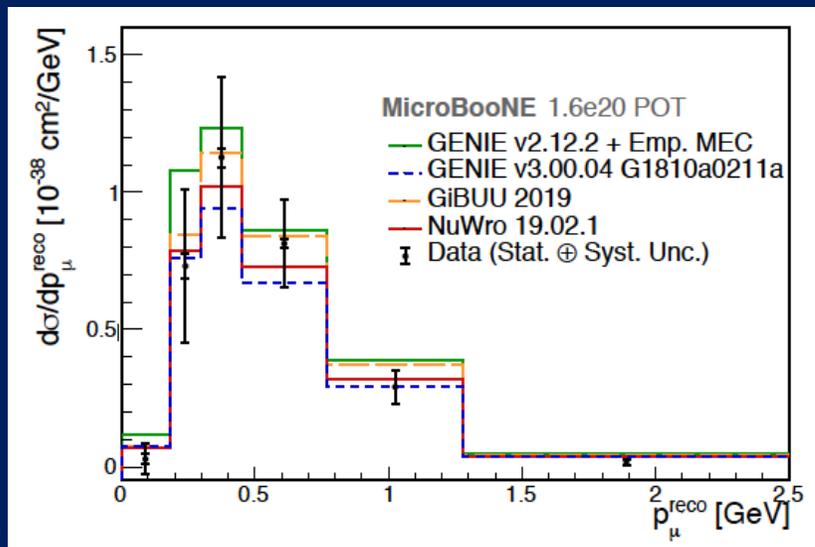
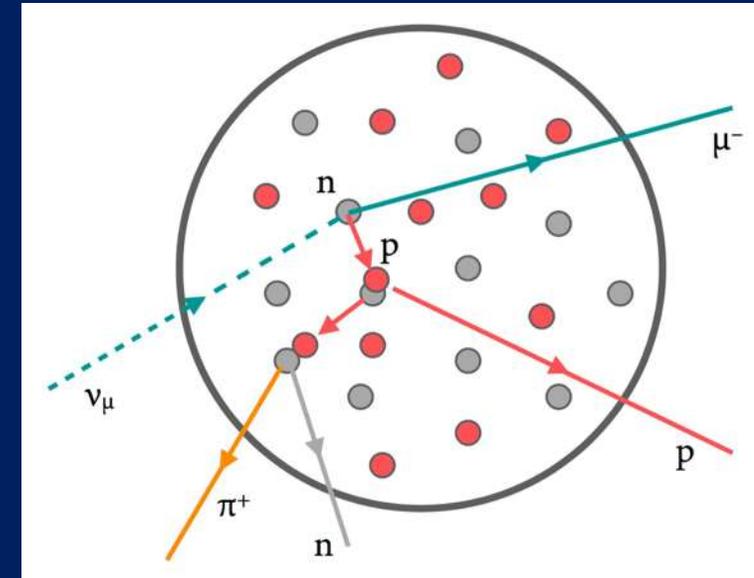
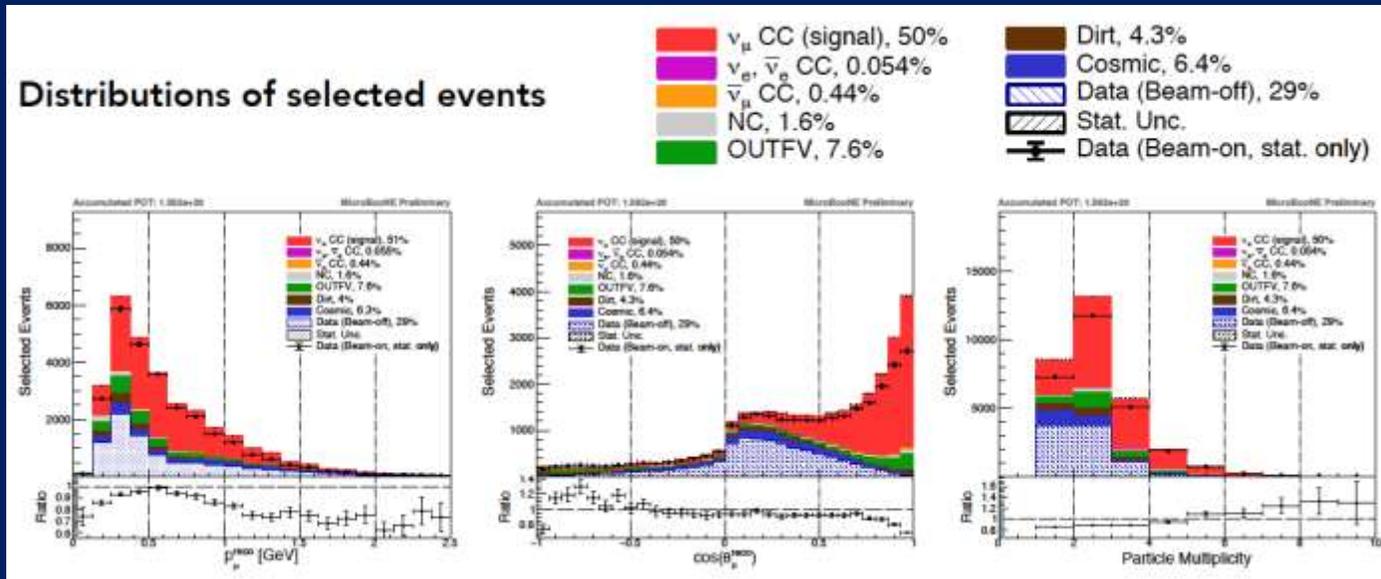
$$\text{CC } \nu_\mu \rightarrow \mu + p + \pi^0$$



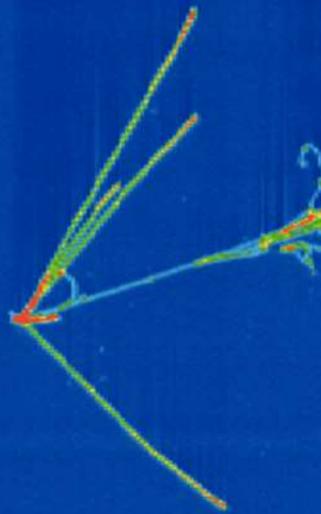
[Phys. Rev. D99, 091102\(R\) \(2019\)](https://arxiv.org/abs/1907.08122)

CC ν_μ double differential cross section

arXiv:1905.09694



μ BooNE



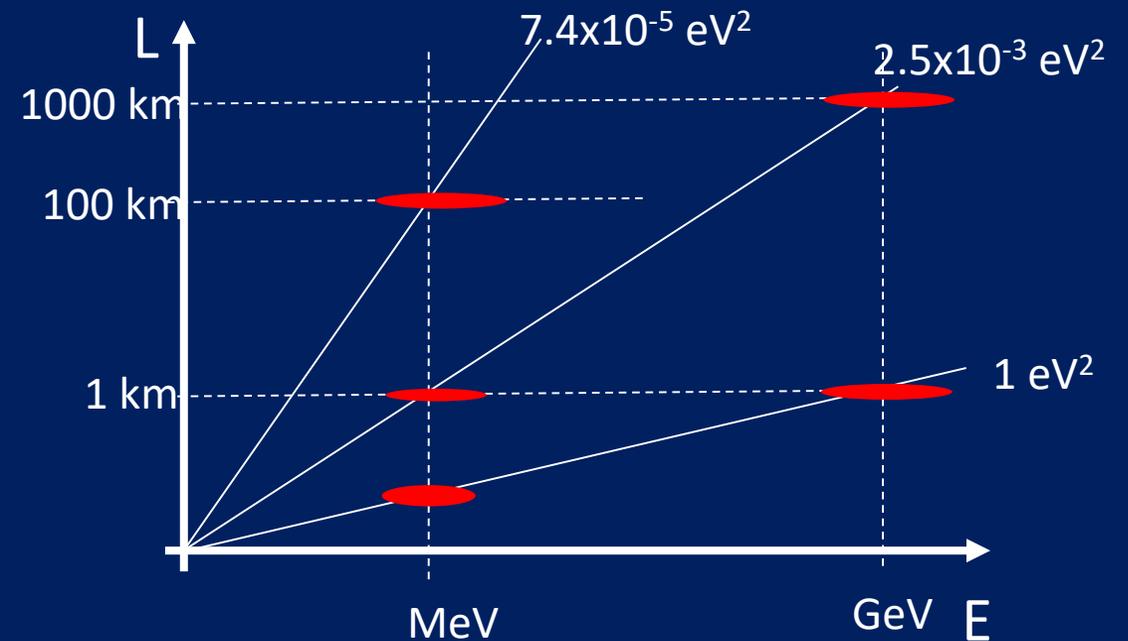
17 cm

$\nu_{\mu} \rightarrow \nu_e ?$

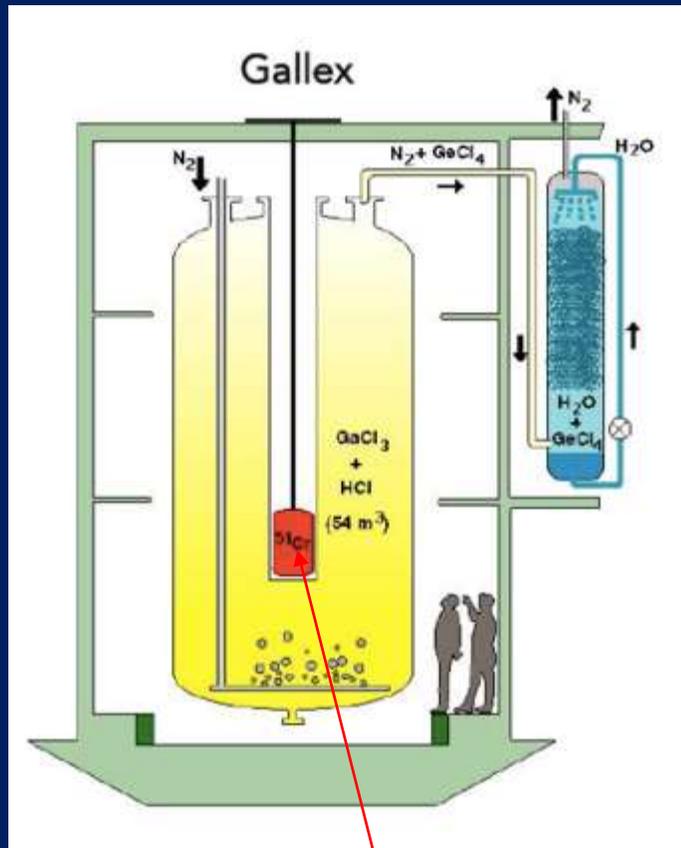
NUMI DATA: RUN 10811, EVENT 2549. APRIL 9, 2017.

Sterile neutrinos: more hints, experiments, searches

- Source experiments
- "Reactor anomaly"
- Dedicated reactors neutrino experiments
- Disappearance analyses at long baseline experiments

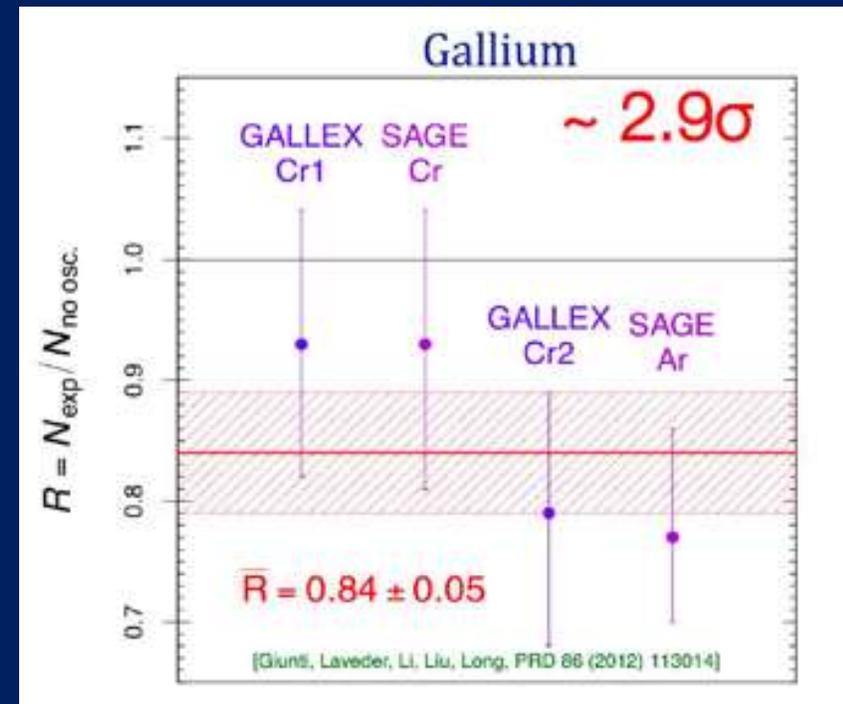


Radioactive source experiments



^{51}Cr

Electron
neutrinos
disappearing

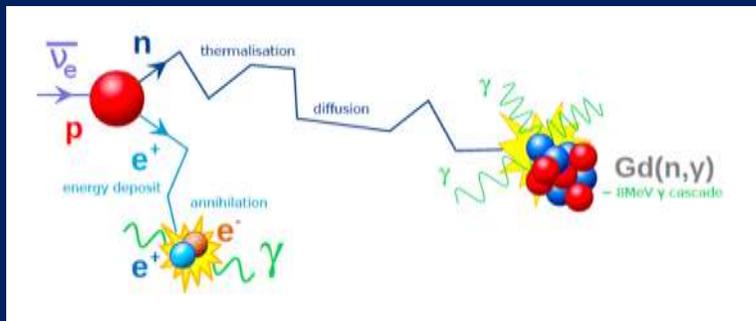
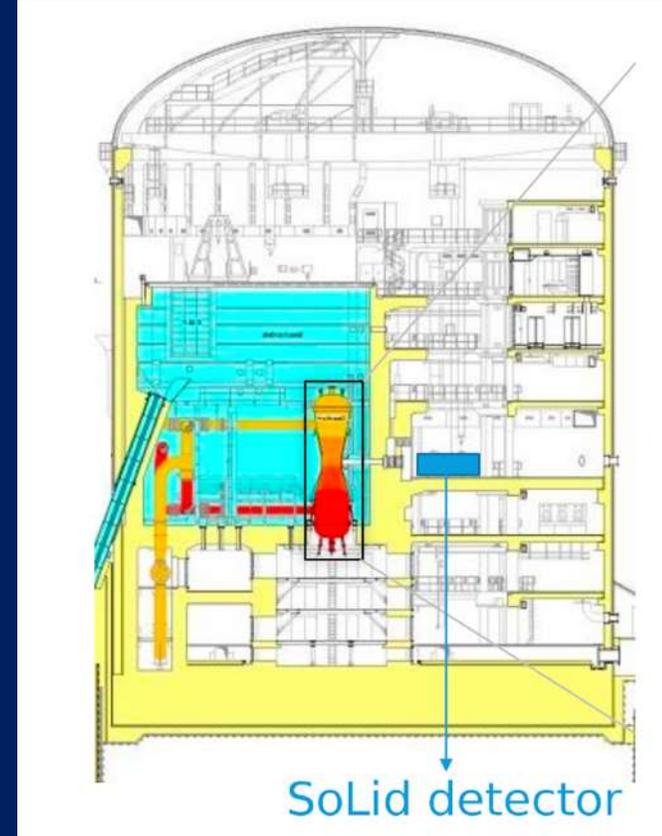
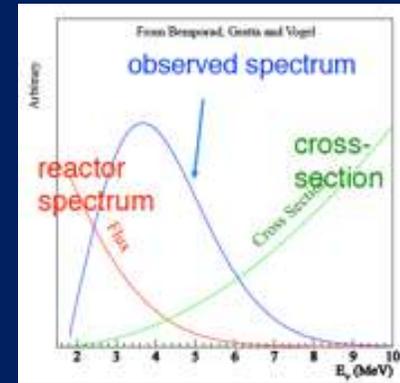
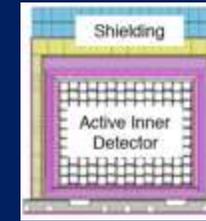


Reactor experiments:

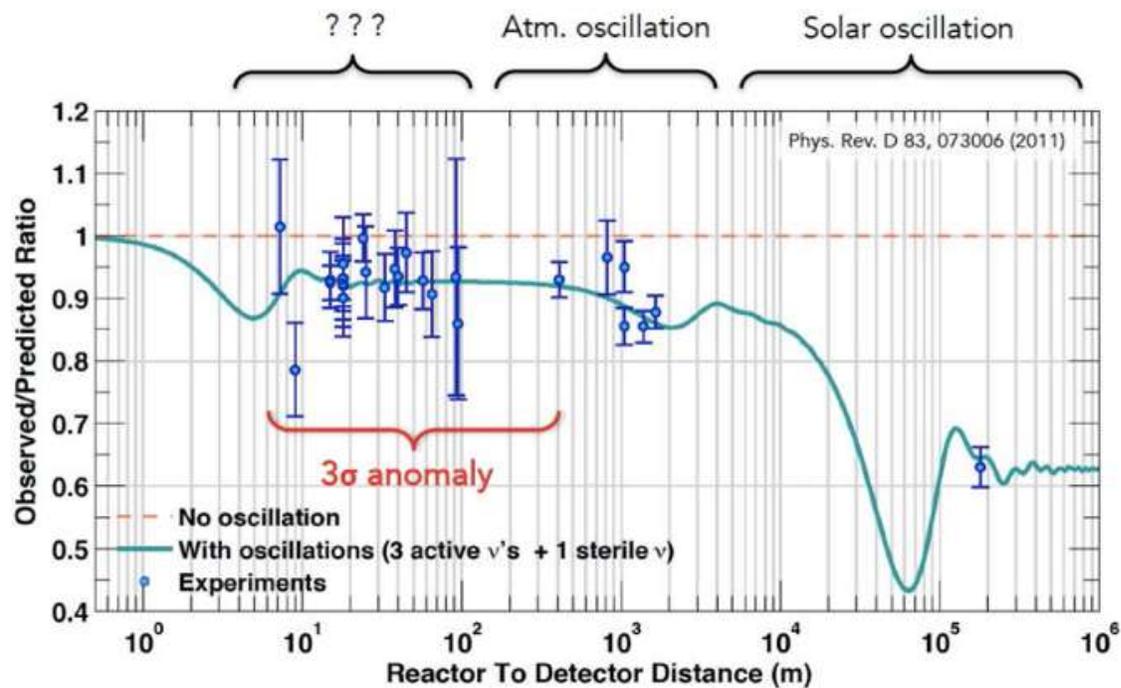
Pure $\bar{\nu}_e$ source produced by fissions in ^{235}U , ^{238}U , ^{239}Pu , ^{241}Pu



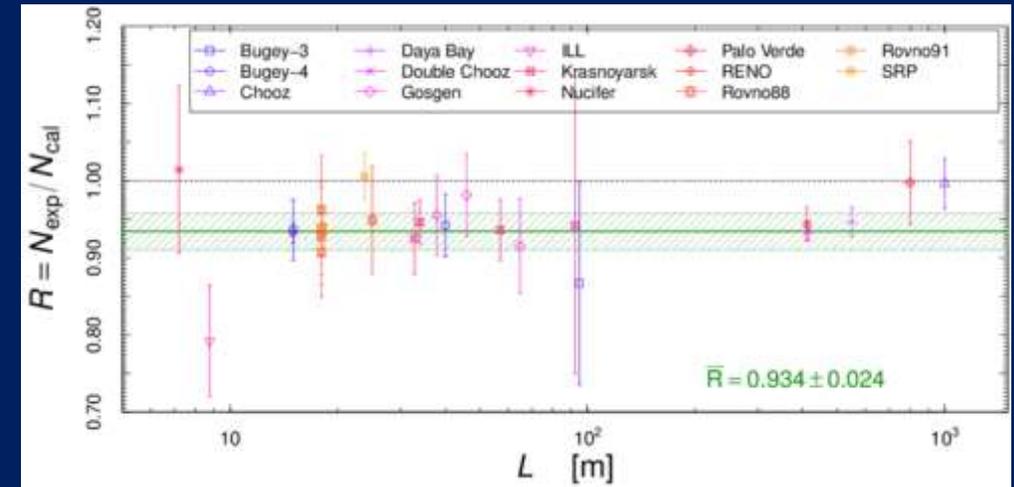
$\bar{\nu}_e$ disappearance
baseline $\sim m$



At short baselines experiments see fewer neutrinos than expected (red line)



G. Mention, M. Fechner, Th. Lasserre, Th. A. Mueller, D.Lhuillier, M. Cribier, and A. Letourneau, Phys. Rev. D 83, 073006 (2011)



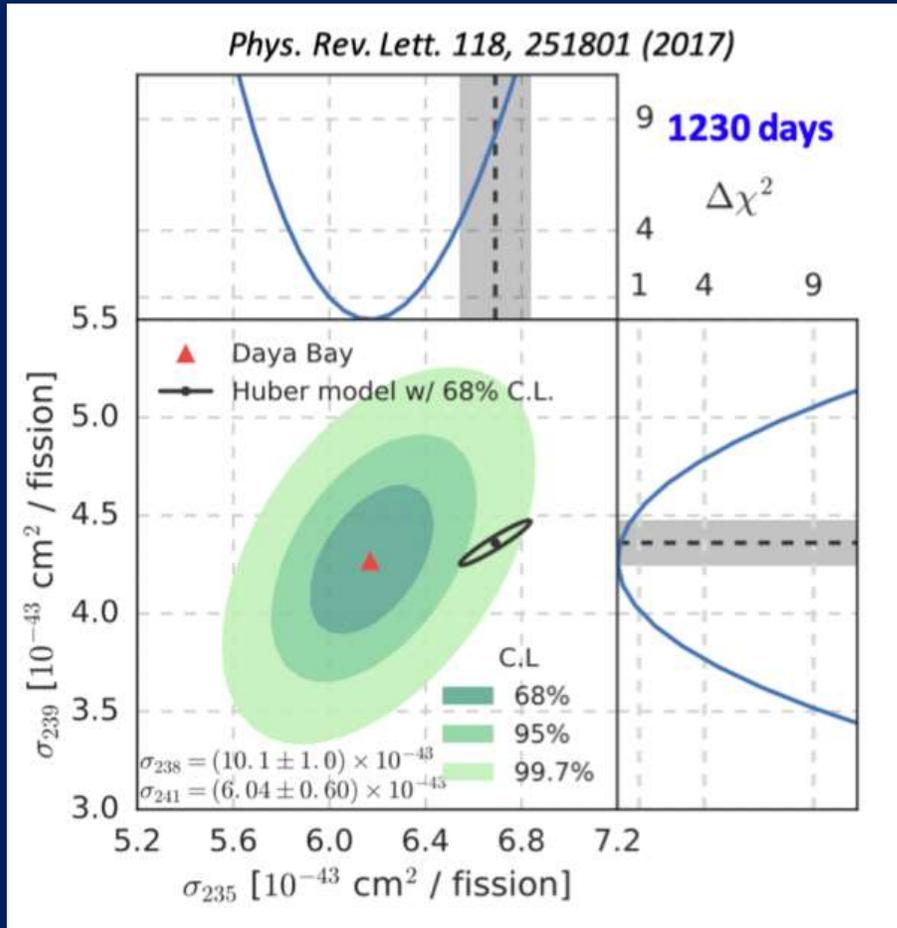
Experiment	Type	Channel	Significance
LSND	DAR	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ CC	3.8σ
MiniBooNE	SBL accelerator	$\nu_\mu \rightarrow \nu_e$ CC	4.5σ
MiniBooNE	SBL accelerator	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ CC	2.8σ
GALLEX/SAGE	Source - e capture	ν_e disappearance	2.8σ
Reactors	Beta-decay	$\bar{\nu}_e$ disappearance	3.0σ

K. N. Abazajian et al. "Light Sterile Neutrinos: A Whitepaper", arXiv:1204.5379 [hep-ph], (2012)

Taken individually, each anomaly is not significant enough to be convincing....

most commonly interpreted as hint for one or more new “sterile” neutrino.

The reactor anomaly can be explained

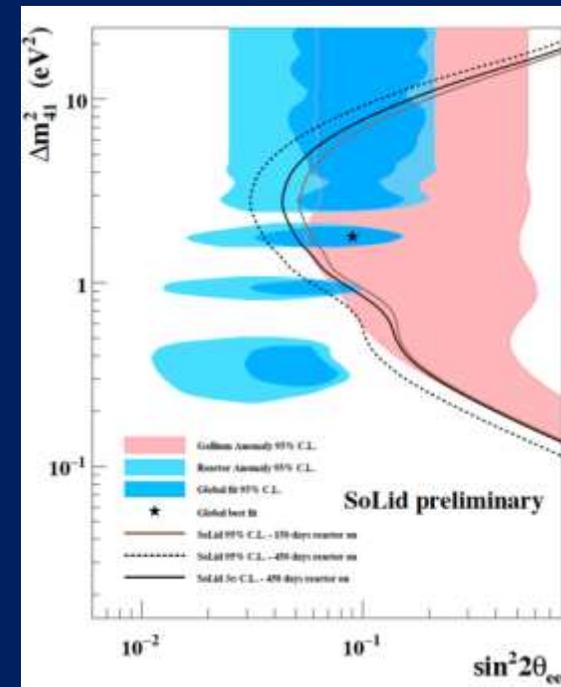
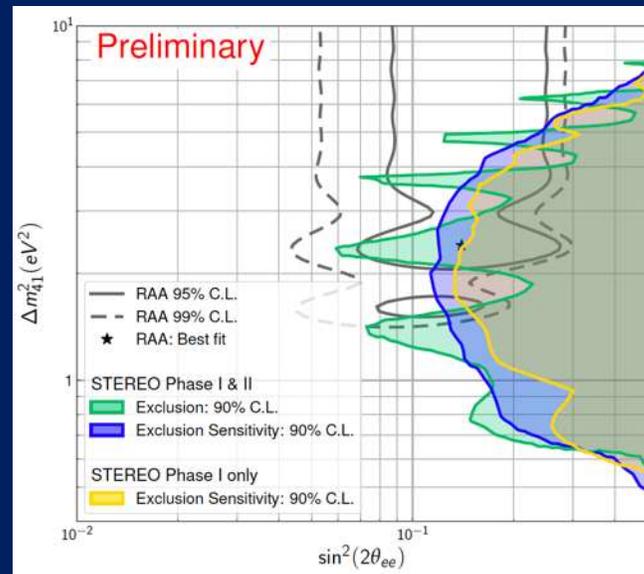
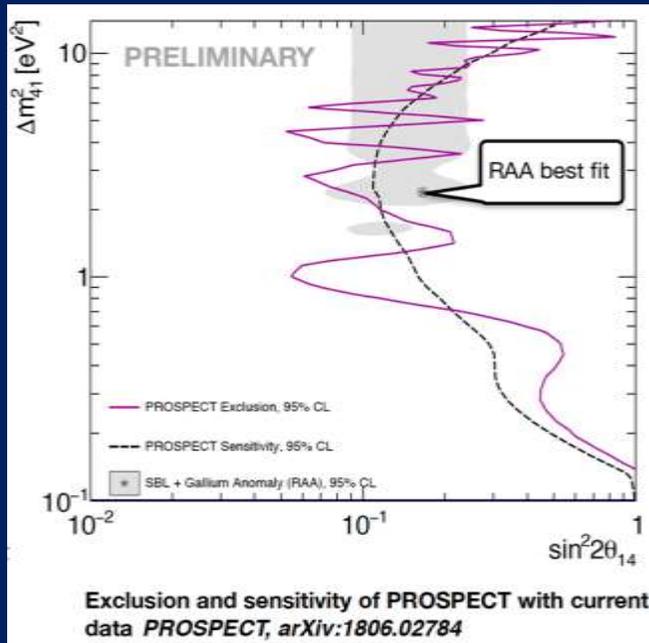
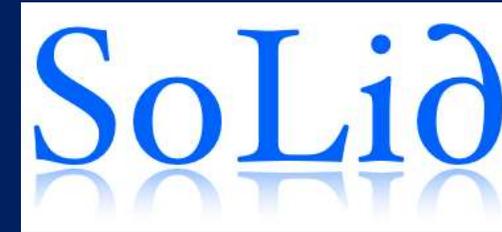


Daya Bay (and RENO) results suggest that the neutrino flux from ^{235}U ~ 3 sigma below what expected from fission models so far

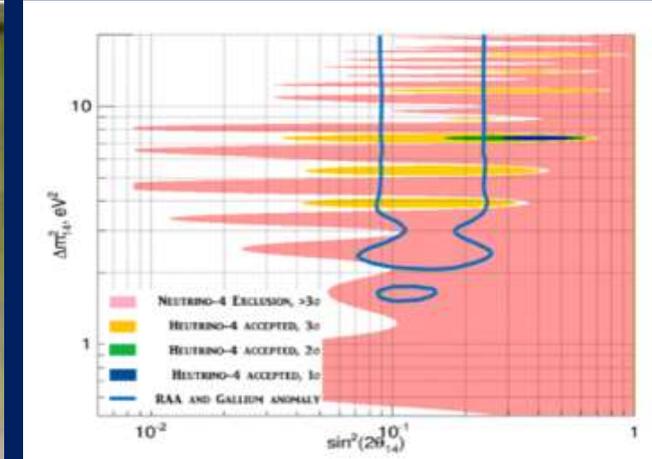
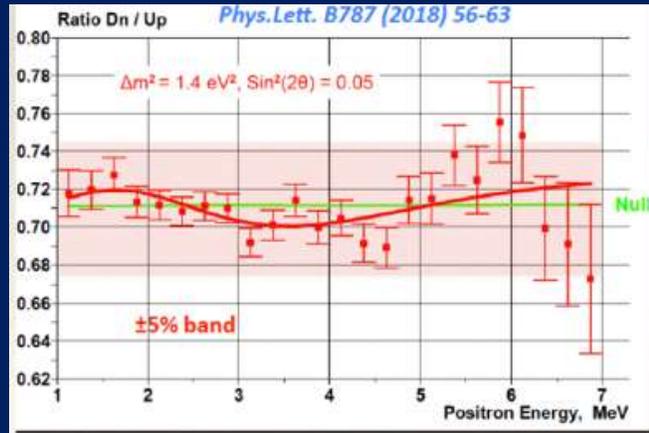
→ Need experiments at Highly Enriched Uranium reactors (20-90% ^{235}U), i.e. research reactors to thoroughly test this



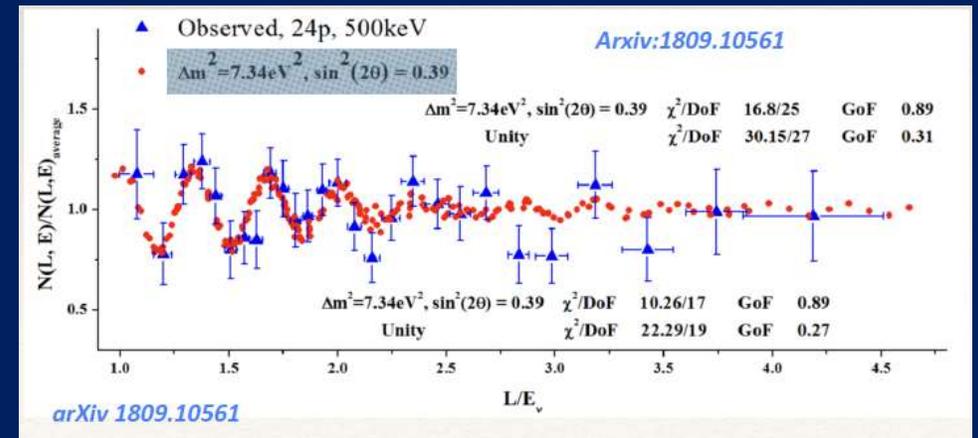
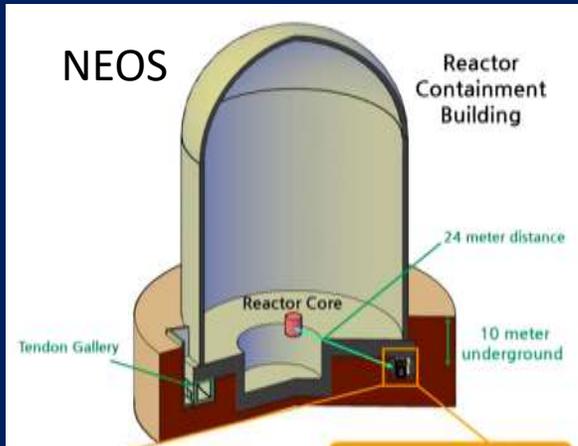
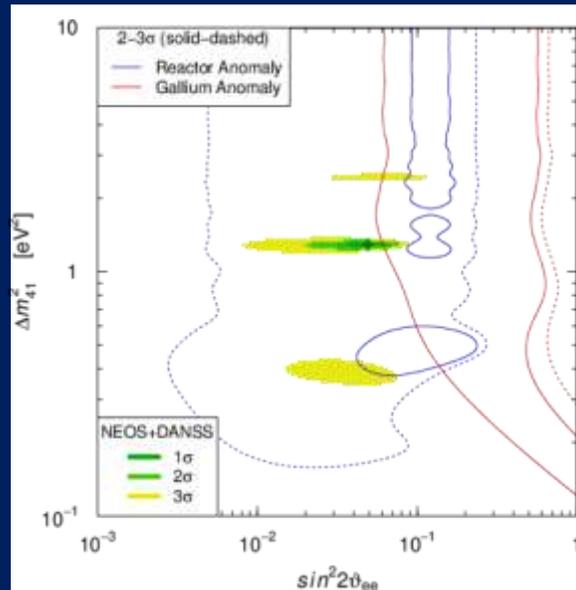
Compact reactors Highly enriched ^{235}U



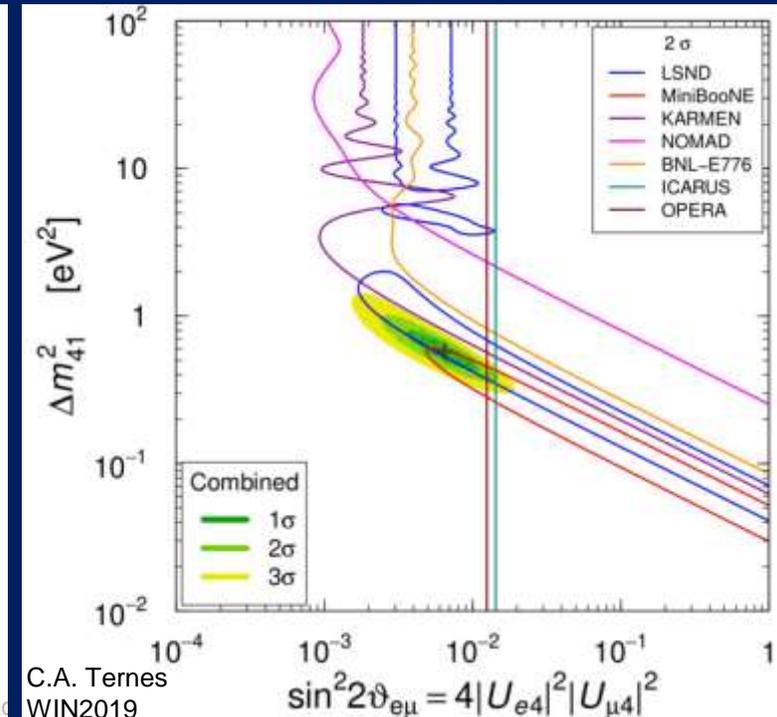
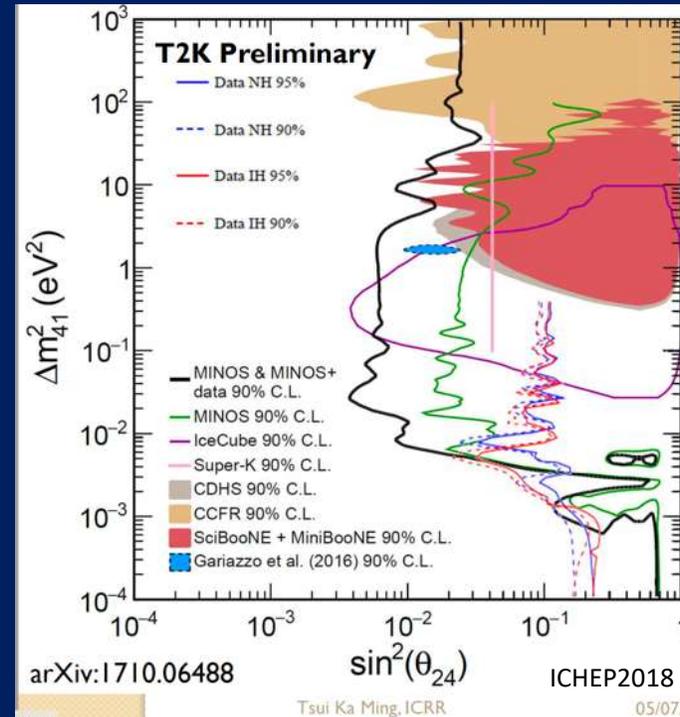
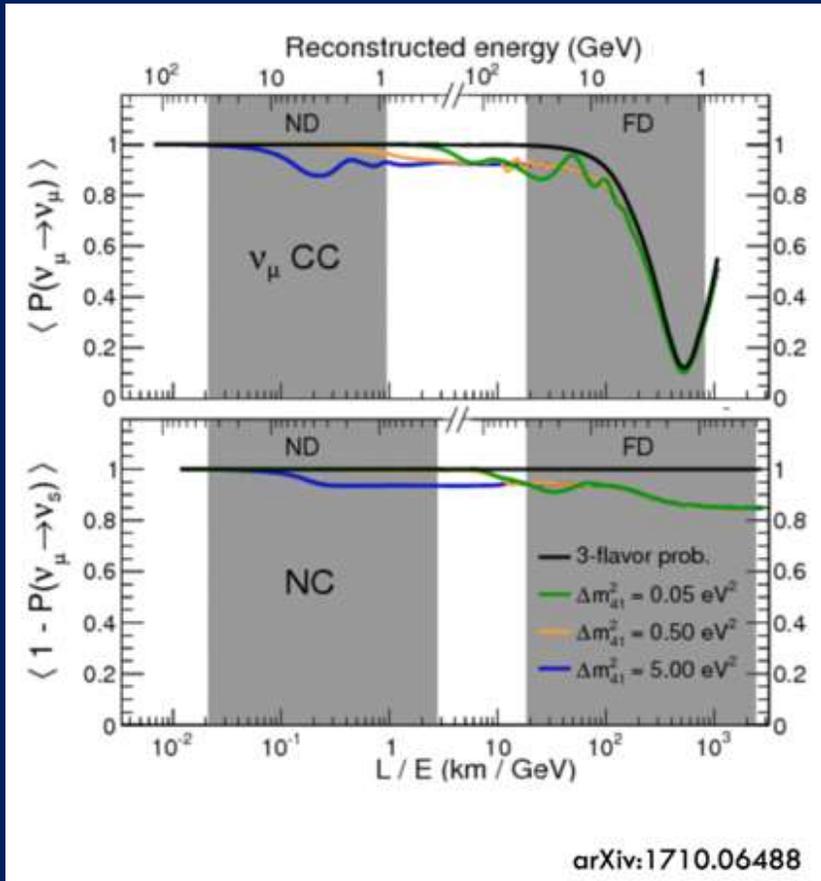
More reactor sterile neutrino searches



Neutrino-4



Sterile Neutrinos and long baseline experiments

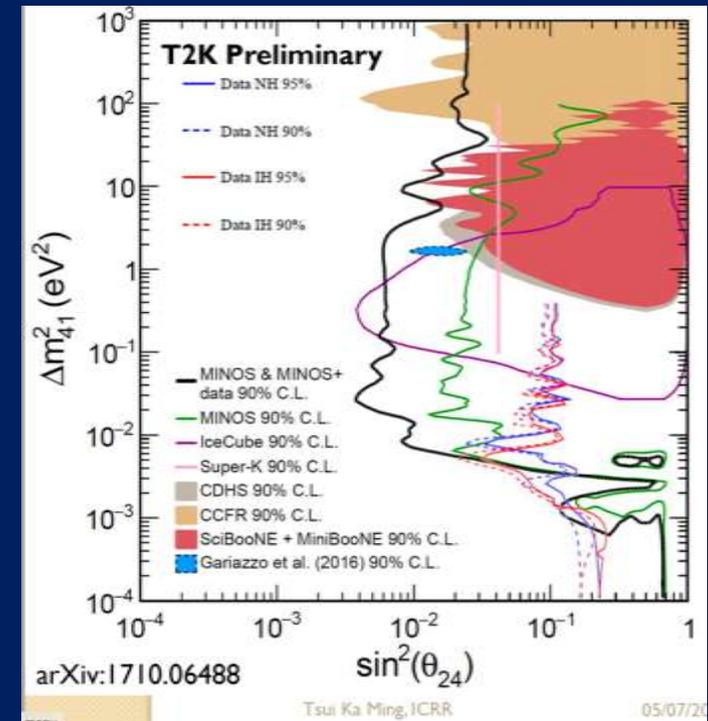
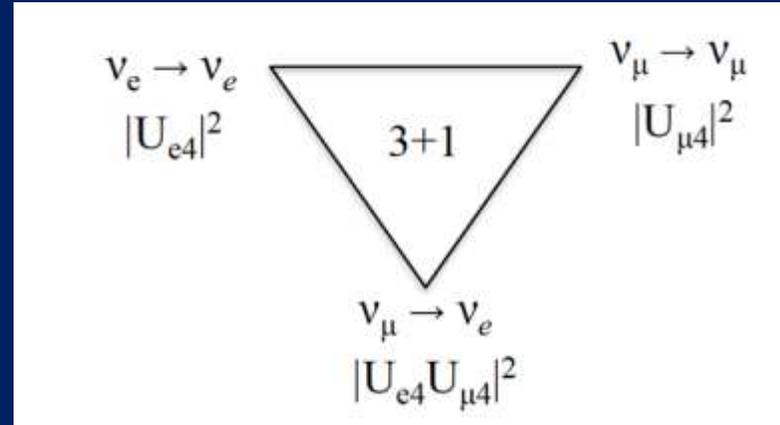
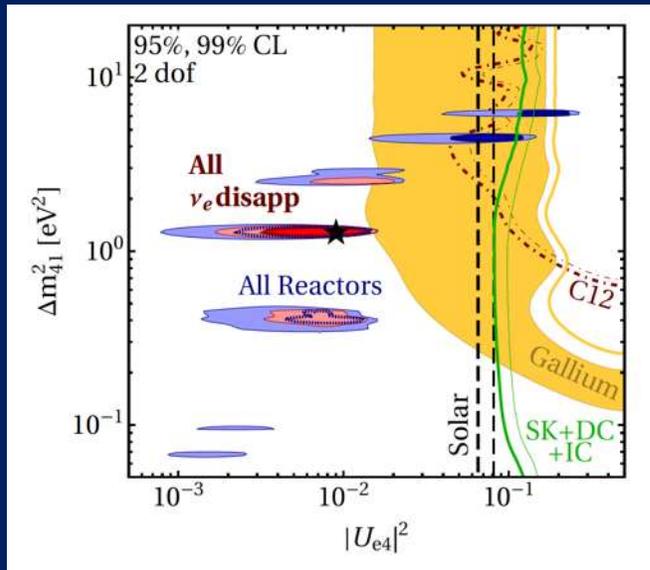


$$\nu_\mu \rightarrow \nu_\mu : |U_{\mu 4}|^2 = \sin^2 \theta_{24} \cos^2 \theta_{14}$$

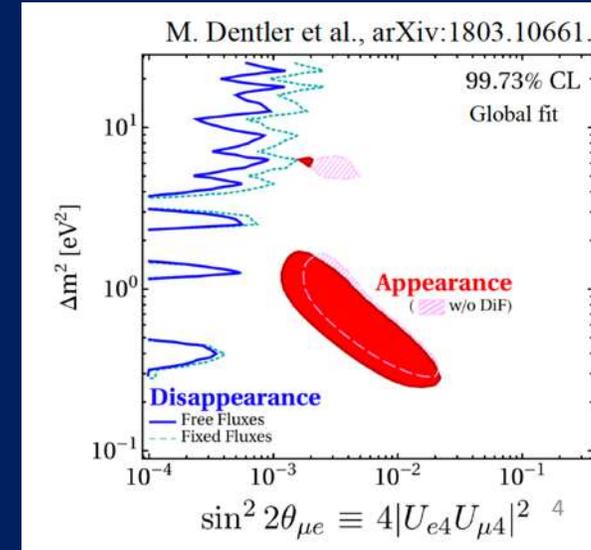
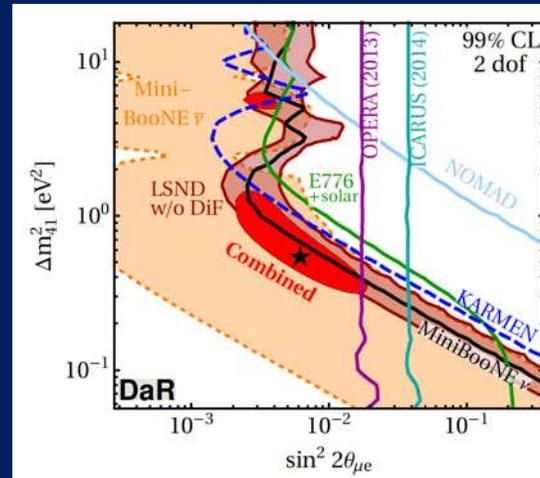
Different oscillation signatures
in near and far detectors

Combination and tensions

All ν_e disappearance arXiv:1803.10661

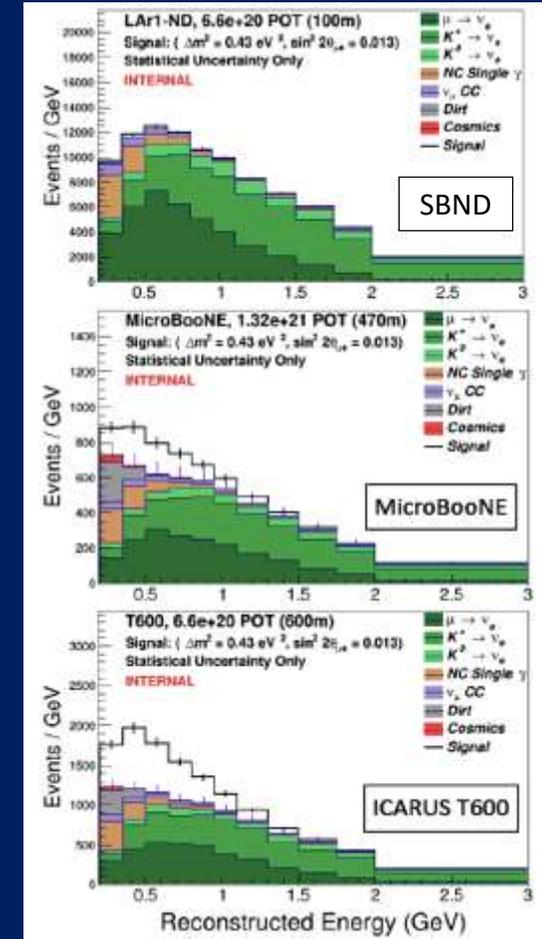
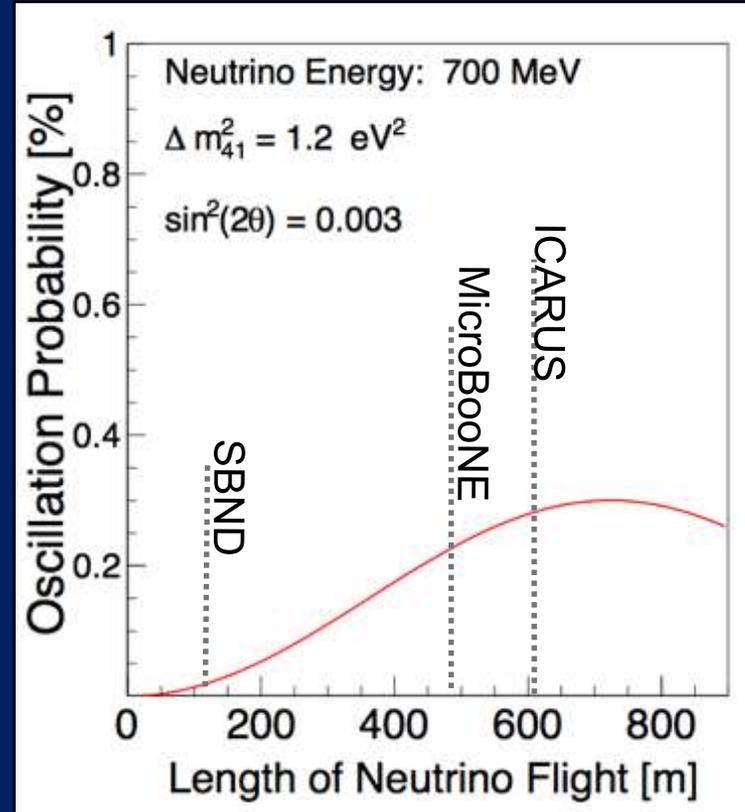


ν_e appearance arXiv:1803.10661

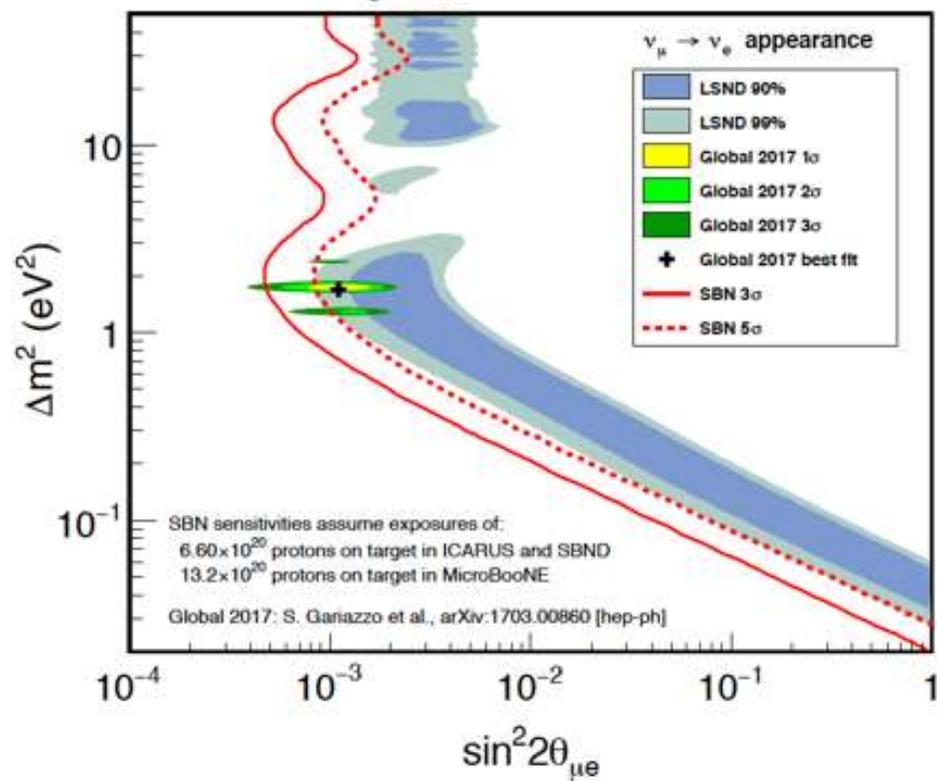




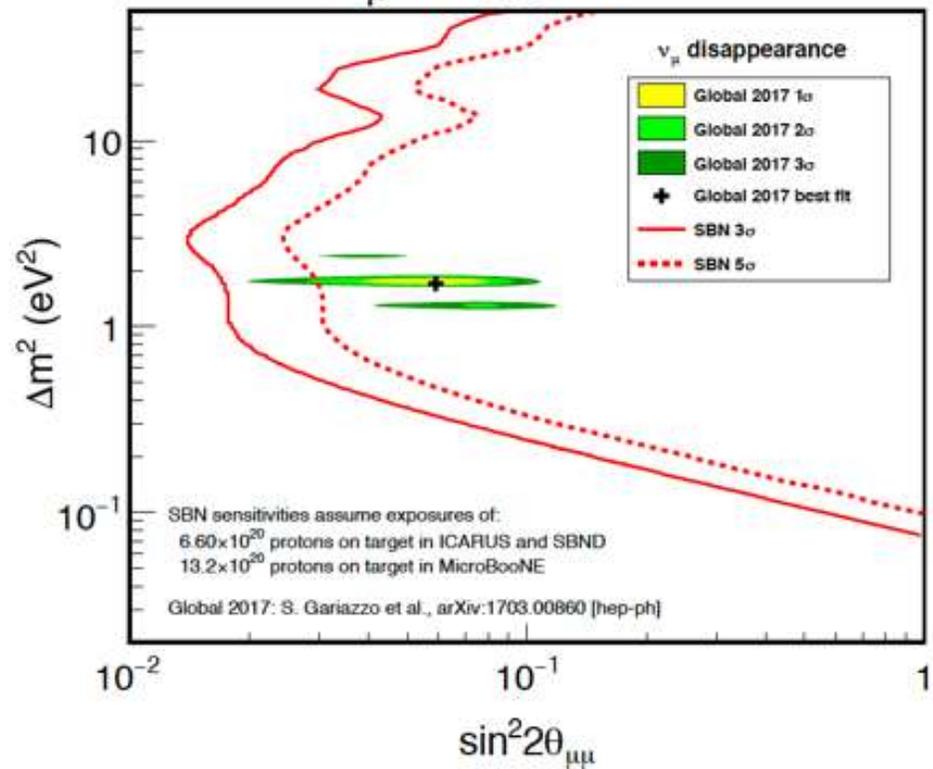
Three-detector setup with LArTPCs



ν_e appearance

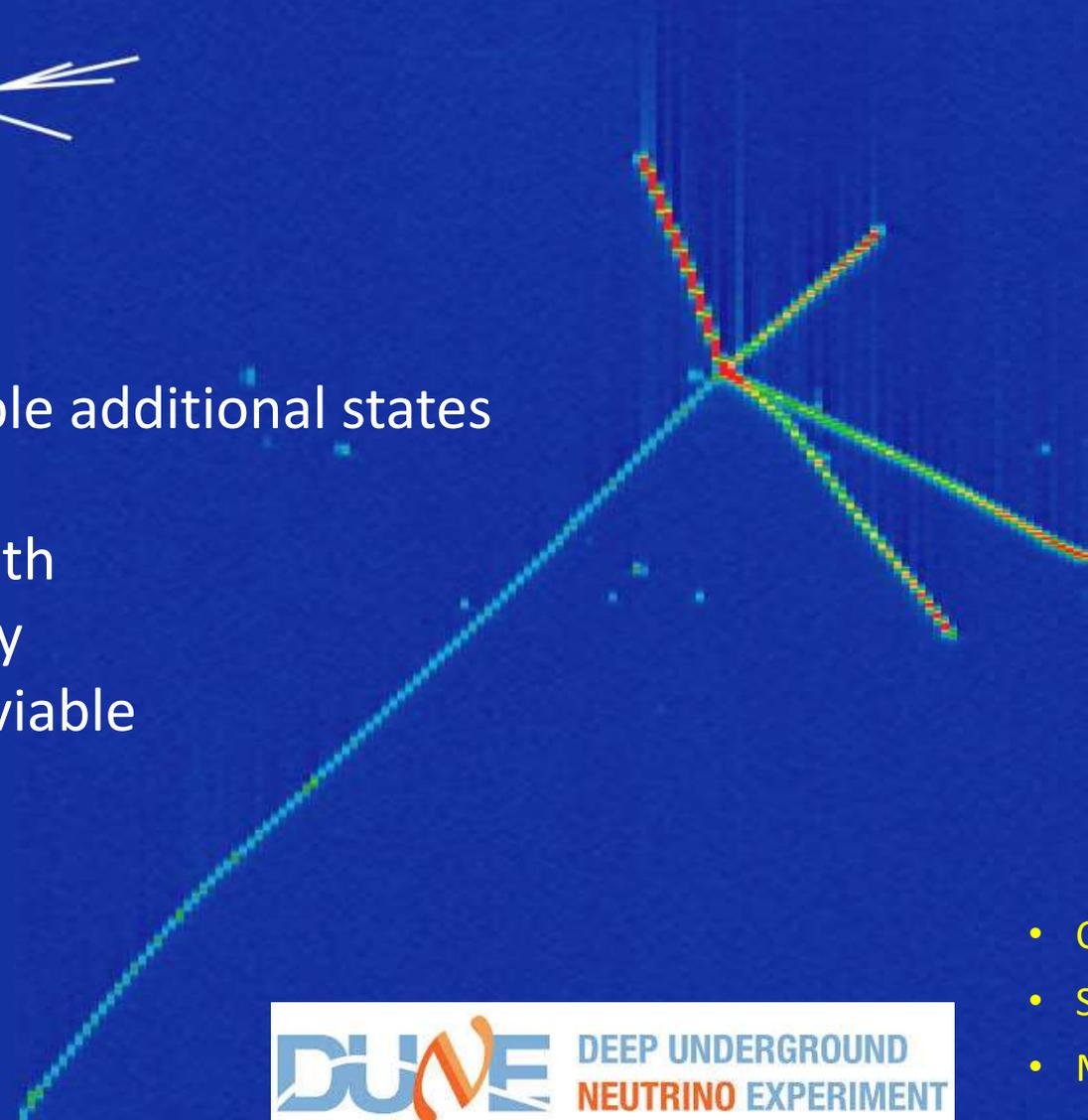


ν_μ disappearance



μ BooNE

- Neutrinos are interesting (known physics BSM)
- Anomalies point to possible additional states (Sterile neutrinos)
- Follow up experiments with new technology underway
- Sterile neutrinos are still viable



10 cm



- CP violation in the lepton sector
- Sign of Δm^2_{31} , "Mass hierarchy"
- Majorana or Dirac neutrinos
- Absolute neutrino masses

+ + + + + + +

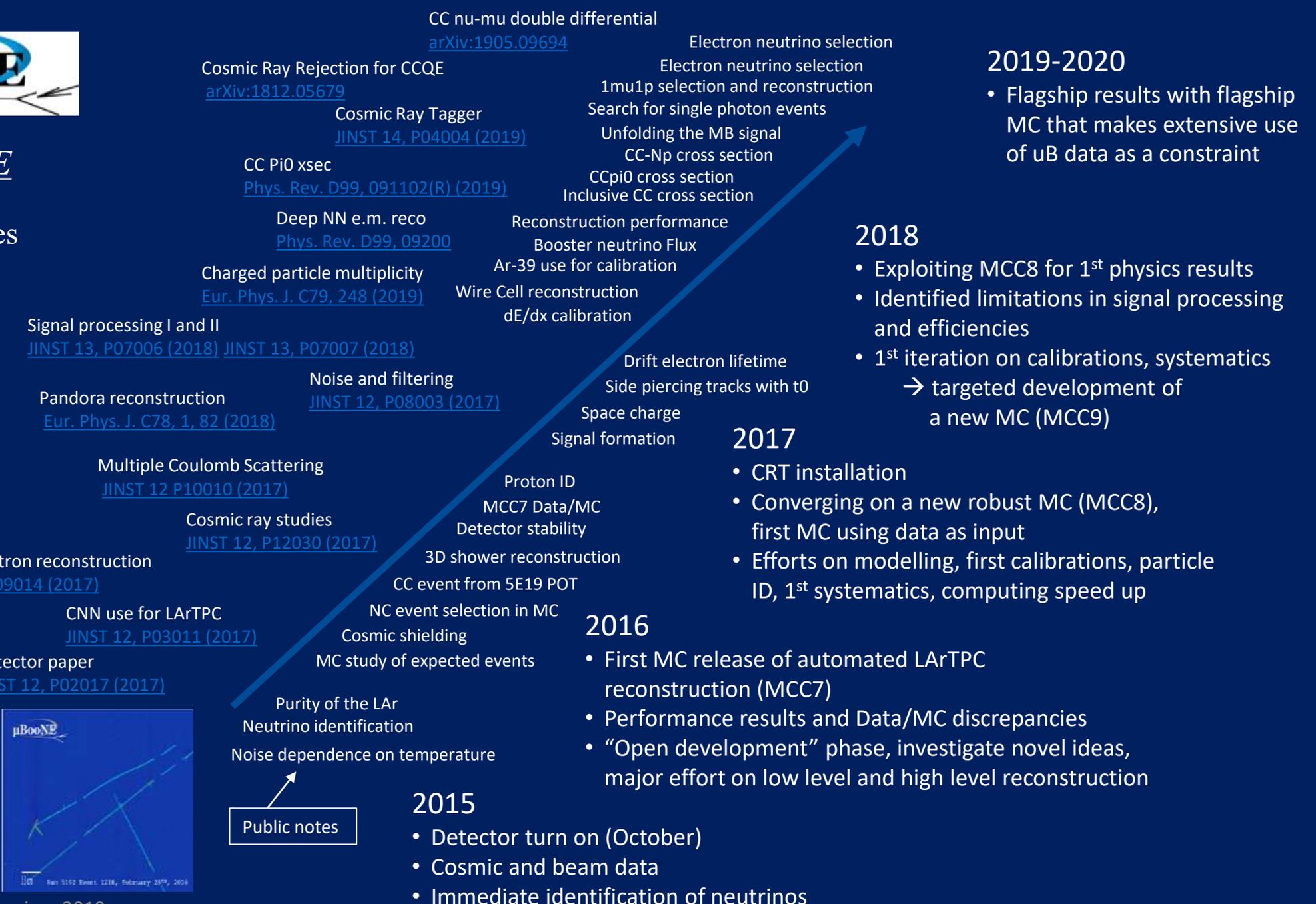
BACKUP

+ + + + + + +

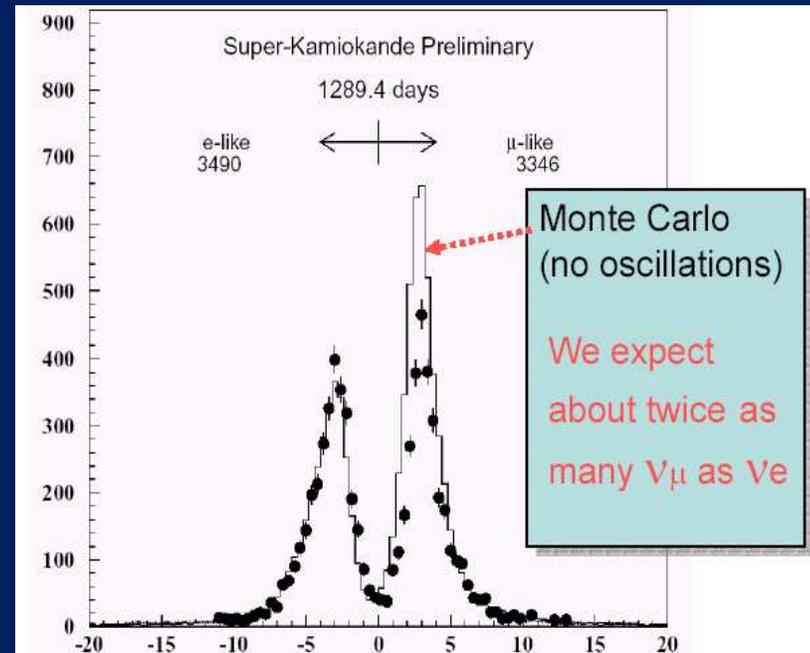
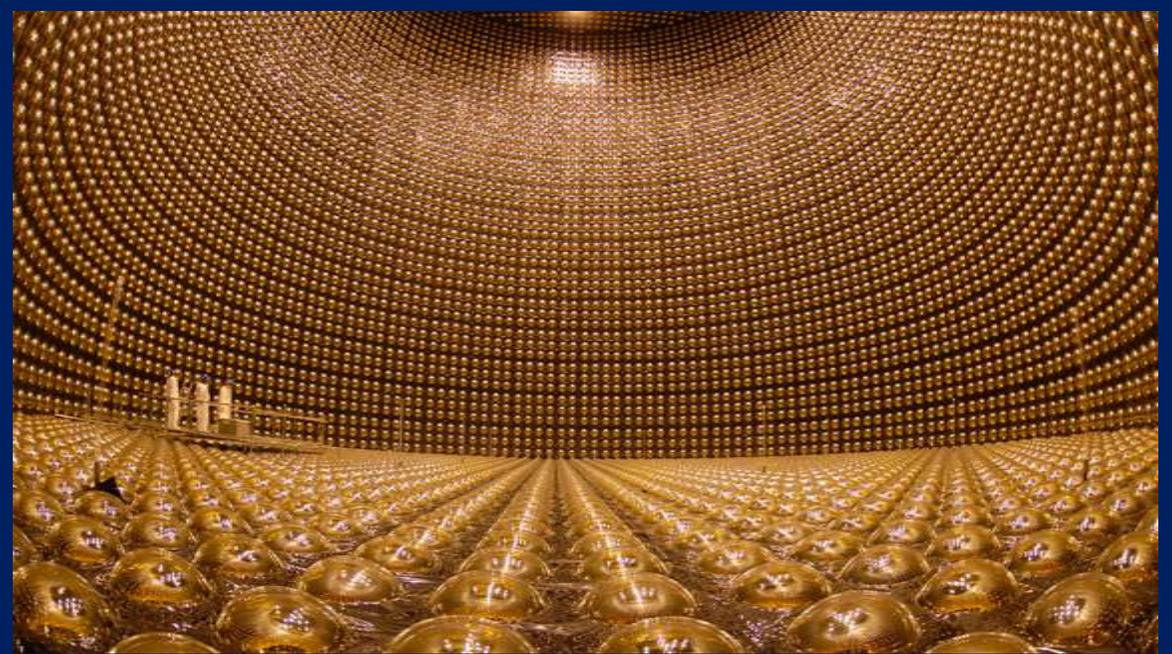
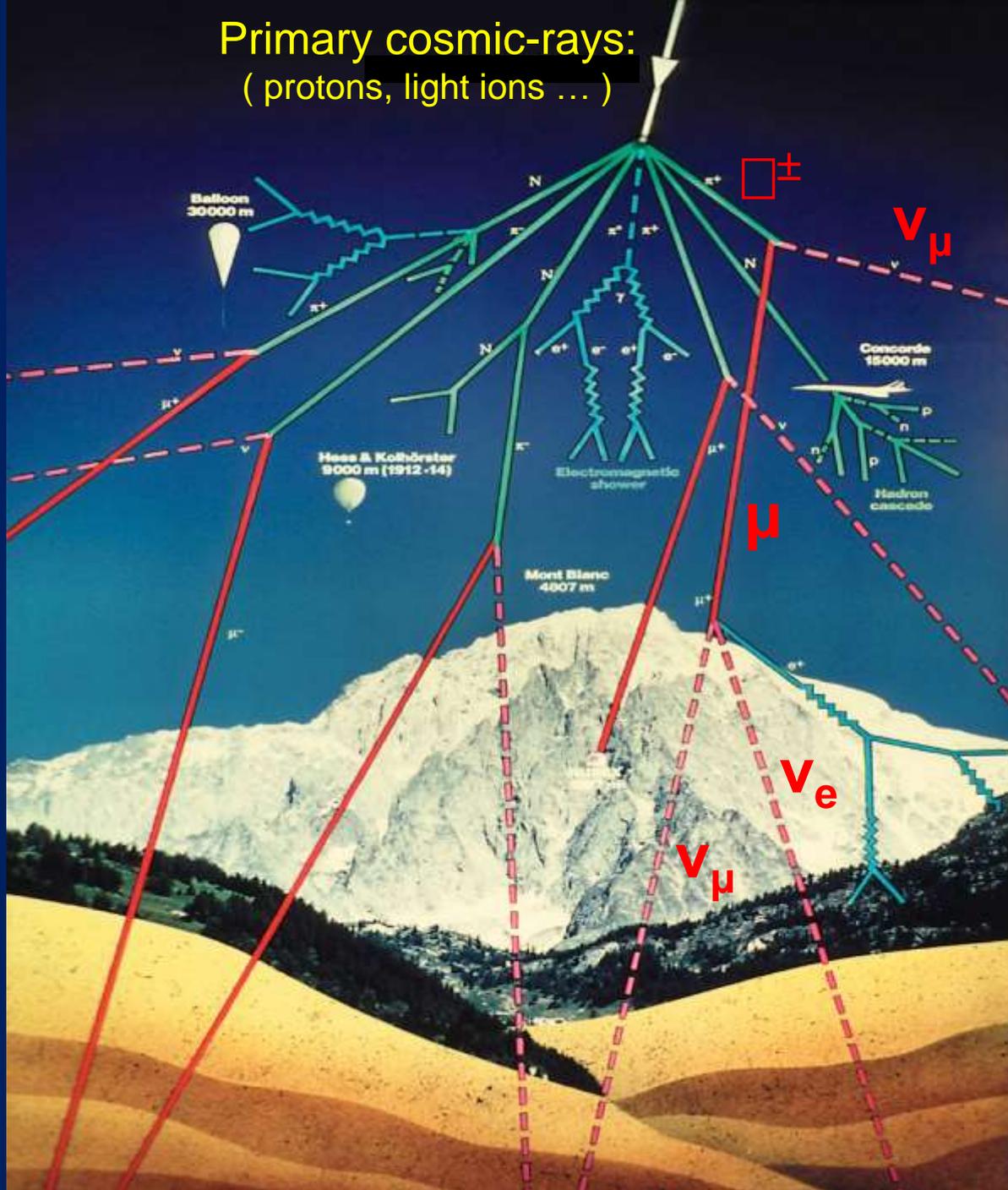


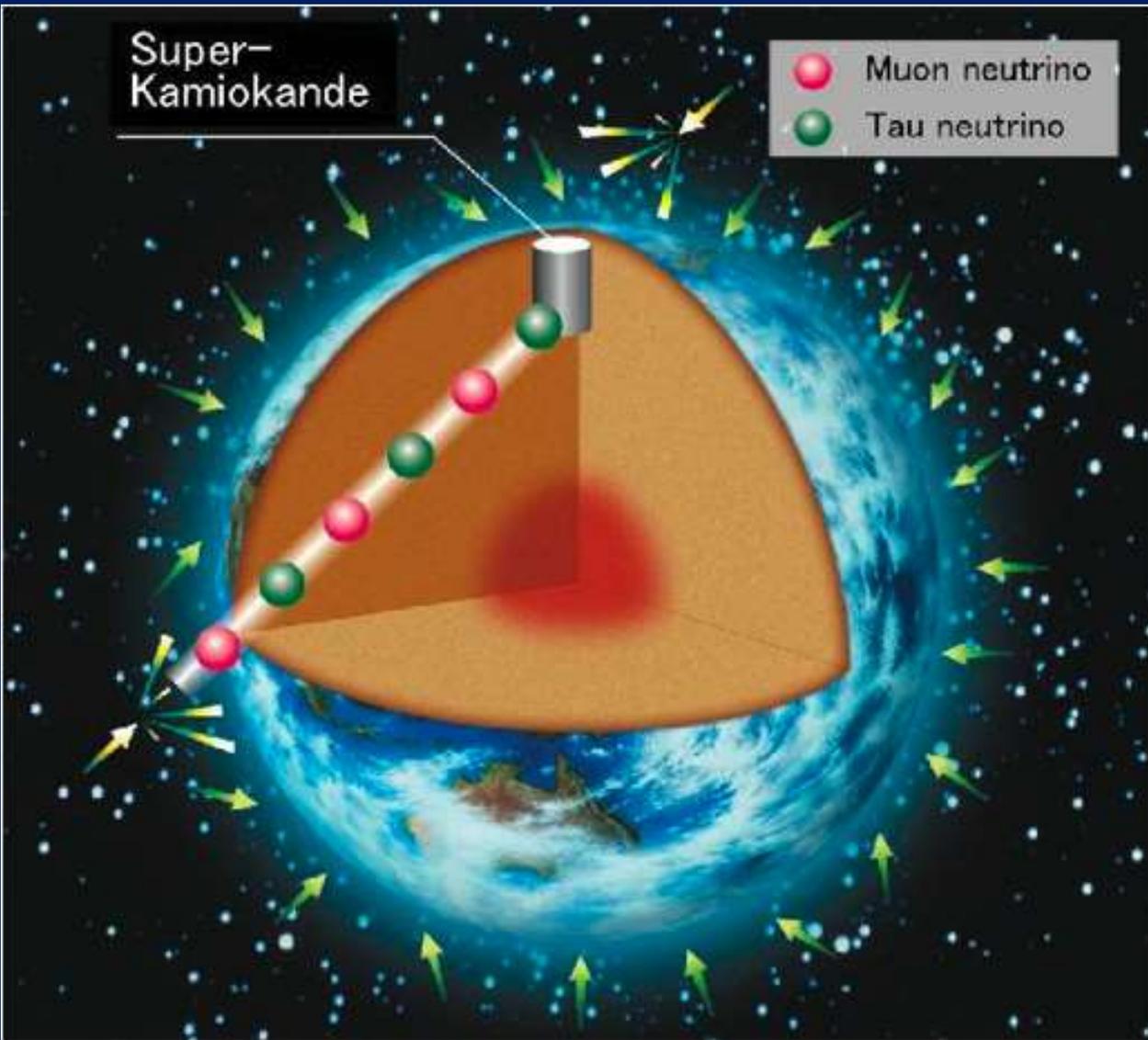
MicroBooNE

15 papers
33 public notes

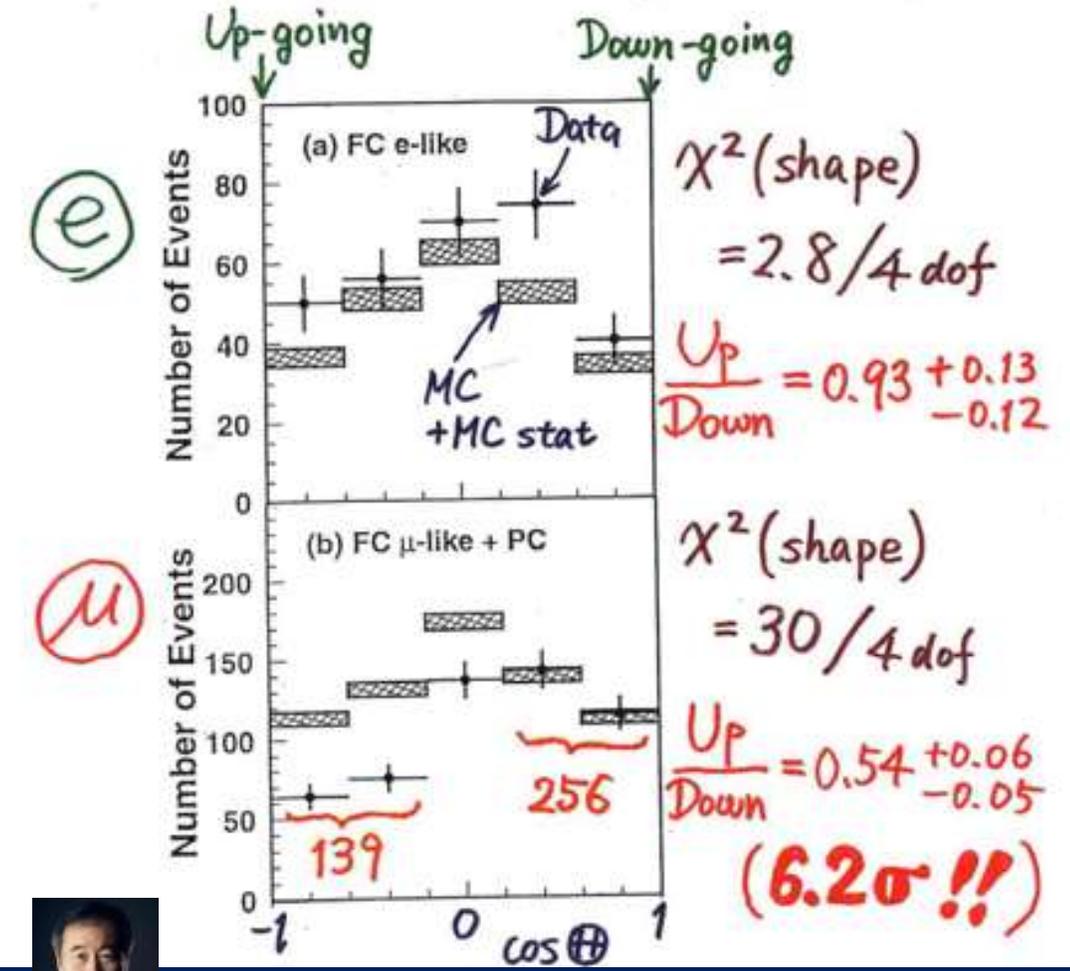


Primary cosmic-rays:
(protons, light ions ...)



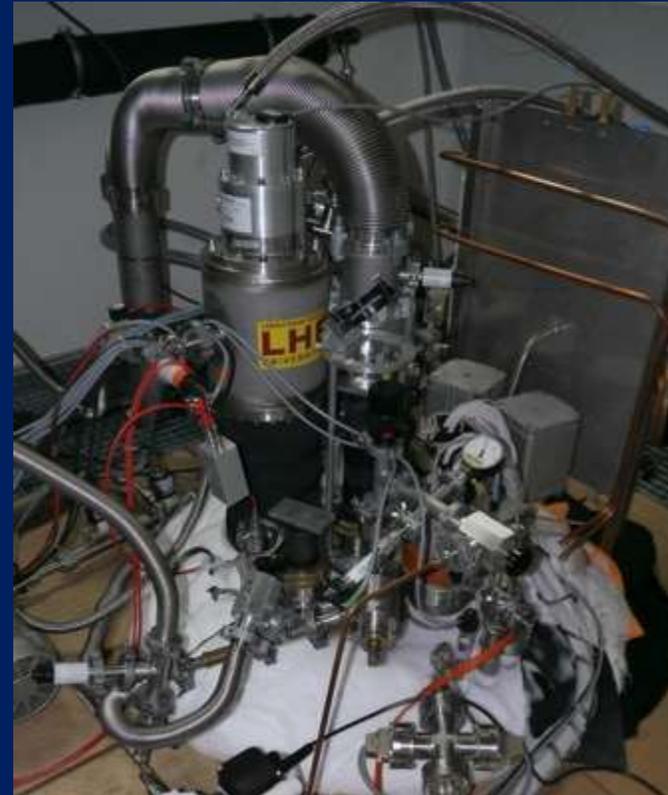


Zenith angle dependence (Multi-GeV)

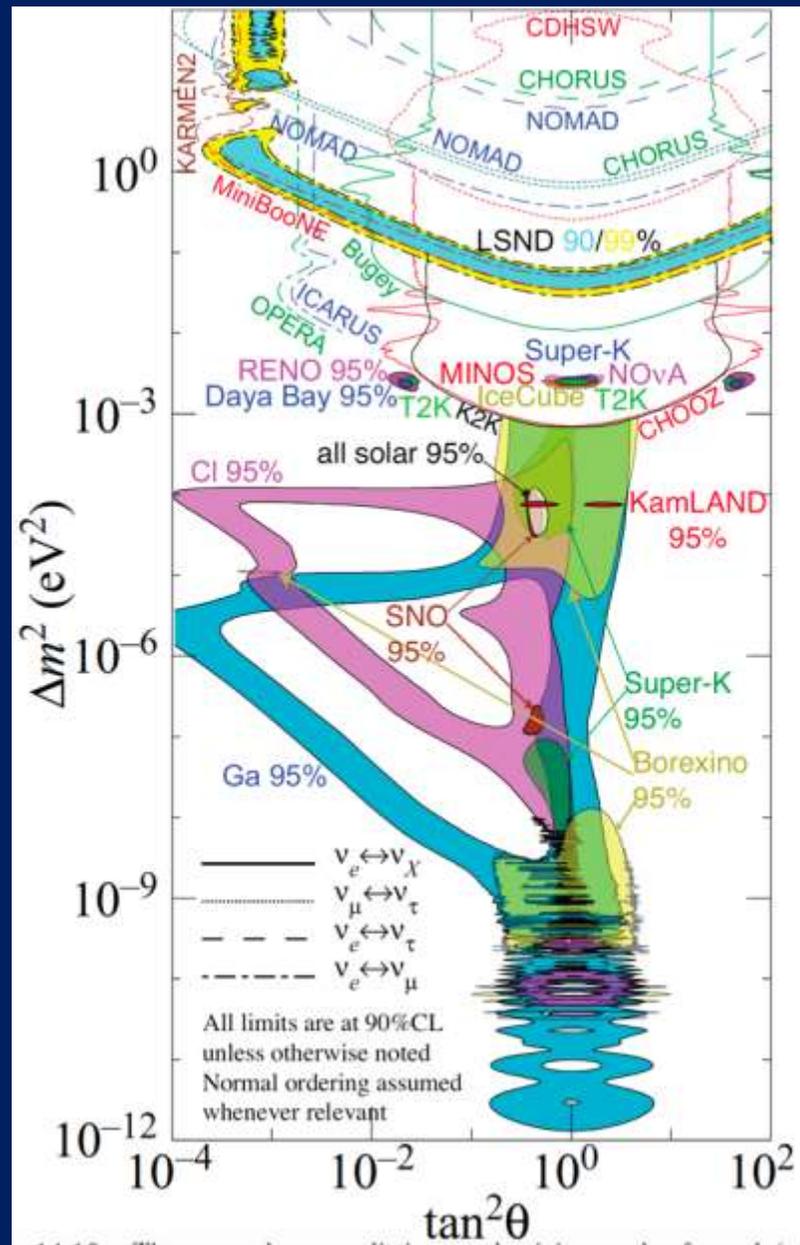


Takaaki Kajita, Neutrino '98

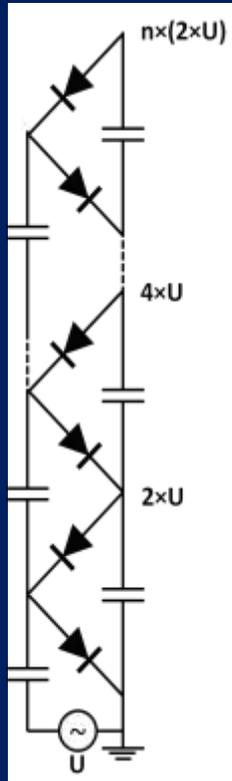
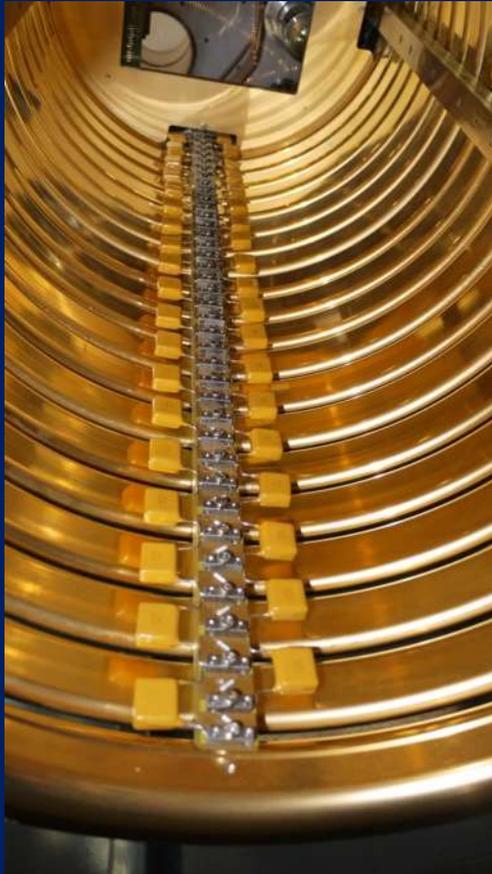
LArTPC development in Bern



ArgonTUBE
2009-2015

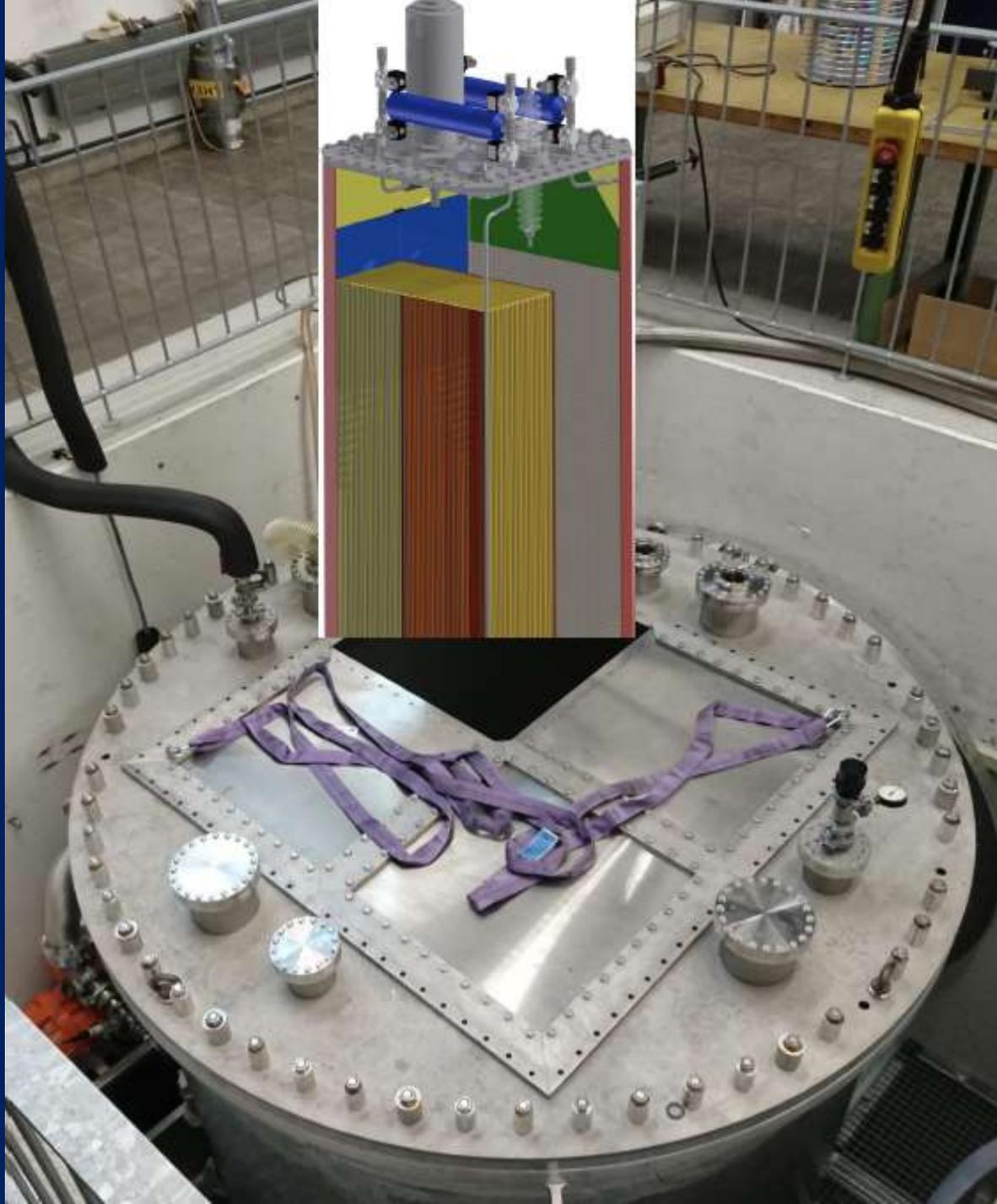


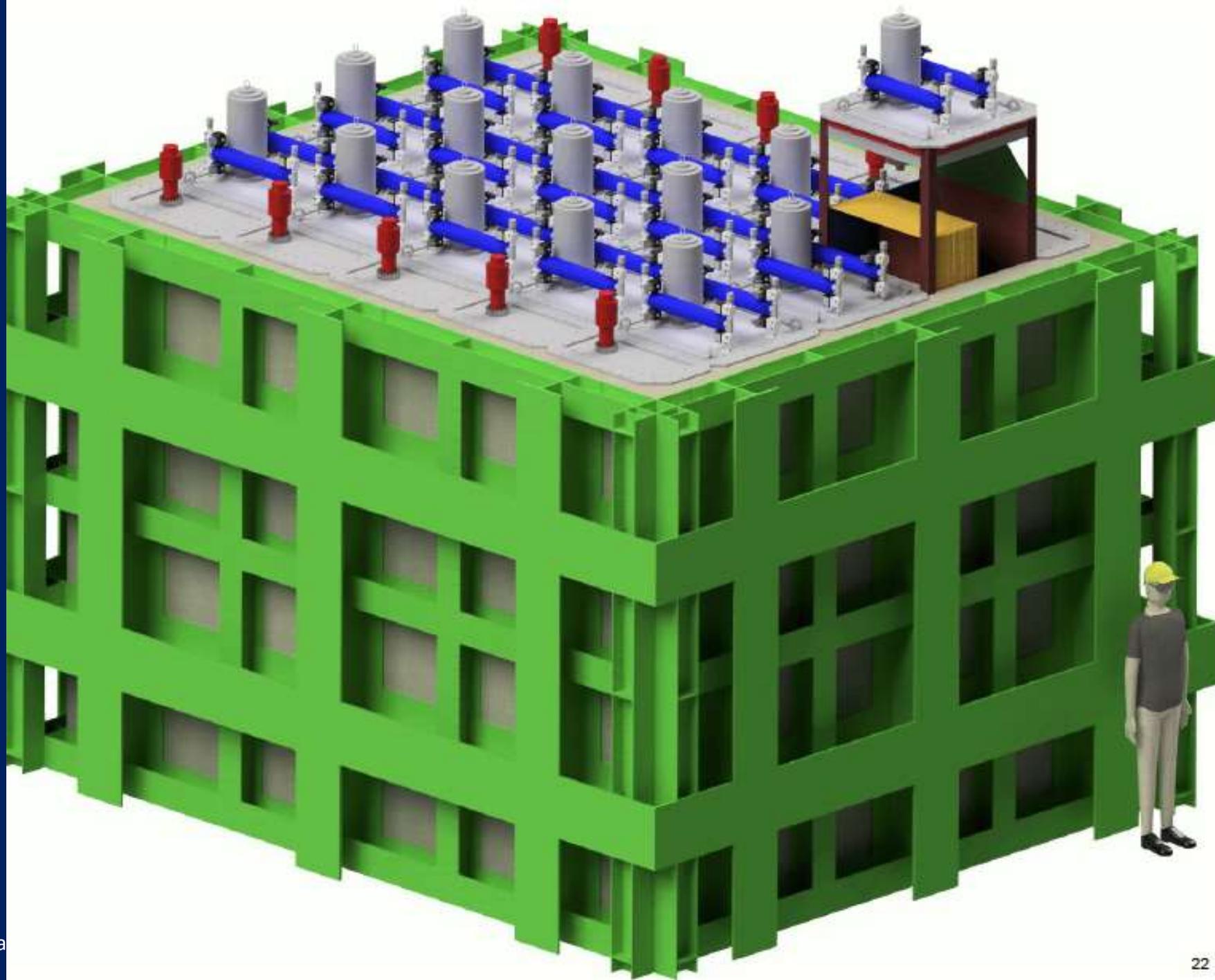
- Greinacher circuit: 125 stages, input 4 kV AC
- COMSOL finite element analysis software to optimize the geometry of the field-shaping rings
- Goal: drift field of 1 kV/cm
- Reached 170 kV (0.34 kV/cm)

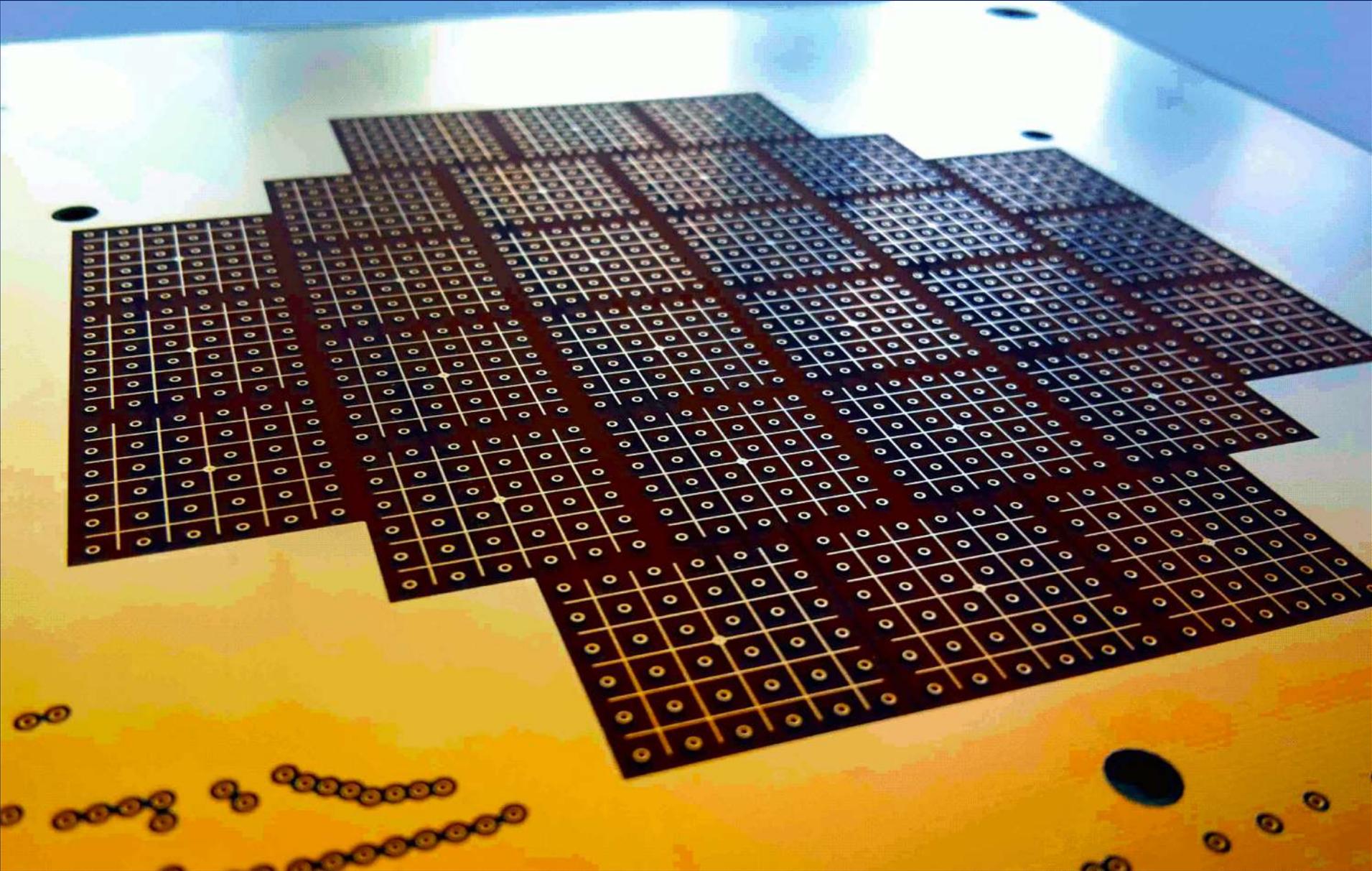


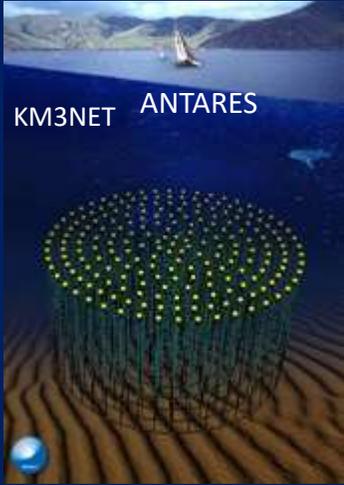
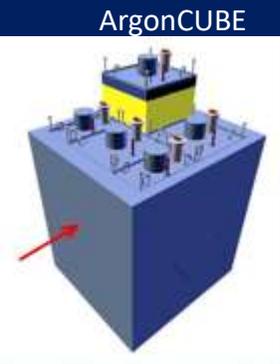
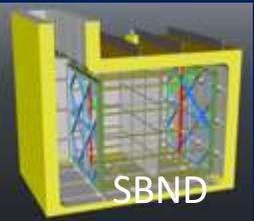
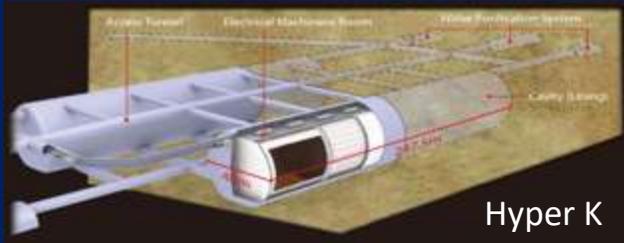
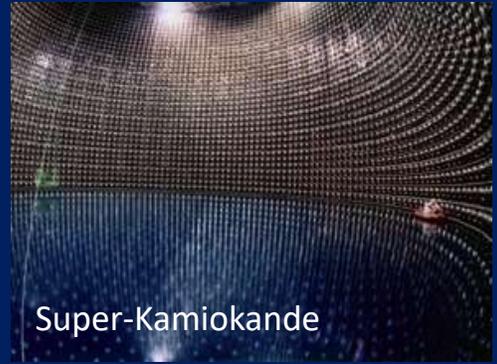
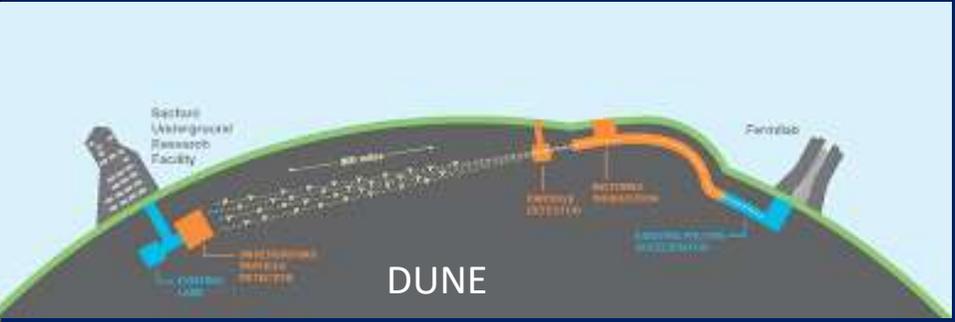
MicroBooNE in USA (2015 —)

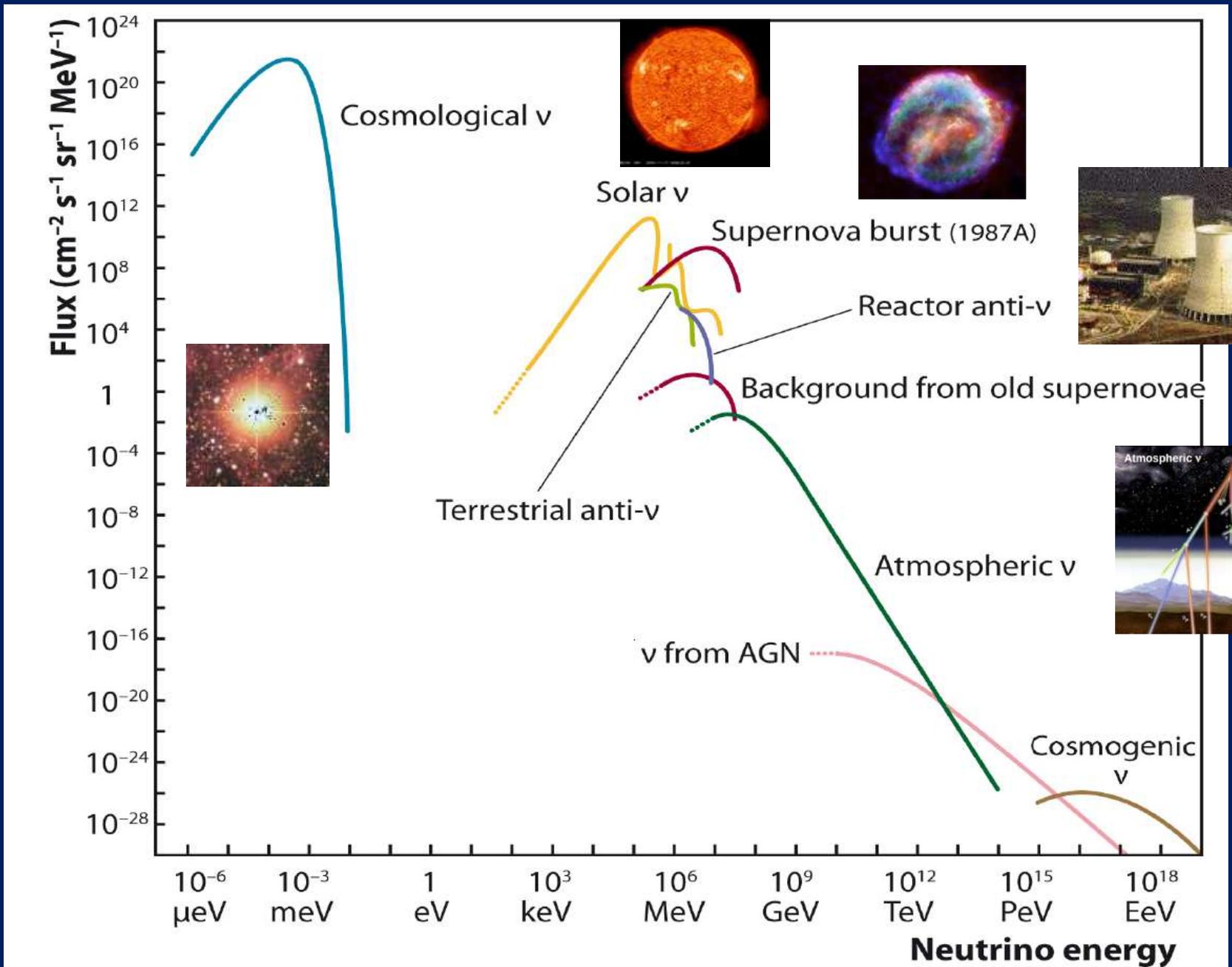


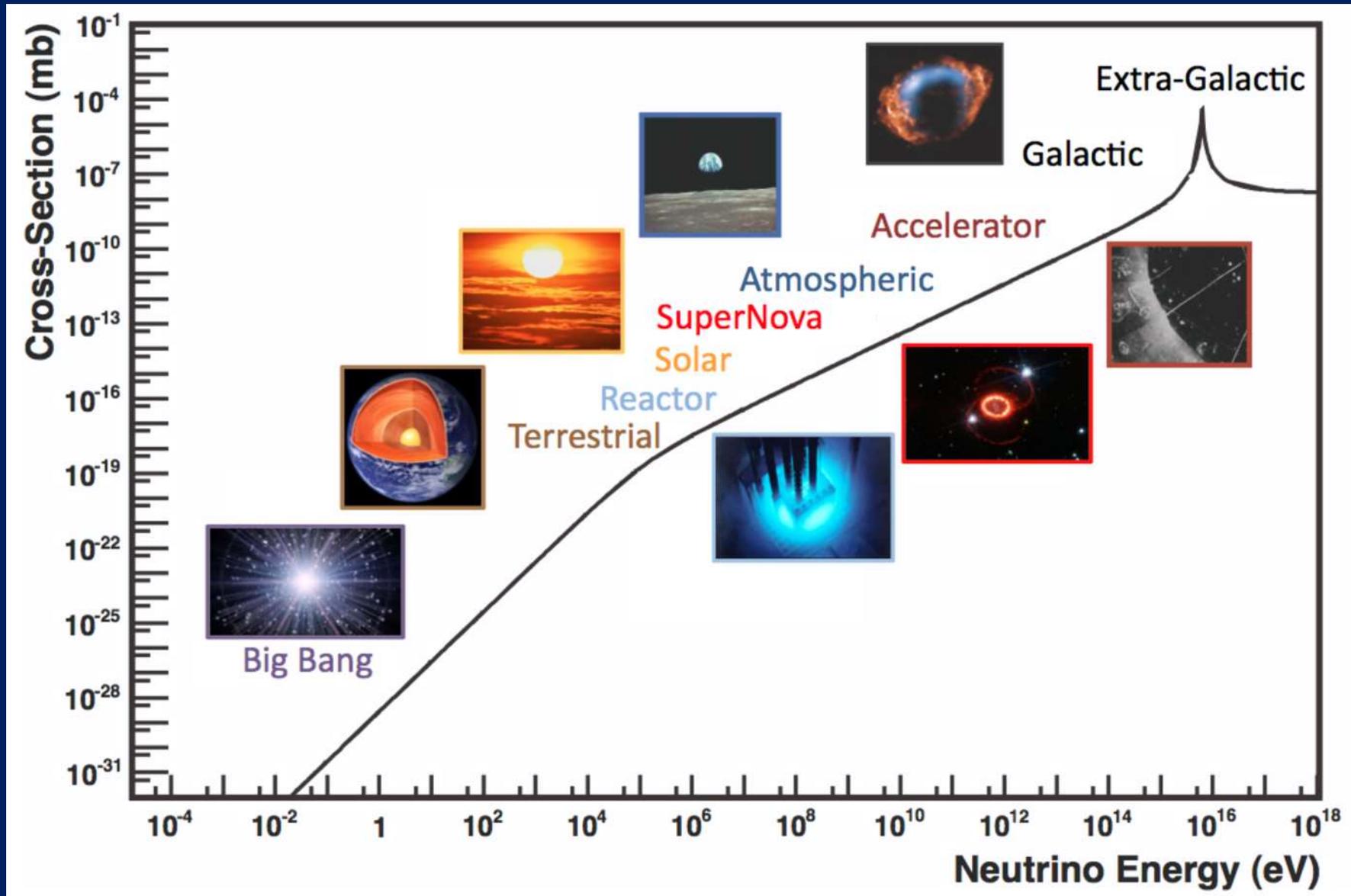




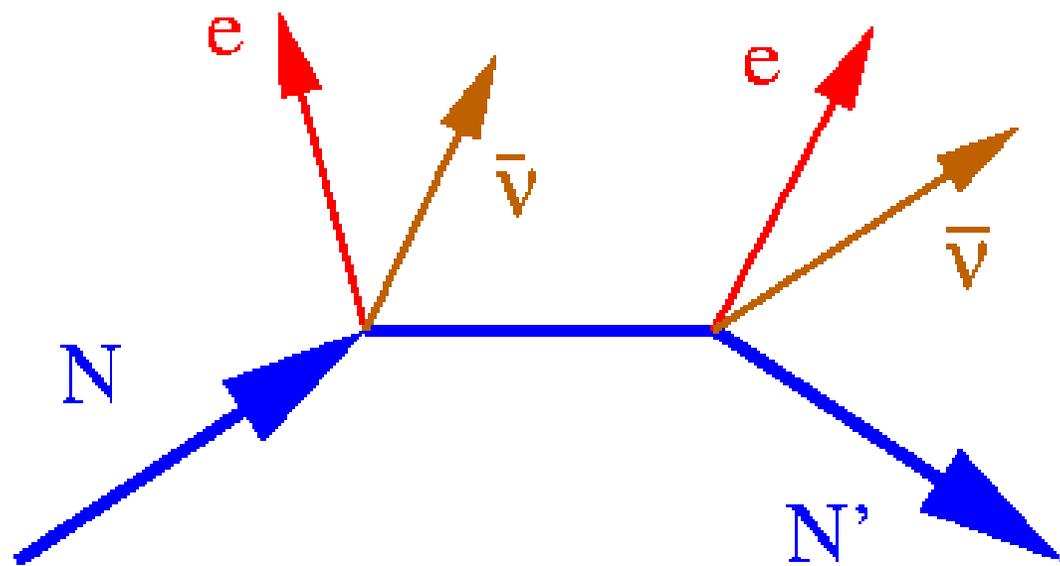




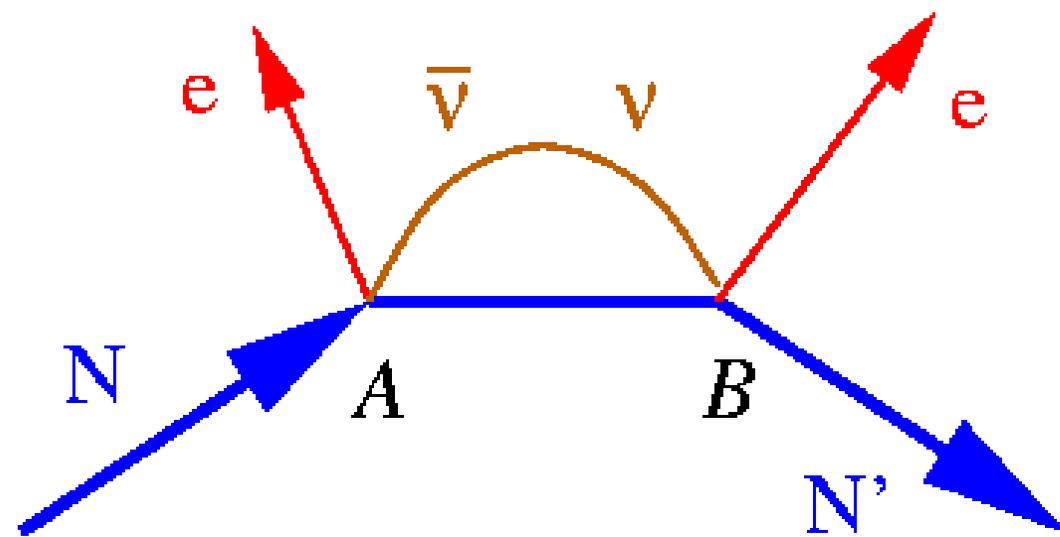




a) $2\nu \beta\beta$



b) $0\nu \beta\beta$



Overview of Karlsruhe TRitium Neutrino Experiment

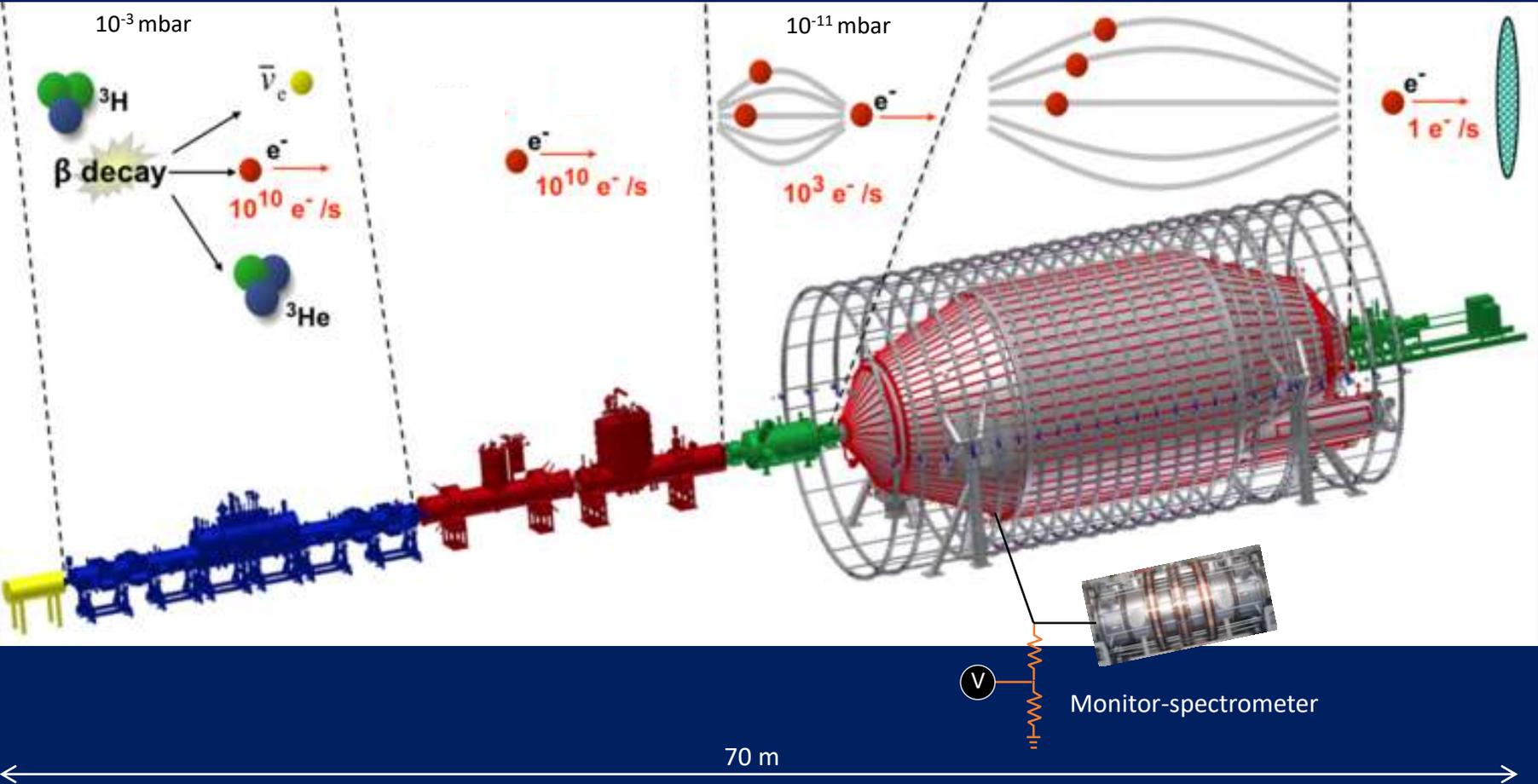
Windowless gaseous source

Transport section

Pre-spectrometer

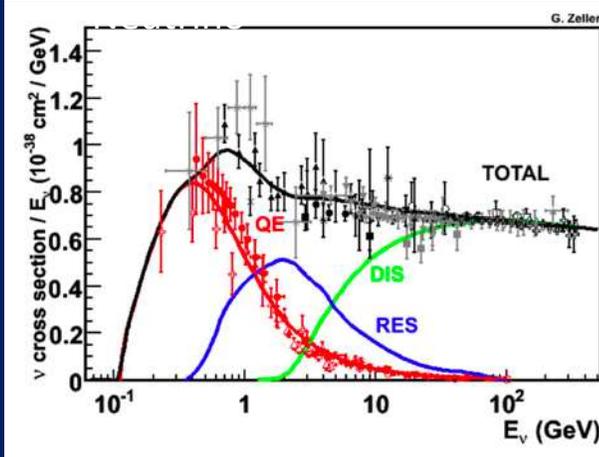
Main-spectrometer

Detector

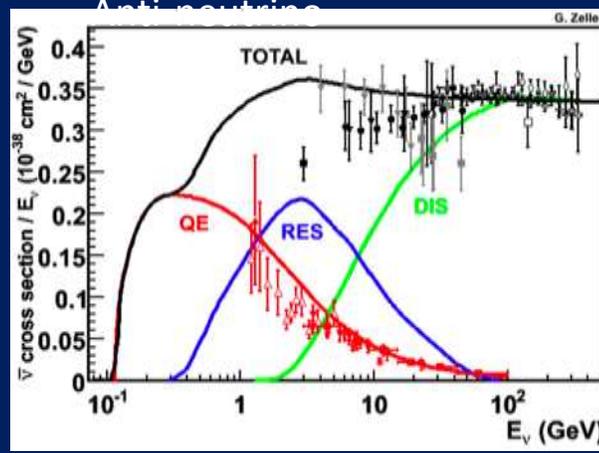


$$P(\nu_\mu \rightarrow \nu_\mu) \sim 1 - \cos^4(\theta_{13}) \sin^2(2\theta_{23}) \sin^2(\Delta m_{31}^2 L/4E) \\ - \sin^2(2\theta_{13}) \sin^2(\theta_{23}) \sin^2(\Delta m_{31}^2 L/4E)$$

Interaction regime

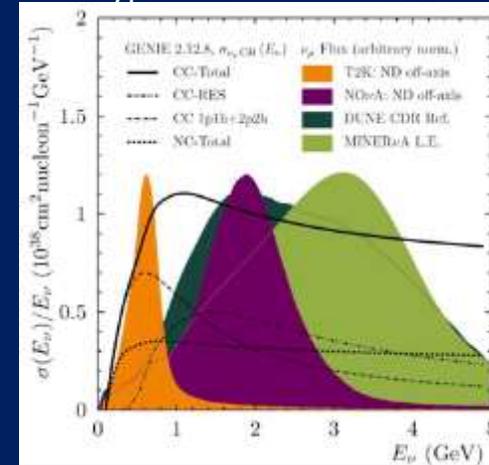


Rev. Mod. Phys. 84, 1307–1341 (2012)

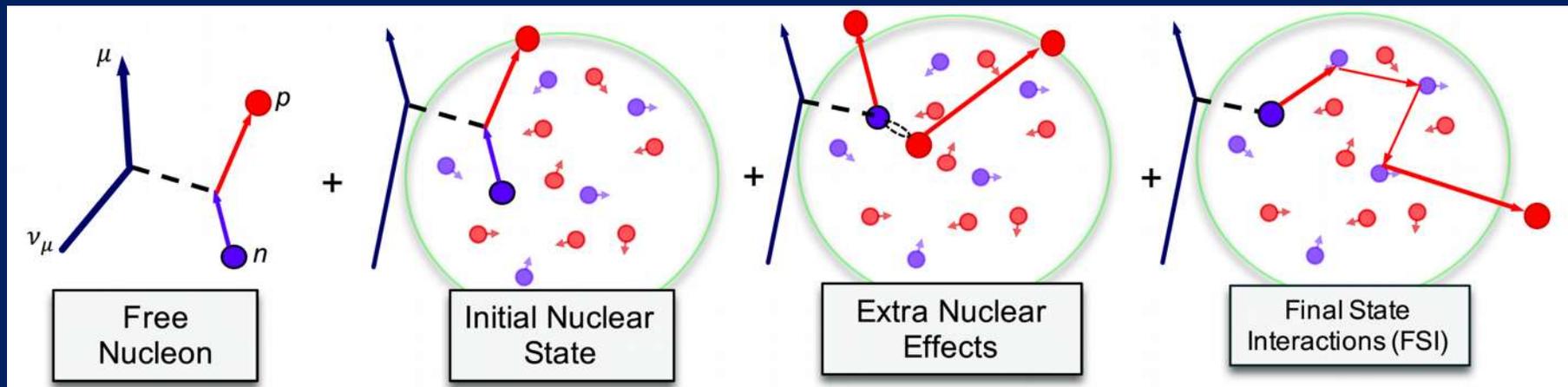


Anti-neutrino

Typical neutrino beams



Nucleus



3+1 neutrino oscillations

We extend the mixing matrix

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \Rightarrow \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

APPEARANCE

DISAPPEARANCE

$$P_{\alpha\beta}^{\text{SBL}} \approx \sin^2(2\theta_{\alpha\beta}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

↓

$$\sin^2(2\theta_{\alpha\beta}) = 4|U_{\alpha4}|^2|U_{\beta4}|^2$$

$$P_{\alpha\alpha}^{\text{SBL}} \approx 1 - \sin^2(2\theta_{\alpha\alpha}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

↙

$$\sin^2(2\theta_{\alpha\alpha}) = 4|U_{\alpha4}|^2(1 - |U_{\alpha4}|^2)$$

$$\nu_{\mu} \rightarrow \nu_e : \sin^2(2\theta_{\mu e}) = 4|U_{e4}|^2|U_{\mu4}|^2$$

$$\nu_e \rightarrow \nu_e : |U_{e4}|^2 = \sin^2 \theta_{14}$$

@Reactors and Gallium

@LSND, Karmen, MiniBooNE,
Opera

$$\nu_{\mu} \rightarrow \nu_{\mu} : |U_{\mu4}|^2 = \sin^2 \theta_{24} \cos^2 \theta_{14}$$

@atmospherics and accelerators

Mixing (rotation) matrix

$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} \\ -\sin\theta_{12} & \cos\theta_{12} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

$$|\nu_e(0)\rangle = \cos\theta_{12} \cdot |\nu_1\rangle + \sin\theta_{12} \cdot |\nu_2\rangle$$

Next, introduce time evolution (**propagation in space**)

$$|\nu_e(t)\rangle = \cos\theta_{12} \cdot e^{-iE_1 t} |\nu_1\rangle + \sin\theta_{12} \cdot e^{-iE_2 t} |\nu_2\rangle$$

$$P(n_a \rightarrow n_b) = \sin^2 2\theta_{ij} \cdot \sin^2 \left(1.27 D m_{ij}^2 \frac{L}{E} \right)$$